Sensitivities to secret neutrino interaction at meson decay and short baseline neutrino experiments

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Overview

Introduction

- 2 Pion and Kaon rare decay experiments
- 3 Short-baseline neutrino experiments
- A new meson decay mode





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Introduction

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Physics beyond SM

- Neutrinos are massive, therefore we need to extend the standard model
- We can introduce a new gauge interaction in which neutrinos are the only charged particles in contrast to electromagnetic interaction
- So called secret neutrino interaction
- $g_{\alpha\beta}Z'_{\mu}\bar{\nu}_{\alpha}\gamma^{\mu}\nu_{\beta}$
- Motivated by νDM models

Chared mesons decay experiments

- Test of the standard model and beyond the standard model
- Leptonic meson decay is a two-body decay, and monochromatic, and is suppressed by m_l^2/m_M^2
- Leptonic meson decay is used to test lepton flavor universality
- Used to constrain new neutrino interactions
- We investigate the sensitivity of the leptonic decay of the charged meson to the new gauge interaction

Short-Baseline neutrino experiments

- Measurement of neutrino interaction cross-section
- Measurement of neutrino flux
- The standard model measurements
- Test of beyond three neutrino oscillation scheme such as sterile neutrino, Non-standard neutrino interaction, etc
- We investigate the sensitivity of short baseline neutrino experiments to the new gauge boson



$$\mathcal{M} = G_F f_M V_{qq'} p_\rho (u_1 \gamma^\rho P_L v_2) \tag{1}$$

$$P_L \text{ is chirality projection matrix, } P_L = \frac{1-\gamma^5}{2}$$

$$f_M \text{ is meson structure constant}$$

$$V_{qq'} \text{ is element of the CKM matrix}$$

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$$|\mathcal{M}|^2 = G_F^2 f_M^2 V_{qq'}^2 p_\rho p_\sigma(u_1 \gamma^\rho P_L v_2) (v_2 P_R \gamma^\sigma u_1)$$
⁽²⁾

$$\sum_{spins} |\mathcal{M}|^2 = G_F^2 f_M^2 V_{qq'}^2 p_\rho p_\sigma \operatorname{Tr}(\gamma^\rho P_L k_2 P_R \gamma^\sigma (k_1 + m_\mu))$$
(3)

$$\sum_{spins} |\mathcal{M}|^2 = G_F^2 f_M^2 V_{qq'}^2 \operatorname{Tr}(\not p \not k_2 P_R \not p \not k_1) = G_F^2 f_M^2 V_{qq'}^2 m_M^2 m_l^2 (1 - \frac{m_l^2}{m_M^2}) \quad (4)$$

$$\Gamma(M \longrightarrow l_{\alpha}\nu_{\alpha}) = \frac{G_F^2 f_M^2 V_{qq'}^2}{4\pi m_M^3} m_l^2 (m_M^2 - m_l^2)^2$$
(5)

 G_F is Fermi constant, f_M is Meson decay constant, and V_{qq^\prime} is element of CKM matrix

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Leptonic branching ratio

$$R_{M} = \frac{\Gamma(M \longrightarrow e\nu_{e})}{\Gamma(M \longrightarrow \mu\nu_{\mu})}$$

$$R_{M}^{SM} = \left(\frac{m_{e}}{m_{\mu}}\right)^{2} \left(\frac{m_{M}^{2} - m_{e}^{2}}{m_{M}^{2} - m_{\mu}^{2}}\right)^{2} (1 + \delta R_{QED}),$$

$$\delta R_{QED} \text{ is electromagnetic correction}$$

$$m_{e} = 0.5109989461(13)MeV$$

$$m_{\mu} = 105.6583745(24)$$

$$(6)$$

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Leptonic branching ratio of pion decay

$$egin{aligned} R_\pi &= rac{\Gamma(\pi \longrightarrow e
u_e)}{\Gamma(\pi \longrightarrow \mu
u_\mu)} = (1.2352 \pm 0.0002) imes 10^{-4} \ m_\pi &= 139.57018 \pm 0.00035 MeV \ \pi^+ ext{ is made of } u ar{d} ext{ and } \pi^- ext{ is made of } ar{u} d \end{aligned}$$

leptonic branching ratio of kaon decay

$$R_{\mathcal{K}} = \frac{\Gamma(\mathcal{K} \longrightarrow e\nu_{e})}{\Gamma(\mathcal{K} \longrightarrow \mu\nu_{\mu})} = (2.477 \pm 0.0001) \times 10^{-5}$$
(9)

$$m_K = 493.667 \pm 0.013 MeV$$

 K^+ is made of $u\bar{s}$ and K^- is made of $\bar{u}s$

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Pion and Kaon rare decay experiments

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https://na62.web.cern.ch/na62/Home/NA62Detector.html

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- In the North Area of the SPS accelerator at CERN
- Study rare decays of charged kaons
- In two body decay spectrum is monochromatic $P = p_1 + p_2$ $p_1^2 = (P - p_2)^2 = M^2 + m_2^2 - 2E_2M = m_1^2$ $E_2 = \frac{M^2 + m_2^2 - m_1^2}{2M}$ $E_1 = M - E_2 = \frac{M^2 + m_1^2 - m_2^2}{2M}$
- Kinematic identification: $M_{miss}^2(I) = (P_K P_I)^2$

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- NA62 criteria is $-M_1^2 < M_{miss}^2(I) < M_2^2$ $M_1^2 \sim 0.013$ to 0.016 $M_2^2 \sim 0.010$ to 0.013
- \sim 150000 reconstructed ${\it K}^{\pm} \longrightarrow e^{\pm} \nu_{e}$ events
- with 11% background
- kaons decay in flight with 74 \pm 1.4 GeV energy
- reconstructed lepton momentum is in the range of 13 to 65 GeV
- $R_{K} = (2.488 \pm 0.010) \times 10^{-5}$



arXiv:1212.4012v2

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PIENU



arXiv:1506.05845v2

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- $\pi^+
 ightarrow e^+
 u$, $E_{e^+}=69.8~{
 m MeV}$
- $\pi^+
 ightarrow \mu^+
 u$, $\mu^+
 ightarrow e^+
 u$, $E_{e^+} = 0.5 52.8~{
 m MeV}$
- pion at rest
- Energy cut equal to 53 MeV
- $R_{\pi} = [1.2344 \pm 0.0023(stat) \pm 0.0019(Syst)] \times 10^{-4}$
- In the near future statistical uncertainty reduces by a factor of 3

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arXiv:1506.05845v2

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E949

- Brookhaven National Laboratory, USA
- 1.70×10^{12} stopped kaon
- $Br(K^+ \to \mu^+ \nu \nu \nu) < 2.4 \times 10^{-6}$ at 90 % C.L. , arXiv:1606.09054 [hep-ex]



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Other constraints

- Old TRIUMF, (Phys. Rev. D **37** (1988) 1131), $R = \frac{\Gamma(\pi \to e\nu Z')}{\Gamma(\pi \to \mu\nu)} < 4 \times 10^{-6}$
- Nucl. Phys. B **149** (1979) 365, $Br(K^+ \to e^+ \nu \nu \nu) < 6 \times 10^{-5}$ at 90 % C.L.
- KLOE, TREK, 1970s measurement of R_K

Short-baseline neutrino experiments

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DUNE near detector



https://sciencesprings.files.wordpress.com/2019/10/fnal-dune-neardetector.png

DUNE near detector applications

- Neutrino flux measurement
- Neutrino cross-section measurements
- Determination of parton and gluon structure function
- Test of sum rules
- Nuclear structure of Ar, Pb, Fe, and C targets (Form factors, etc)
- QCD measurements
- Weak angle measurement
- Beyond SM measurements

FASER and FASER ν



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$FASER\nu$

- Detect collider neutrinos for the first time
- Detector is located 480 m from ATLAS IP
- FASER ν is an emulsion detector
- Detect tau neutrinos
- Neutrino cross-section measurements and nucleon structure
- Charm meson production and decay and study Quantum Chromodynamics
- Reduction of atmospheric neutrino flux uncertainaties
- Test of physics beyond the standard model
- FASER ν 2 and FLARE are future detectors at high luminosity LHC

FASER ν 2 and FLARE



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SHIP experiment (Beam Dump Facility)



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$$\mathcal{M} = g_{\alpha i} G_F f_M V_{qq'} k_\mu \varepsilon_\alpha(p) \bar{u}(l) \gamma^\mu P_L \frac{Q}{Q^2} \gamma^\alpha P_L v(q), \qquad (10)$$

$$\sum_{spins} |\mathcal{M}|^2 = g_{\alpha i}^2 G_F^2 f_M^2 V_{qq'}^2 \left(\sum_{polarization} \varepsilon_{\alpha}^*(p) \varepsilon_{\beta}(p) \right)$$

$$tr \left(v(q) \bar{v}(q) P_R \gamma^{\alpha} \frac{Q^2 + m_I Q}{Q^2} u(I) \bar{u}(I) \frac{Q^2 + m_I Q}{Q^2} \gamma_{\beta} P_L \right)$$
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$$\sum_{spins} |\mathcal{M}|^{2} = \left(\sum_{i} g_{\alpha i}^{2}\right) G_{F}^{2} f_{M}^{2} V_{qq'}^{2} \left(g_{\alpha\beta} - \frac{p_{\alpha}p_{\beta}}{m_{Z}^{2}}\right)$$

$$tr \left(\oint P_{R} \gamma^{\alpha} \frac{Q^{2} + m_{I} Q}{Q^{2}} (I + m_{I}) \frac{Q^{2} + m_{I} Q}{Q^{2}} \gamma_{\beta} P_{L} \right)$$
(12)

neglecting the lepton mass in the case of electron

$$\sum_{spins} |\mathcal{M}|^{2} = \left(\sum_{i} g_{\alpha i}^{2}\right) G_{F}^{2} f_{M}^{2} V_{qq'}^{2} \left(m_{M}^{2} + m_{Z'}^{2} - 2m_{M} E_{Z'} + \frac{(m_{M}^{2} - m_{Z'}^{2} - 2m_{M} E_{I})(m_{M}^{2} - m_{Z'}^{2} - 2m_{M} E_{\nu})}{m_{Z'}^{2}}\right)$$
(13)

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$$\Gamma(M \longrightarrow l_{\alpha}\nu Z') = \frac{1}{64\pi^{3}m_{M}} \int_{E_{l}^{min}}^{E_{l}^{max}} \int_{E_{\nu}^{min}}^{E_{\nu}^{max}} dE_{l}dE_{\nu} \sum_{spins} |\mathcal{M}|^{2}.$$
(14)

$$E_{l}^{min} = m_{l}, \qquad E_{l}^{max} = \frac{m_{M}^{2} - m_{Z'}^{2}}{2m_{M}},$$

$$E_{\nu}^{min} = \frac{m_{M}^{2} - m_{Z'}^{2} - 2m_{M}E_{l}}{2m_{M}}, \qquad E_{\nu}^{max} = \frac{m_{M}^{2} - m_{Z'}^{2} - 2m_{K}E_{l}}{2(m_{M} - 2E_{l})}$$

$$(M \longrightarrow e\nu_{\alpha}Z') = \frac{g_{e\alpha}^{2}G_{F}^{2}\cos^{2}(\theta_{C})f_{M}^{2}}{6144\pi^{3}m_{M}^{3}m_{Z'}^{2}} \left(m_{M}^{8} + 72m_{M}^{4}m_{Z'}^{4} - 64m_{M}^{2}m_{Z'}^{6} + 24\left(3m_{M}^{4}m_{Z'}^{4} + 4m_{M}^{2}m_{Z'}^{6}\right)\log\left(\frac{m_{Z'}}{m_{M}}\right) - 9m_{Z'}^{8}\right).$$
(15)

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- Considering the lepton mass in the case of meson decay to the muon caused 5% correction
- In this case the decay rate is calculated numerically
- In the case of meson decay to electron, the effect of electron mass is negligible

- $\Gamma_{SM}(M \longrightarrow l_{\alpha}\nu)$, $\Gamma(M \longrightarrow l_{\alpha}\nu Z')$, experimental R_M
- In the case of experimental that R_M is measured, we defined $\chi^2_K = \frac{(R_K^{EXP} R_K^{SM} R_K^{NEW})^2}{\sigma^2}$
- In the case of NA62, because of kinematic criteria, $E_e > 233.68 MeV$
- We presented analitical formula for the spectrum of the Z' and produced neutrinos from meson and Z' decay, arXiv:1702.04187 [hep-ph].

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Constraining the coupling from decay to electron



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Constraining the coupling from decay to electron



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Constraining the coupling from decay to muon



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DUNE ND Spectrum



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Neutrino flux at FASER ν



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Constraining the coupling



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Constraining the coupling



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Constraining the coupling



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Constraining the coupling from $\mathsf{FASER}\nu$



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Summary

- Meson decay and short-baseline neutrino experiments constrain neutrinophilic interaction
- NA62 and E949 sets the most stringent bound on the coupling for 4 $MeV < m_{Z^\prime} < 50~MeV$
- DUNE ND will make comparable or more stringent constraints and has flavor sensitivity
- FASER ν sets more stringent constraints for 50 $MeV < m_{Z'} < 150 MeV$
- FASER ν 2, FLARE, and SHIP sets one order of magnitude more stringent constraints in 4 $MeV < m_{Z'} < 500 MeV$

Thank you for your attention.

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