



How to measure the W Mass: A Theory Perspective

Joshua Isaacson In Collaboration with: Yao Fu and C.-P. Yuan Based on: arxiv:2205.02788 KIAS: W Mass Workshop 24 June 2022



Standard Model: W Mass

Standard Model EW Fit

$$\begin{split} M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right) &= \frac{\pi \alpha}{\sqrt{2}G_F} \left(1 + \Delta r\right) \\ \Delta r &= \Delta \alpha - \frac{c_W^2}{s_W^2} \Delta \rho + \Delta r_{\rm rem} \left(M_H\right) \,, \end{split}$$

where s_W^2 is the Weinberg angle, $\Delta \alpha$ is the correction to α from the light fermions, $\Delta \rho$ is the correction to the ρ parameter, and $\Delta r_{\rm rem}$ contains all corrections containing the Higgs mass.

Parameter	Fit Result
G_{μ} [GeV ⁻²]	1.1663787×10^{-5}
$\alpha(0)^{-1}$	137.035999139
$\Delta lpha_{\sf had}^{(5)}(M_Z^2)$	0.027627 ± 0.000096
M_Z [GeV]	91.1883 ± 0.0021
M_H [GeV]	125.21 ± 0.12
$m_t \; [{ m GeV}]$	172.75 ± 0.44
M_W [GeV]	80.3591 ± 0.0052

Table reproduced from: HEPFit Group (2112.07274).



Experimental Measurements

- CDF Run II results most precise
- 7σ tension with SM
- 3σ tension between CDF-II and ATLAS result
- $\bullet~$ Missing LHCb result: 80,354 $\pm~$ 32 MeV



Figure reproduced from CDF-II measurement (Science 376, 170).

Extracting W Mass from Data



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- Can't measure invariant mass directly due to neutrino
- Look at sensitive observables

•
$$M_T = \sqrt{2} \left(p_T^{\ell} p_T^{\nu} - \vec{p}_T^{\ell} \cdot \vec{p}_T^{\nu} \right)$$

• p_T^{ℓ}
• p_T^{ν} with $(\vec{p}_T^{\nu} = -\vec{p}_T^{\ell} - \vec{u}_T)$

- Requires precise theory calculation
- Fit theory templates with varying M_W



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Theory Calculation



Breakdown of Fixed Order

- Perturbative series has terms proportional to $\alpha_s^n \log^m\left(\frac{p_T^2}{M_W^2}\right)$, $m \leq 2n$
- As $p_T^W \to 0$ the series no longer converges
- Need to include corrections to all orders by resumming the series



Analytic vs. Numeric Resummation

Analytic:

- Formal resummation (focus here on *b*-space CSS resummation)
- Pros:
 - High precision and accuracy
- Cons:
 - Inclusive only
 - Numerically expensive
- $\bullet~{\rm Used}$ by CDF to obtain M_W

Numerical

- Parton Showers (Pythia, Sherpa, Herwig, Dire, Vincia)
- Pros:
 - Exclusive final states
 - Quick
- Cons:
 - Currently only LL with some subleading effects included
- ${\, \bullet \,}$ Used by ATLAS to obtain M_W



Parton Showers

Evolution Equation

$$\frac{df_a\left(x,t\right)}{d\ln t} = \sum_{b=q,g} \int_0^1 \frac{dz}{z} \frac{\alpha_s}{2\pi} \left[P_{ab}\left(z\right)\right]_+ f_b\left(\frac{x}{z},t\right)$$

- $f_{a}\left(x,t
 ight)$ is the obervable being evolved
- $P_{ab}\left(z
 ight)$ is the evolution (splitting) kernel
- Solve using Markovian Monte-Carlo algorithms
- Treat P_{ab} as a probability
- Virtual corrections defined at kinematic endpoints by + prescription

 \overline{d}

Resummation

$$\frac{d\sigma_{\rm res}}{Q^2 d^2 \vec{q_T} dy d\Omega} = \sigma \int \frac{d^2 b}{(2\pi)^2} e^{i \vec{q_T} \cdot \vec{b}} \tilde{W},$$

$$\tilde{W} = e^{-S(b)} C \otimes f(x_A, C_3/b) C \otimes f(x_B, C_3/b)$$

$$S(b) = \int_{\frac{C_1^2}{b^2}}^{C_2^2 Q^2} \frac{d\bar{\mu}^2}{\bar{\mu}^2} \left[\ln\left(\frac{C_2^2 Q^2}{\bar{\mu}^2}\right) A(\bar{\mu}) + B(\bar{\mu}) \right]$$

[Collins, Soper, Sterman, '85] [...]



Resummation

$$\frac{d\sigma_{\text{res}}}{dQ^2 d^2 \vec{q_T} dy d\Omega} = \sigma \int \frac{d^2 b}{(2\pi)^2} e^{i \vec{q_T} \cdot \vec{b}} \tilde{W},$$
$$\tilde{W} = \left\langle e^{-S(b)} \quad C \otimes f(x_A, C_3/b) C \otimes f(x_B, C_3/b) \right\rangle$$
$$S(b) \neq \int_{\frac{C_1^2}{b^2}}^{C_2^2 Q^2} \frac{d\bar{\mu}^2}{\bar{\mu}^2} \left[\ln \left(\frac{C_2^2 Q^2}{\bar{\mu}^2} \right) A(\bar{\mu}) + B(\bar{\mu}) \right]$$

• Electroweak cross section /

[Collins, Soper, Sterman, '85] [...]



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- Electroweak cross section
- Sudakov factor /



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- Sudakov factor
- Collinear factors -

Resummation

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• Electroweak cross section
• Sudakov factor
• Collinear factors
• Perturbative Coefficients (A, B, C)

[Collins, Soper, Sterman, '85] [...]

J. Isaacson



		Anomalous D	imension	
Order	Boundary Condition	γ_i (non-cusp)	Γ_{cusp}, β	Fixed Order Matching
LL	1	-	1-loop	-
NLL	1	1-loop	2-loop	-
NLL' (+ NLO)	α_s	1-loop	2-loop	α_s
NNLL $(+ NLO)$	α_s	2-loop	3-loop	$lpha_s$
NNLL' (+ NNLO)	α_s^2	2-loop	3-loop	α_s^2
$N^{3}LL (+ NNLO)$	α_s^2	3-loop	4-loop	α_s^2
$N^{3}LL'(+ N^{3}LO)$	$\alpha_s^{\bar{3}}$	3-loop	4-loop	$\alpha_s^{\bar{3}}$
$N^4LL (+ N^3LO)$	α_s^3	4-loop	5-loop	α_s^{3}

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● ■ Accuracy used by CDF

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- ■ Accuracy used by CDF
- Current accuracy available in ResBos code

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- ■ Accuracy used by CDF
- Current accuracy available in ResBos code
- ■ All terms known to this accuracy



Non-Perturbative Fit

$$S(b) = \int_{\frac{C_1^2}{\bar{\mu}^2}}^{C_2^2 Q^2} \frac{d\bar{\mu}^2}{\bar{\mu}^2} \left[\ln\left(\frac{C_2^2 Q^2}{\bar{\mu}^2}\right) A(\bar{\mu}) + B(\bar{\mu}) \right]$$

- Lower limit goes to zero as b goes to infinity
- Requires evaluation of $\alpha_s(C_1/b)$ which is non-perturbative
- Need to introduce a non-perturbative cutoff (*b**-prescription):

$$b^* = \frac{b}{\sqrt{1 + \frac{b^2}{b_{\max}^2}}}$$

BLNY Form

$$S_{NP}(b) = -b^2 \left(g_1 + g_2 \log \left(\frac{Q}{2Q_0} \right) + g_1 g_3 \log(100x_1 x_2) \right)$$



- g_2 tuned to reproduce CDF-II p_T^Z
- M_W vs. M_Z captured in Q dependence
- No flavor dependence included
- No consideration of uncertainty from changing form, but expected to be small



NOTE: SIYY2 is the same functional form as BLNY, but with $b_{max} = 1.5 \text{ GeV}^{-1}$



Flavor Dependence

- Study on flavor dependence for $\sqrt{s}=7~{\rm TeV}~{\rm LHC}$
- $S_{NP}(b) = -b^2(g_a + g_{evo}\log(Q^2/Q_0^2))$, where g_a is the flavor dependent piece
- Found shift could be up to 10 MeV
- Additional studies are required to validate
- Unclear what the global shift would be

Set	u_v	d_v	u_s	d_s	others
1	0.34	0.26	0.46	0.59	0.32
2	0.34	0.46	0.56	0.32	0.51
3	0.55	0.34	0.33	0.55	0.30
4	0.53	0.49	0.37	0.22	0.52
5	0.42	0.38	0.29	0.57	0.27

Set	ΔM_W^+		ΔM_W^-	
	M_T	p_T^ℓ	M_T	p_T^ℓ
1	0	-1	-2	3
2	0	-6	-2	0
3	-1	9	-2	-4
4	0	0	-2	-4
5	0	4	-1	-3

Table reproduced from: Phys. Letters B 788 (2019) 542-545

Speed Improvements



$$\int_{0}^{\infty} dx f(x) J_{n}(x) = \pi \sum_{j=1}^{N} w_{nj} f\left(\frac{\pi}{h} \psi\left(x_{nj}\right)\right) J_{n}\left(\frac{\pi}{h} \psi\left(x_{nj}\right)\right) \psi'\left(x_{nj}\right) + \left[I_{nN+1} + \mathcal{O}\left(e^{-c/h}\right)\right],$$
$$\psi(t) = t \tanh\left(\frac{\pi}{2}\sinh\left(t\right)\right), \quad x_{nj} = h\xi_{nj}, \quad w_{nj} = \frac{2}{\pi^{2}\xi_{nj}J_{n+1}\left(\pi\xi_{nj}\right)}$$

[1906.05949]

$$\int_{0}^{\infty} dx f(x) J_{n}(x) = \pi \sum_{j=1}^{N} w_{nj} f\left(\frac{\pi}{h}\psi\left(x_{nj}\right)\right) J_{n}\left(\frac{\pi}{h}\psi\left(x_{nj}\right)\right) \psi'\left(x_{nj}\right) + \begin{bmatrix} I_{nN+1} + \mathcal{O}\left(e^{-c/h}\right) \end{bmatrix},$$
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• Truncation error

[1906.05949]

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$$\bullet \text{ Truncation error}$$
$$\bullet \text{ Finite step size error}$$



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$$\bullet \text{ Truncation error}$$
$$\bullet \text{ Finite step size error}$$
$$\bullet \text{ Evaluation points}$$

$$\int_{0}^{\infty} dx f(x) J_{n}(x) = \pi \sum_{j=1}^{N} w_{nj} f\left(\frac{\pi}{h} \psi(x_{nj})\right) J_{n}\left(\frac{\pi}{h} \psi(x_{nj})\right) \psi'(x_{nj}) + \left[I_{nN+1} + \mathcal{O}\left(e^{-c/h}\right)\right],$$
$$\psi(t) = t \tanh\left(\frac{\pi}{2}\sinh(t)\right), \quad x_{nj} = h\xi_{nj}, \quad w_{nj} = \frac{2}{\pi^{2}\xi_{nj}J_{n+1}(\pi\xi_{nj})}$$
$$\text{ Truncation error}$$
$$\text{ Finite step size error}$$
$$\text{ Evaluation points}$$
$$\text{ Quadrature weights}$$

[1906.05949]

Ogata Quadrature









Angular Coefficients



lepton plane

$$\frac{\mathrm{d}\sigma}{\mathrm{d}p_{\mathrm{T}}^2\,\mathrm{d}y^2\,\mathrm{d}m^2\,\mathrm{d}\cos\theta\,\mathrm{d}\phi} = \frac{3}{16\pi}\frac{\mathrm{d}\sigma^{U+L}}{\mathrm{d}p_{\mathrm{T}}^2\,\mathrm{d}y^2\,\mathrm{d}m^2}$$

$$\left\{ (1+\cos^2\theta) + \frac{1}{2}\,A_0(1-3\cos^2\theta) + A_1\,\sin2\theta\,\cos\phi + \frac{1}{2}\,A_2\,\sin^2\theta\,\cos2\phi + A_3\,\sin\theta\,\cos\phi + A_4\,\cos\theta + A_5\,\sin^2\theta\,\sin2\phi + A_6\,\sin2\theta\,\sin\phi + A_7\,\sin\theta\,\sin\phi \right\},$$

$$\langle P(\cos\theta,\phi)\rangle = \frac{\int P(\cos\theta,\phi) d\sigma(\cos\theta,\phi) d\cos\theta d\phi}{\int d\sigma(\cos\theta,\phi) d\cos\theta d\phi}.$$

$$\begin{split} &\langle \frac{1}{2}(1-3\cos^2\theta)\rangle = \frac{3}{20}(A_0-\frac{2}{3}); \quad \langle\sin 2\theta\cos\phi\rangle = \frac{1}{5}A_1; \quad \langle\sin^2\theta\cos2\phi\rangle = \frac{1}{10}A_2; \\ &\langle\sin\theta\cos\phi\rangle = \frac{1}{4}A_3; \quad \langle\cos\phi\rangle = \frac{1}{4}A_4; \quad \langle\sin^2\theta\sin2\phi\rangle = \frac{1}{5}A_5; \\ &\langle\sin2\theta\sin\phi\rangle = \frac{1}{5}A_6; \quad \langle\sin\theta\sin\phi\rangle = \frac{1}{4}A_7. \end{split}$$

- Well known issue with angular coefficients in the ResBos code at NNLO (No issue with matching to NLO)
- CDF-II only used the NLO so the angular functions are exact to that order
- ResBos only included NNLO corrections to the total rate, but not to the angular functions
- This is an issue with matching to an incomplete NNLO calculation, and not an issue with the resummation or the matching to fixed order
- Only effects larger $p_T \ (p_T > 30 \ {\rm GeV}, \ {\rm CDF} \ {\rm has} \ {\rm a} \ {\rm cut} \ {\rm of} \ p_T < 15 \ {\rm GeV})$
- Has been resolved via matching to MCFM (preliminary results next slides)



NOTE: Uncertainties are purely statistical for ResBos + MCFM



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Results



$P_T(Z)/P_T(W)$



- Ratio is stable to higher order corrections at small p_T
- Scale uncertainty only using correlated prediction
- Need to investigate the CDF estimated uncertainty from this ratio



Methodology

Our Procedure:

- Generate pseudodata using N³LL+NNLO prediction
- Tune NNLL+NLO prediction to reproduce $p_T(Z)$ data
- Validate tuned result against $p_T(W)$ data
- Use tuned result to generate mass templates
- Extract W mass from template fit for each observable
- Calculate the mass shift from the input value for pseudodata

Details:

- Pseudodata $M_W = 80,358 \text{ MeV}$
- Cuts:
 - $p_T(Z) < 15 \text{ GeV}, p_T(W) < 15 \text{ GeV}$
 - $30 < p_T(\ell) < 55 GeV$, $30 < p_T(\nu) < 55 \text{ GeV}$
 - $|\eta(\ell)| < 1$
 - $66 < M_{\ell\ell} < 116$ GeV (Z events), $60 < m_T < 100$ GeV (W events)
- Number of Events:
 - 1,811,700 $W \to e \nu$
 - 66,180 $Z \rightarrow ee$
 - 2,424,486 $W
 ightarrow \mu
 u$
 - 238,534 $Z
 ightarrow \mu \mu$



Tuning to Pseudodata



Tuned result:

- Fit to $p_T(Z) < 15~{\rm GeV}$
- $g_2 = 0.662 \text{ GeV}^2$



• $\alpha_S(M_Z) = 0.120$

• Tuned PDF set: CT18NNL0_as_120



ResBos and Resummation

Results







Best Fit: $M_W = 80,386$ MeV

Best Fit: $M_W = 80,388$ MeV

Best Fit: $M_W = 80,389$ MeV

	Mass Shift [MeV]			
Observable	ResBos2	+Detector Effect+FSR		
m_T	1.5 ± 0.5	$0.2 \pm 1.8 \pm 1.0$		
$p_T(\ell)$	3.1 ± 2.1	$4.3 \pm 2.7 \pm 1.3$		
$p_T(u)$	4.5 ± 2.1	$3.0 \pm 3.4 \pm 2.2$		



Future Studies

- $\bullet\,$ Investigate effect of non-perturbative functional form on M_W
- Investigate flavor dependence effects on ${\cal M}_W$ extraction
- Perform detailed study on the $p_T(Z)/p_T(W)$ ratio and its impact on the M_W uncertainty
- Work with experimentalists to better understand detector smearing
- Understand how to properly combine the three observables



Conclusions

- CDF used ResBos code at NNLL+NLO accuracy
- ResBos v2 is able to go to $N^3LL+NNLO$ accuracy
- ResBos2 corrected major criticism of incorrect angular functions in the ResBos code
- Mimic CDF analysis using pseudoexperiments at N $^{3}LL+NNLO$ accuracy
- Find shift to be consistent with 0 MeV and up to 10 MeV (2σ) .



Backup





Detector Smearing:

• Fit functional form (Smearing 1): σ b c

$$\overline{E} = a \oplus \overline{\sqrt{E}} \oplus \overline{E}$$

- Use gaussian with 5%(11%) width for $\ell(\nu)$ (Smearing 2)
- Note results not sensitive to smearing effect if data and theory smeared identically

	Mass Shift [MeV]			
Observable	Smearing 1 Smearing 2			
m_T	$0.2\pm1.8\pm1.0$	$1.0\pm2.1\pm1.3$		
$p_T(\ell)$	$4.3\pm2.7\pm1.3$	$4.5 \pm 2.6 \pm 1.4$		
$p_T(u)$	$3.0\pm3.4\pm2.2$	$3.8\pm4\pm2.7$		

Width Effect:

- Central width: $\Gamma_W = 2.0895 \text{ GeV}$
- NLO width proportional to ${\cal M}^3_W$
- Negligible shift

Width	Mass Shift [MeV]
2.0475 GeV	2.0 ± 0.5
2.1315 GeV	0.3 ± 0.5
NLO	1.2 ± 0.5



	m	n_T	$p_T(\ell)$		$p_T(\ell)$ $p_T(u)$		r(u)
PDF Set	NNLO	NLO	NNLO	NLO	NNLO	NLO	
CT18	0.0 ± 1.3	1.8 ± 1.2	0.0 ± 15.9	2.0 ± 14.3	0.0 ± 15.5	2.9 ± 14.2	
MMHT2014	1.0 ± 0.6	2.6 ± 0.6	6.2 ± 7.8	36.7 ± 7.0	3.9 ± 7.5	36.0 ± 6.7	
NNPDF3.1	1.1 ± 0.3	2.1 ± 0.4	2.1 ± 3.8	13.5 ± 4.9	5.4 ± 3.7	10.0 ± 4.9	
CTEQ6M	N/A	2.8 ± 0.9	N/A	19.0 ± 10.4	N/A	20.9 ± 10.2	

- Central value is shift from 80,385 MeV
- Uncertainty is the PDF uncertainty for the given set
- Need to combine to compare to 3.9 MeV from CDF
- Rough estimates say it is consistent with CDF

PDF Correlations



- PDF-induced correlation of M_W and CT18 NNLO error set vs. x at $Q=100~{\rm GeV}$
- Region around $x = \frac{M_W}{\sqrt{s}}$ dominated by \bar{d}/\bar{u} , d/u and d PDFs

