

Primordial Black Holes from Inflation

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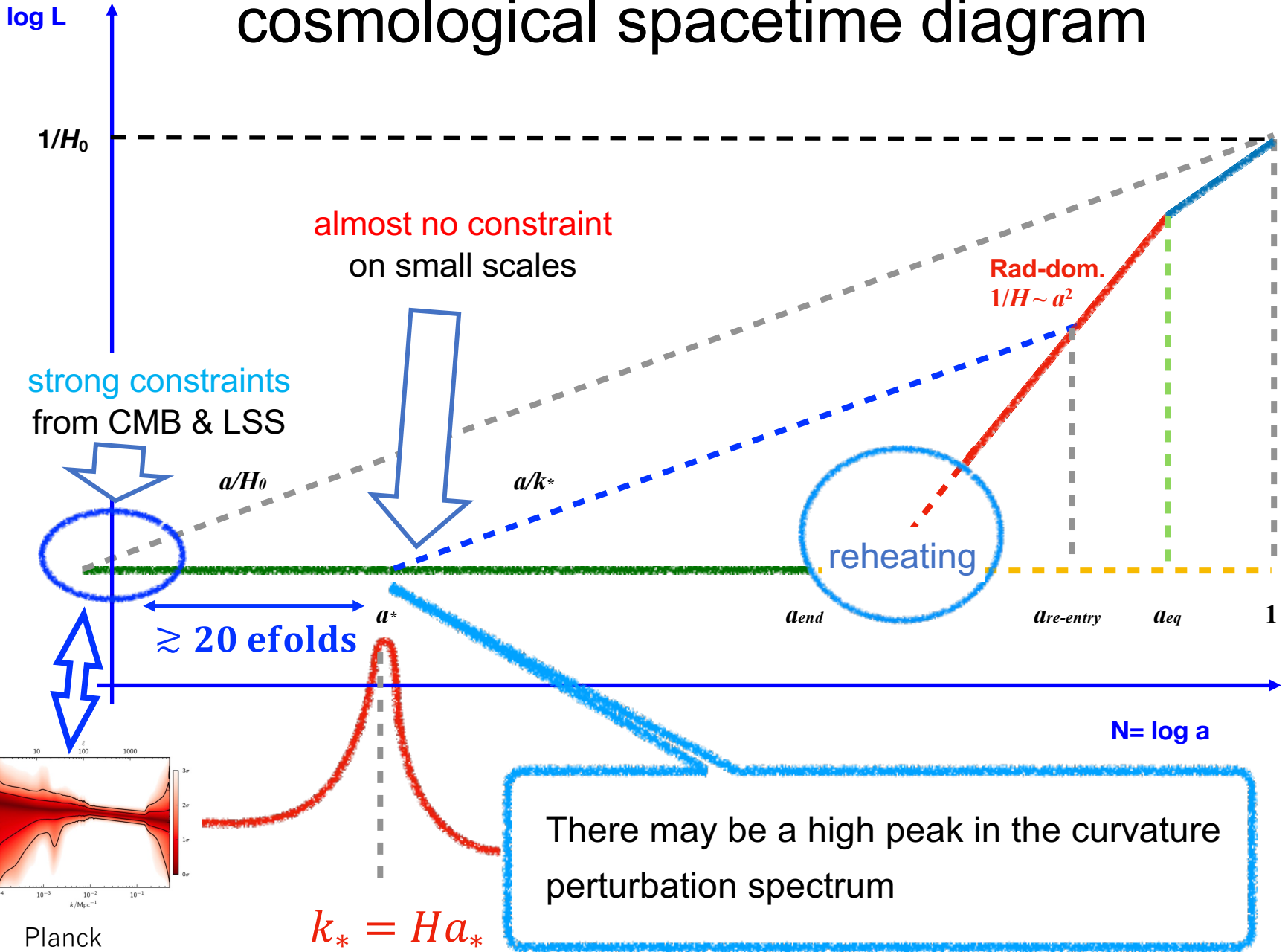
LeCosPA, Taiwan National University

PBH formation

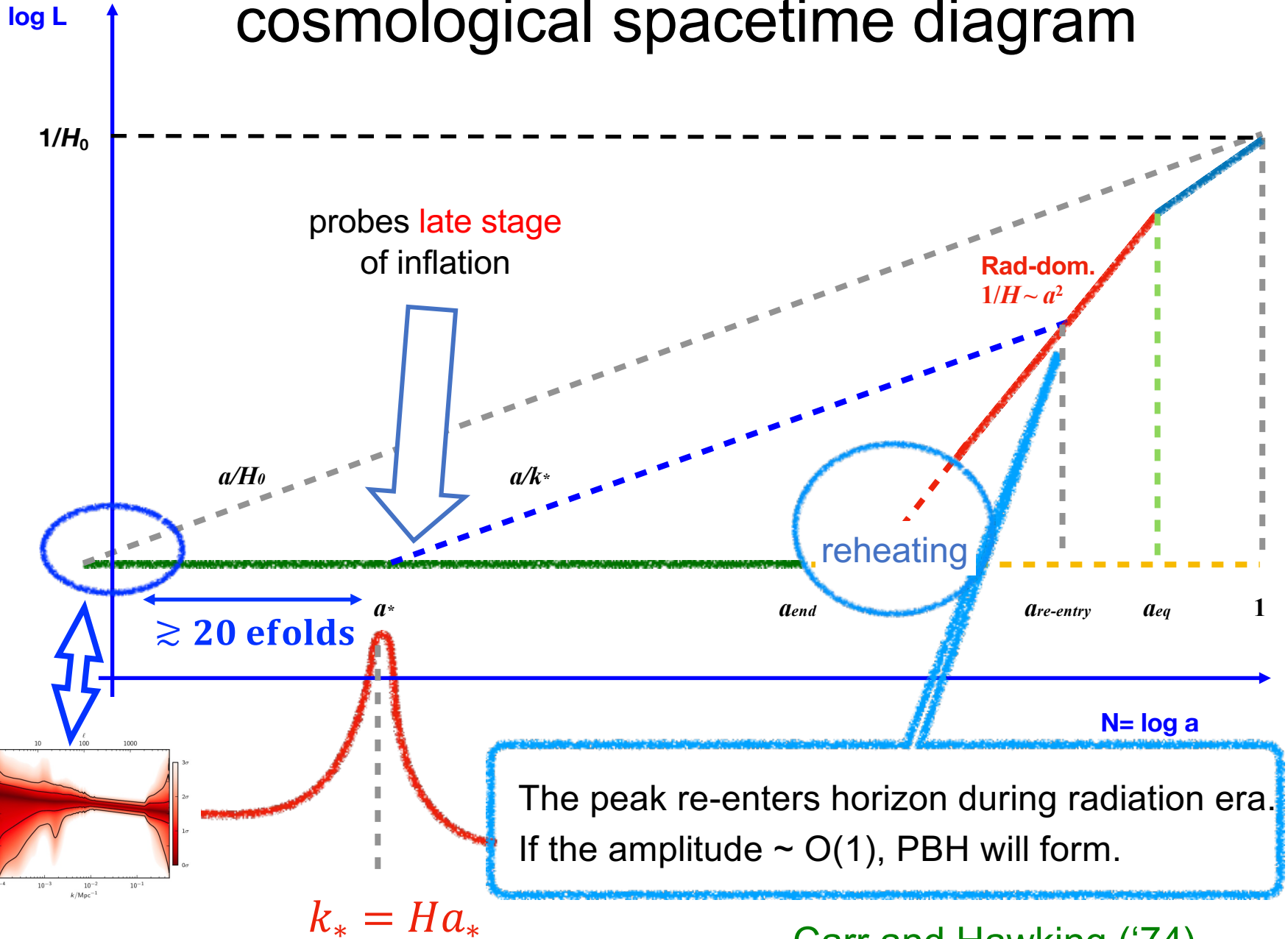
- conventional scenario -

originally proposed by Hawking ('71)

cosmological spacetime diagram



cosmological spacetime diagram

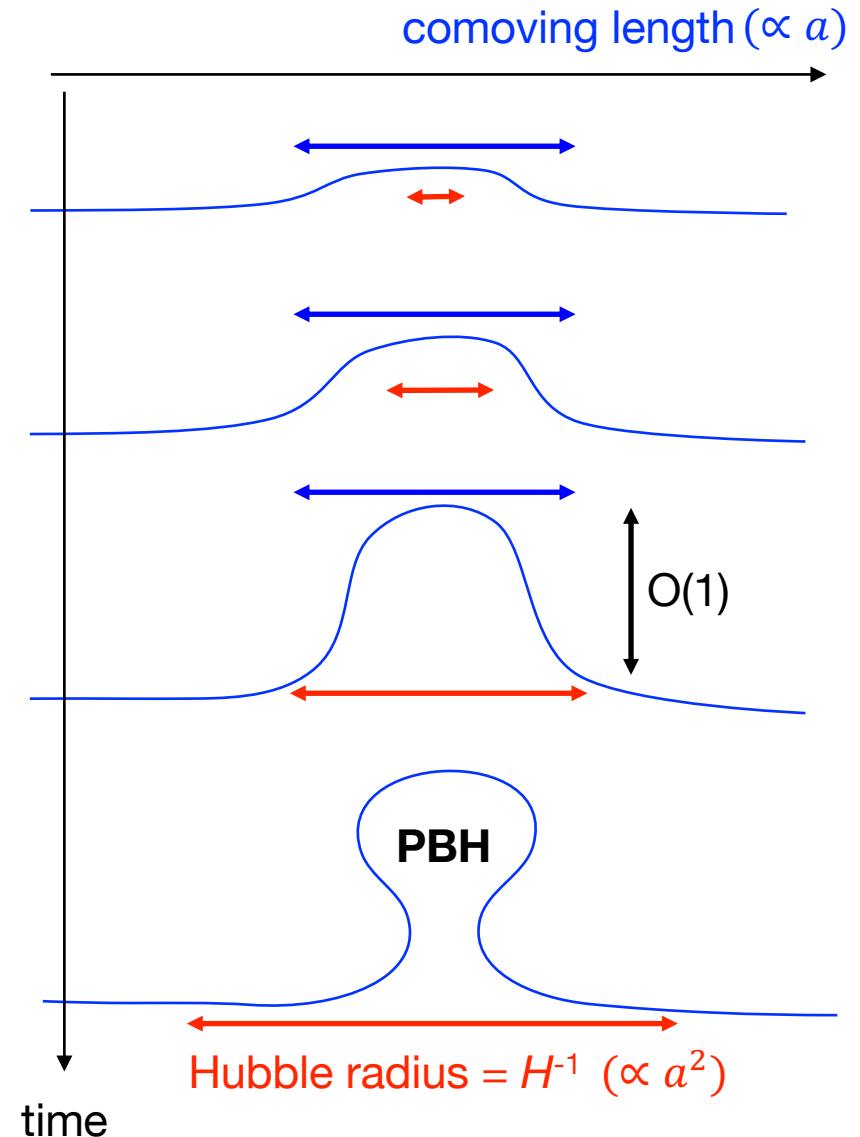


Carr and Hawking ('74)

Conventional PBH formation in a nutshell

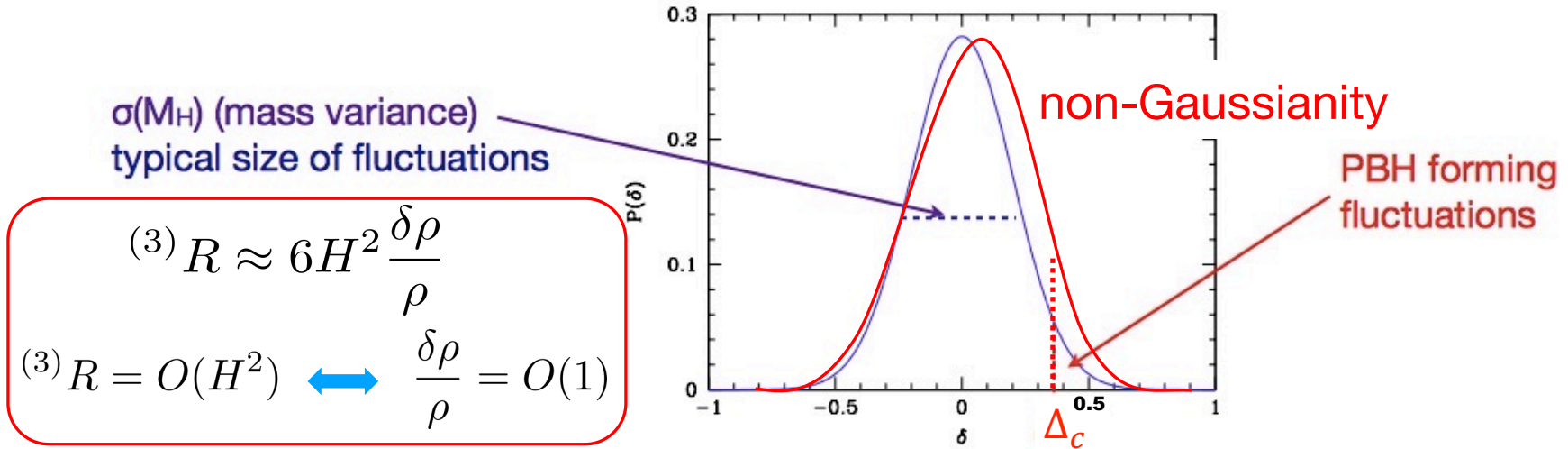
- Primordial Black Holes (PBHs) are those formed in the very early universe, conventionally when the universe was radiation-dominated.
- Presumably they originate from a large positive curvature perturbation produced during inflation (which hence should be a rare event).
- For a BH to form during radiation dominance, the perturbation must be $O(1)$ on the Hubble horizon scale.

$$M_{\text{PBH}} \sim M_{\text{horizon}} \sim \left(\frac{100 \text{ MeV}}{T} \right)^2 M_{\odot} \sim \left(\frac{\ell}{1 \text{ pc}} \right)^2 M_{\odot}$$



β : fraction of ρ that turns into PBHs

for **Gaussian** probability distribution



$${}^{(3)}R \approx 6H^2 \frac{\delta\rho}{\rho}$$

$${}^{(3)}R = O(H^2) \iff \frac{\delta\rho}{\rho} = O(1)$$

- When $\sigma_M \ll \Delta_c$, β can be approximated by exponential:

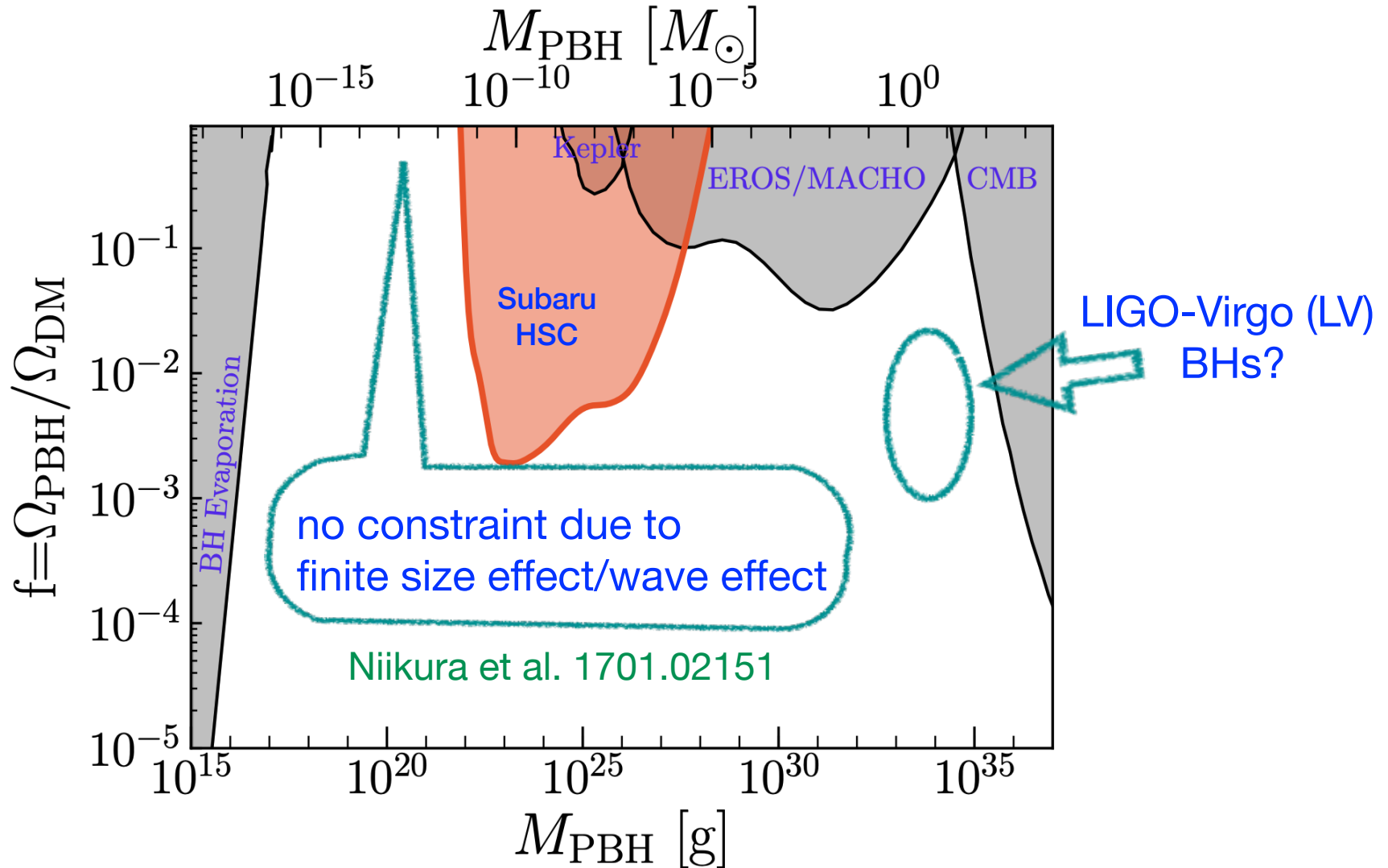
$$\beta \approx \sqrt{2/\pi} \frac{\sigma(M)}{\Delta_c} \exp\left(-\frac{\Delta_c^2}{2\sigma(M)^2}\right) \quad \Delta_c \equiv \left(\frac{\delta\rho_c}{\rho}\right)_{\text{crit}} \sim 0.4$$

Carr, ApJ 201, 1 (1975), ...

- Recent studies indicates **enhanced production**: $\Delta_c \sim 0.2$
using **peak theory** Yoo, Harada, Garriga & Kohri, 1805.03946
- Criterion using **compaction function** ($C \sim GM/R$) may be more relevant
Musco, De Luca, Franciolini & Riotto, 2011.03014

- Non-Gaussianity** may significantly affect β

observational constraints



big window at $M_{\text{PBH}} \approx 10^{17} - 10^{22} \text{g}$ \leftrightarrow $T_{\text{re-entry}} \sim 10^4 - 10^8 \text{ GeV}$

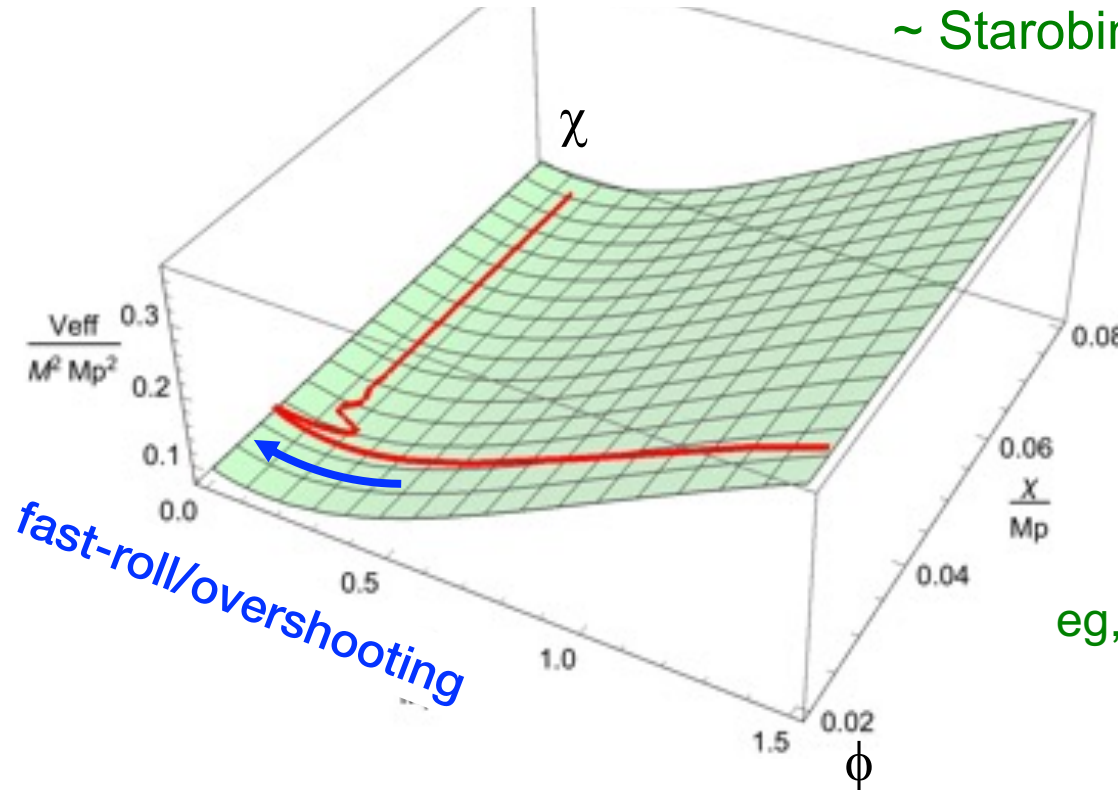
Inflation models

Model 1a: two-field inflation

Pi, Zhang, Huang & MS, 1712.09896

$$S_J = \int d^4x \sqrt{-g} \left\{ \frac{M_{\text{Pl}}^2}{2} \left(R + \frac{R^2}{6M^2} \right) - \frac{1}{2} g^{\mu\nu} \partial_\mu \chi \partial_\nu \chi - V(\chi) - \frac{1}{2} \xi R \chi^2 \right\}.$$

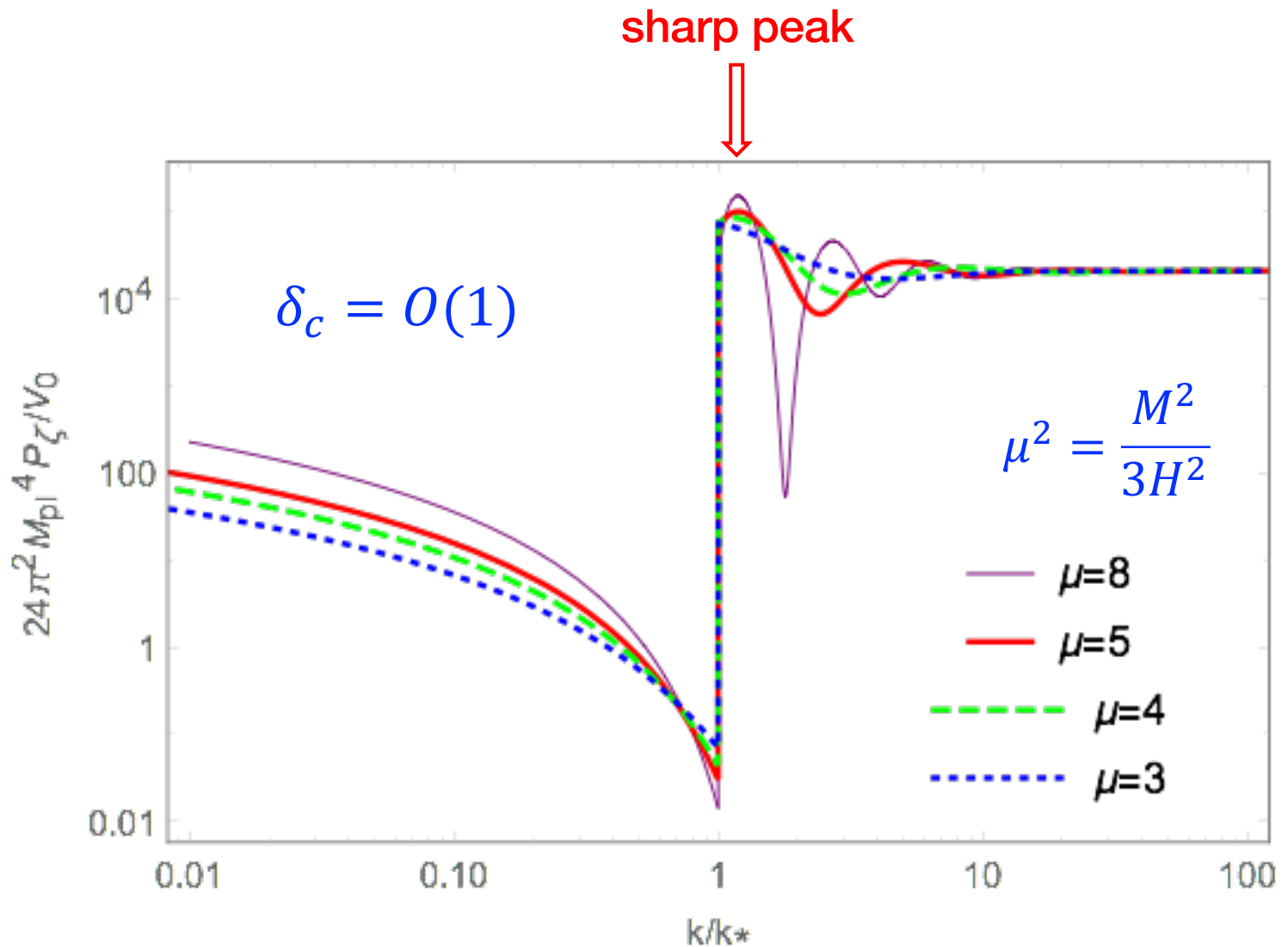
~ Starobinsky (scalaron) + curvaton



many 2-stage models
can account for
PBH formation

eg, Kawasaki et al., 1606.07631

- Scalaron ϕ becomes massive at the end of the 1st stage.
- Field χ plays the role of inflaton at the 2nd stage.



- 2-stage model can produce a sharp peak from the transition

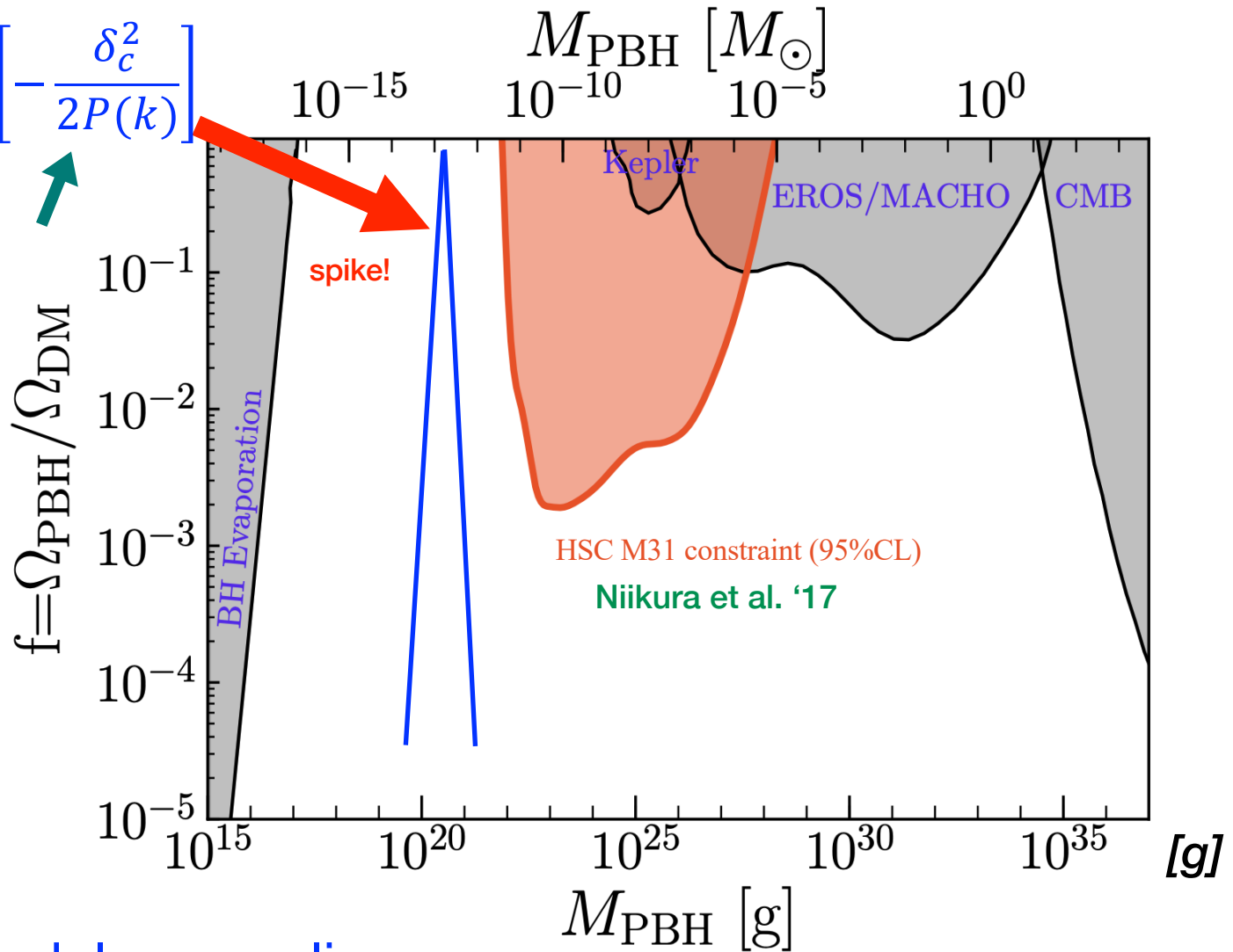
- non-Gaussianity is small in most 2-stage models.

PBH mass function

$$\delta_c = O(1)$$

$$f(M_{\text{PBH}}) \propto \exp\left[-\frac{\delta_c^2}{2P(k)}\right]$$

sharp peak



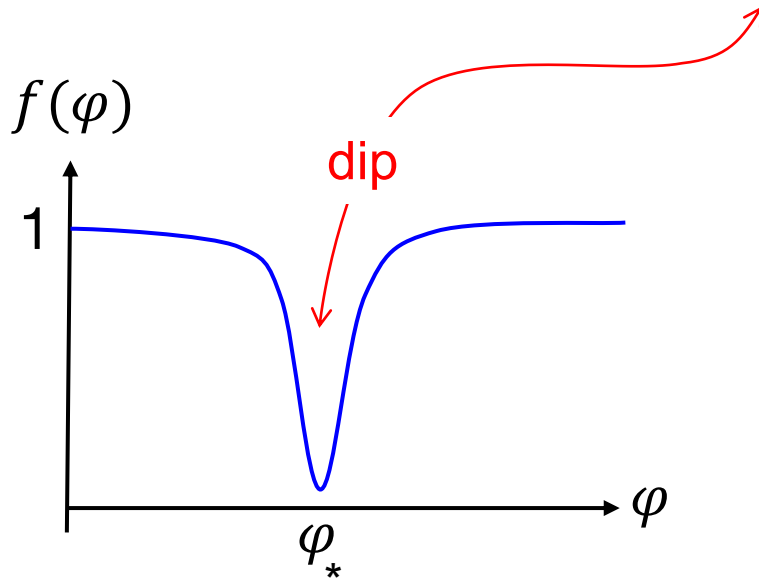
2-stage model can realize

PBH=CDM scenario with a monochromatic PBH mass

Model 1b: non-minimal curvaton

Pi & MS, 2112.12680

$$\mathcal{L} = -\frac{1}{2}(\partial\varphi)^2 - V(\varphi) - \frac{1}{2}\underline{f(\varphi)^2}(\partial\chi)^2 - \frac{1}{2}m_\chi^2\chi^2 \quad (m_\chi^2 \ll H^2)$$



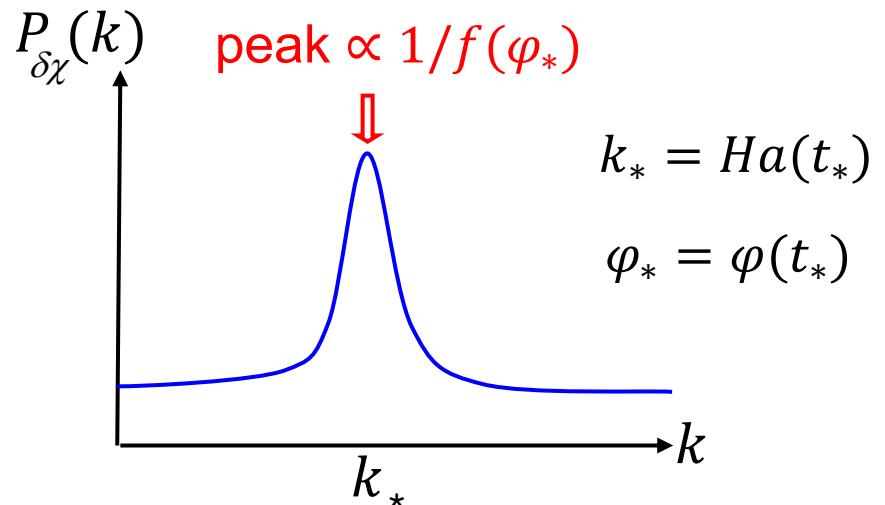
- assume $f \ll 1$ when $\varphi \sim \varphi_*$

- vacuum fluctuation: $f(\varphi)\delta x = \frac{H}{2\pi}$

$$\rightarrow \delta x = \frac{H}{2\pi f}$$

- $\delta\chi$ is **enhanced** at $\varphi = \varphi_*$

\rightarrow leads to PBH formation



Fully non-Gaussian curvature perturbation

$$e^{4\zeta} - \left[\frac{4r}{3+r} \left(1 + \frac{\delta\chi}{\chi} \right)^2 \right] e^\zeta + \left[\frac{3r-3}{3+r} \right] = 0$$

MS, Valiviita & Wands, astro-ph/0607627

ζ = curvature perturbation on **uniform density** slices

$r = \rho_\chi / \rho_{\text{tot}}$ at epoch of curvaton decay

- **PBH formation** is fully **non-perturbatively non-Gaussian**:
criterion $\zeta > \zeta_{\text{cr}} \sim 0.5$ gives a highly nonlinear expression in $\delta \equiv \delta\chi/\chi$



$\Delta(R) : \delta\rho/\rho$
smoothed over
comoving scale R

$$\Delta(R) \approx \frac{4k_*^2 R^2}{9} \zeta(R)$$

for spectrum peaked at $k=k_*$

- **Power spectrum** is well approximated by that of $\delta\chi^2$ for $r \ll 1$.

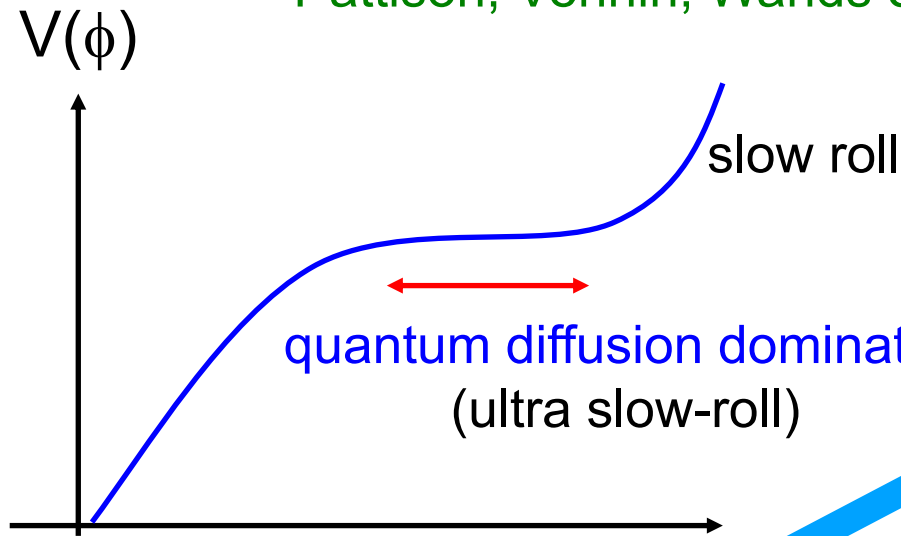
Pi & MS, 2112.12680

$$P_{\delta\chi}(k_*) \longrightarrow \zeta = \frac{\delta\chi}{\chi} + \frac{3}{4r} \left(\frac{\delta\chi}{\chi} \right)^2 \text{ for } r \ll 1 \longrightarrow P_\zeta(k_*)$$

perturbative approximation is valid

Model 2a: potential with inflection point

Pattison, Vennin, Wands & Assadullahi, 2101.05741, ...



$$\delta\phi \sim \frac{\dot{\phi}}{H}$$

vacuum
fluctuations $\frac{H}{2\pi}$

classical
background
motion

$$\mathcal{R} = \delta N = -\frac{H}{\dot{\phi}} \delta\phi \sim \mathcal{O}(1)$$

tail of the PDF can be exponentially enhanced: $e^{-c\mathcal{R}}$ instead of $\exp(-c\mathcal{R}^2)$

→ fully non-Gaussian PBH formation

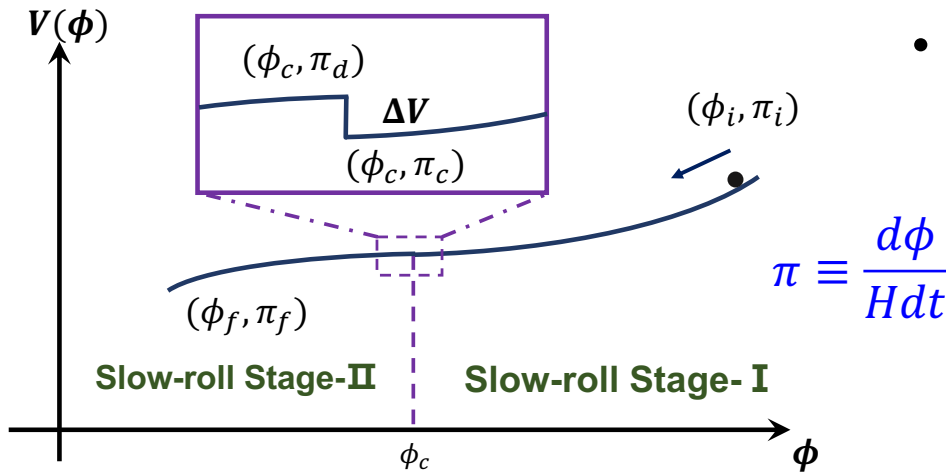
exponential tail is actually quite common in featured potential

Pi & MS, 2211.13932

Model 2b: Amplification by Upward Step

- One Small Step for an Inflaton, One Giant Leap for Inflation -

Cai, Ma, MS, Wang & Zhou, 2112.13836



- energy conservation at the step:

$$\pi_d = -\sqrt{\pi_c^2 - 6\Delta V/V}$$

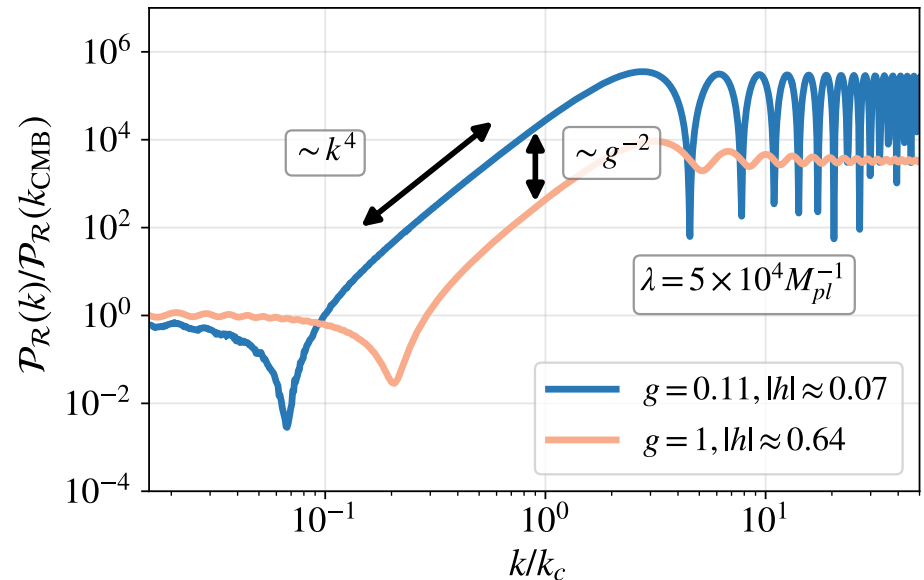
(in $M_{\text{Planck}} = 1$ units)

$$\pi_c = -\sqrt{2\varepsilon_c}$$

ε_c : SR parameter at $\phi = \phi_c$

even for a tiny step, $\Delta V \ll V$,
 $P_R(k)$ is enhanced by $1/g^2$ if

$$g \equiv \frac{\pi_d}{\pi_c} \ll 1$$



non-perturbative non-Gaussianity at tail of distribution

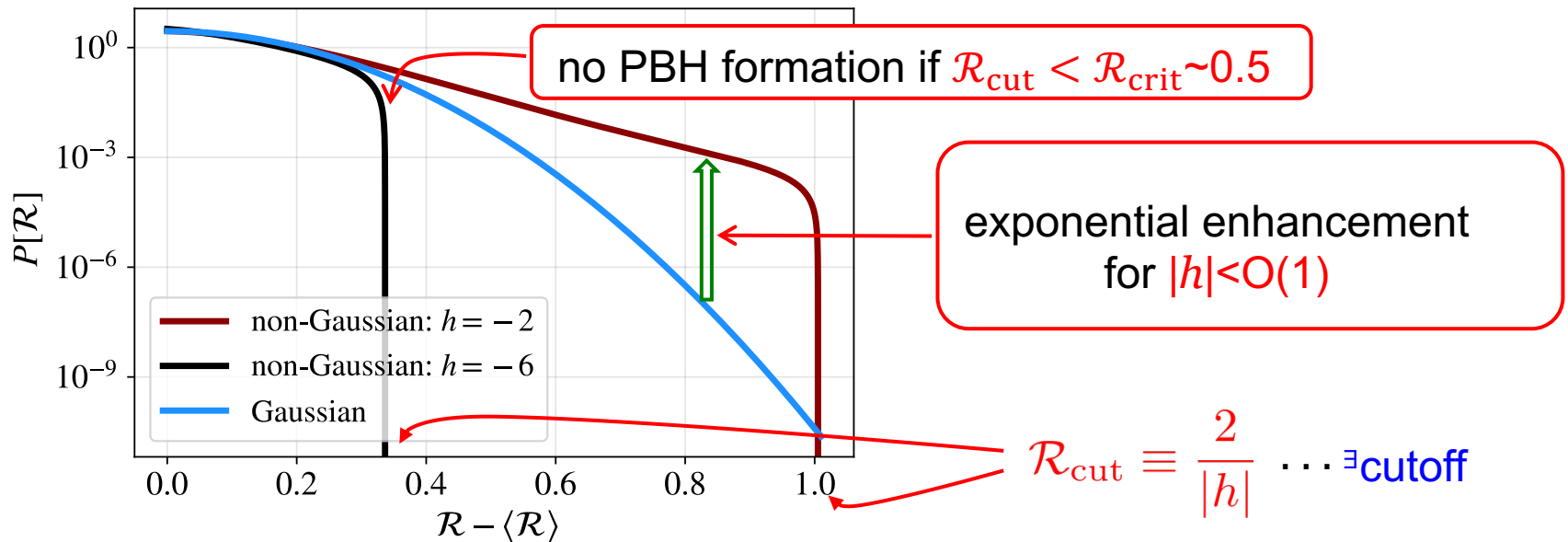
- perturbative non-Gaussianity is small if $-h \equiv \frac{6\sqrt{2\varepsilon_V}}{|\pi_d|} \ll 1$

$$\mathcal{R} = \mathcal{R}_G + \frac{|h|}{4}\mathcal{R}_G^2 + \frac{|h|^2}{8}\mathcal{R}_G^3 + \dots \quad \Longrightarrow \quad \mathcal{P}(k) \approx \mathcal{P}_G(k)$$

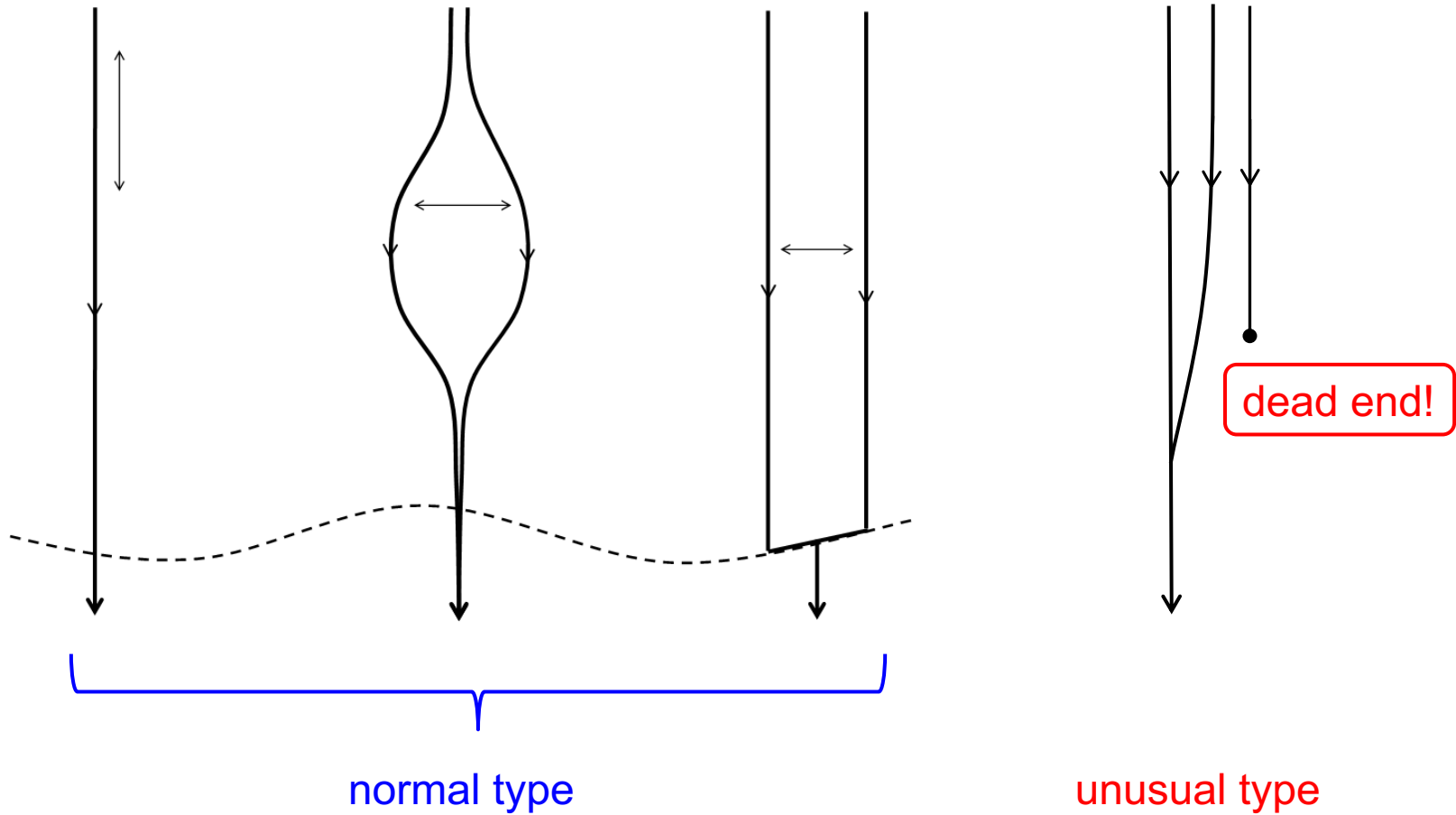
power spectrum is given by Gaussian part

- tail of distribution is extremely non-Gaussian

$$\left| \frac{d\mathcal{R}_G}{d\mathcal{R}} \right| \quad P[\mathcal{R}] = \frac{2 - |h|\mathcal{R}}{\Omega} \exp \left[-\frac{\mathcal{R}^2(4 - |h|\mathcal{R})^2}{32\sigma_{\mathcal{R}}^2} \right]$$



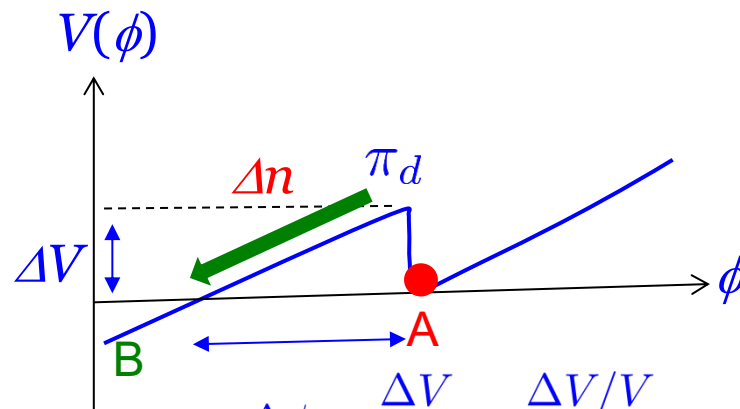
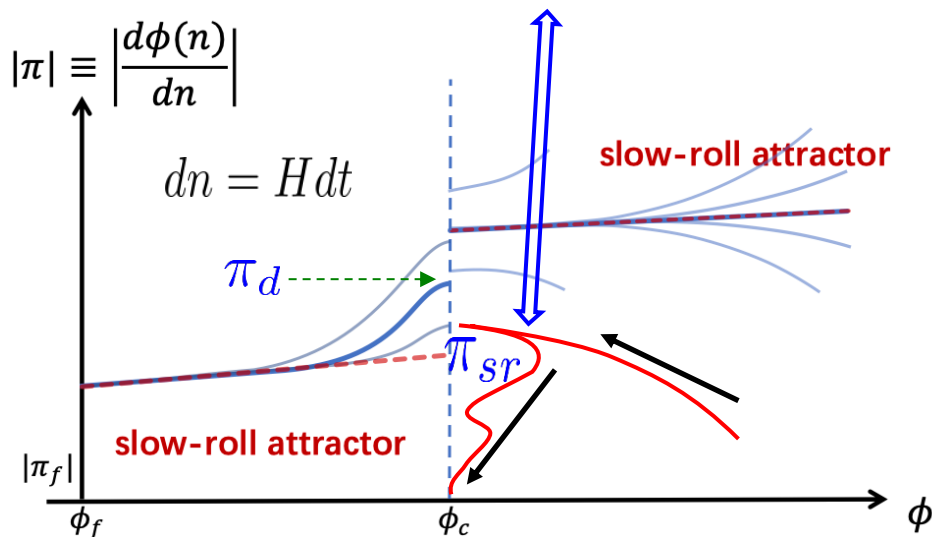
unusual type of field space trajectories



PBH formation during inflation

$$\mathcal{R} \simeq \frac{2}{|h|} \left(1 - \sqrt{1 - |h|\mathcal{R}_G} \right) \quad \text{PDF cutoff at } \mathcal{R} = \mathcal{R}_{\text{cut}} \equiv \frac{2}{|h|}$$

⇒ trajectories that can't climb the step



$$\Delta\phi = \frac{\Delta V}{V'} = \frac{\Delta V/V}{\sqrt{2\epsilon}}$$

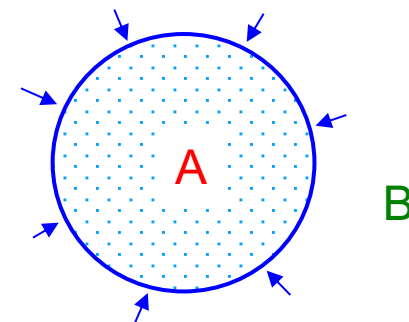
$$\Delta n \simeq \frac{\Delta\phi}{|\pi_{sr}|} = \frac{\Delta V/V}{\pi_{sr}^2}$$

of e-folds region A expands

region stuck at $\phi = \phi_c$ will become PBH!

region **A** expands until $V(\phi)$ surrounding it becomes smaller than $V(\phi_A) = V_0$

$$M_{\text{BH}} \simeq H^{-1} e^{3\Delta n}$$



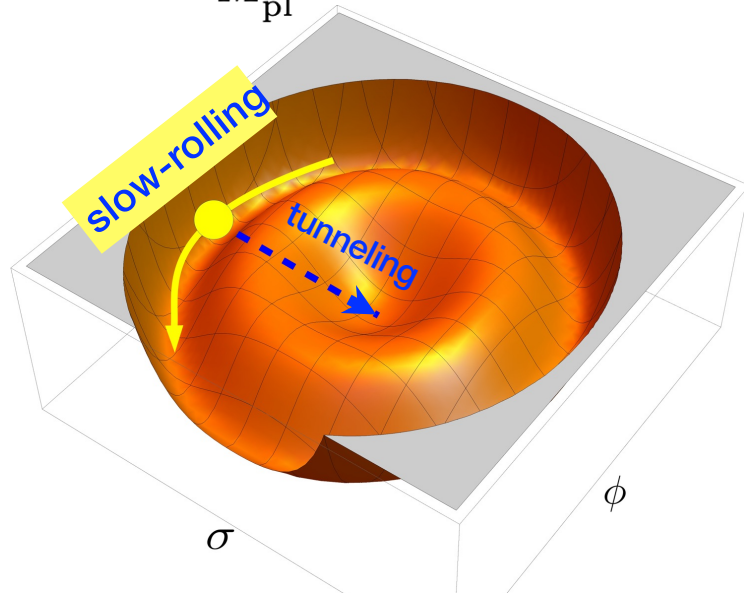
Model 4: PBH-as-MVP scenario

PBH formation during inflation due to vacuum tunneling
(not from curvature perturbation)

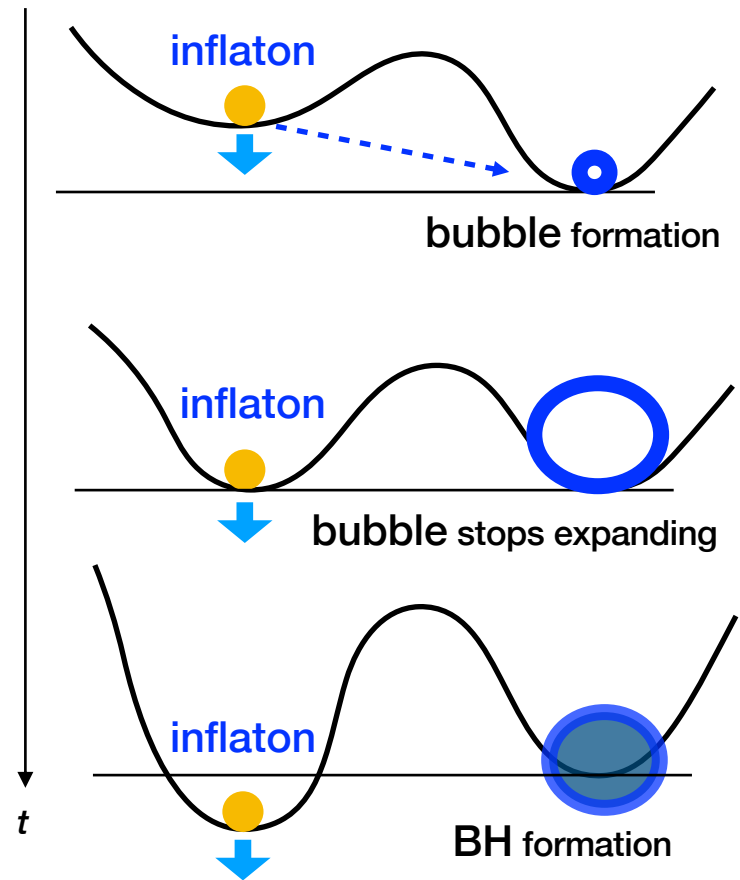
Garriga, Vilenkin & Zhang, 1512.01819, Deng & Vilenkin, 1710.02865,...

example:

$$V(\phi, \sigma) = m^2(\phi^2 + \sigma^2) - a(\phi^2 + \sigma^2)^2 + \frac{c}{M_{\text{pl}}^2}(\phi^2 + \sigma^2)^3 + gM_{\text{pl}}^4 \sin\left(\frac{\phi}{fM_{\text{pl}}}\right)$$

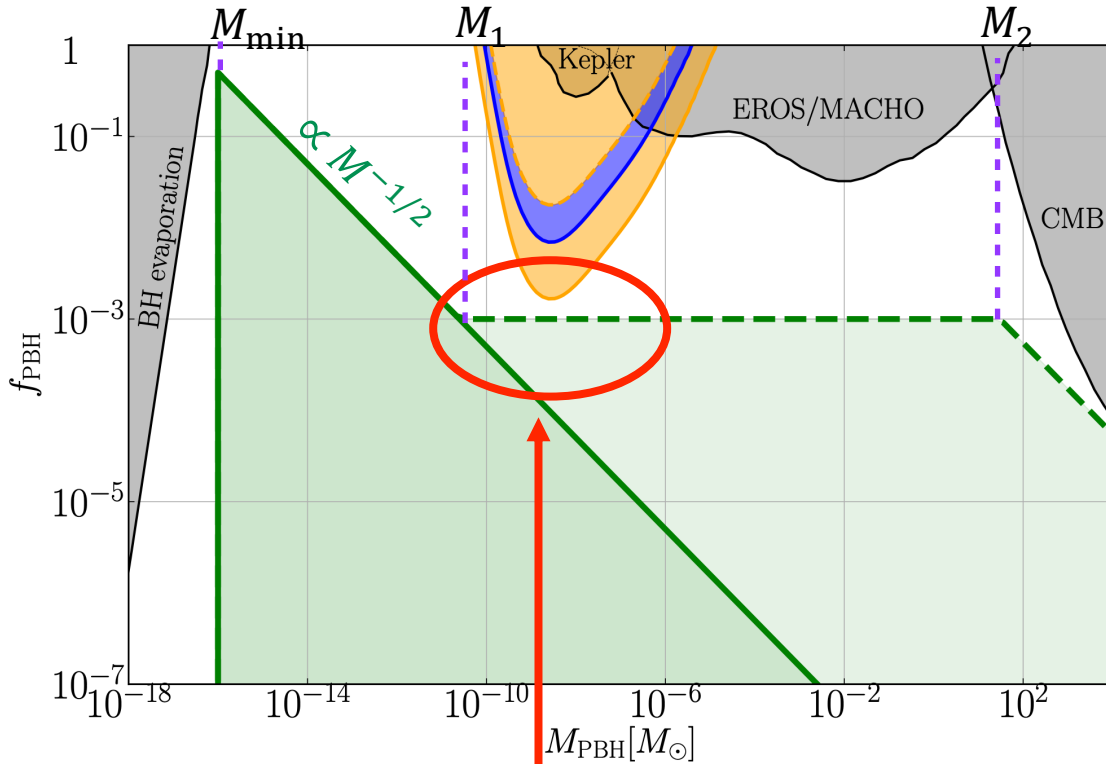


can probe multiverse!



Mass function

Kusenko, MS, Sugiyama, Takada, Takhistov & Vitagliano, 2001.09160



may be tested by Subaru HSC!

analysis is in progress

Implications to GW cosmology?

- for scale M re-entering horizon during radiation-dom stage

$$f(M) = \lambda \left(\frac{M}{M_{\text{min}}} \right)^{-1/2} : M_{\text{min}} < M$$

$M \simeq M_{\text{min}} \cdots$ CDM

- if there is an intermediate matter-dom stage

$$f(M) = \lambda \left(\frac{M_1}{M_{\text{min}}} \right)^{-1/2} : M_1 < M < M_2$$

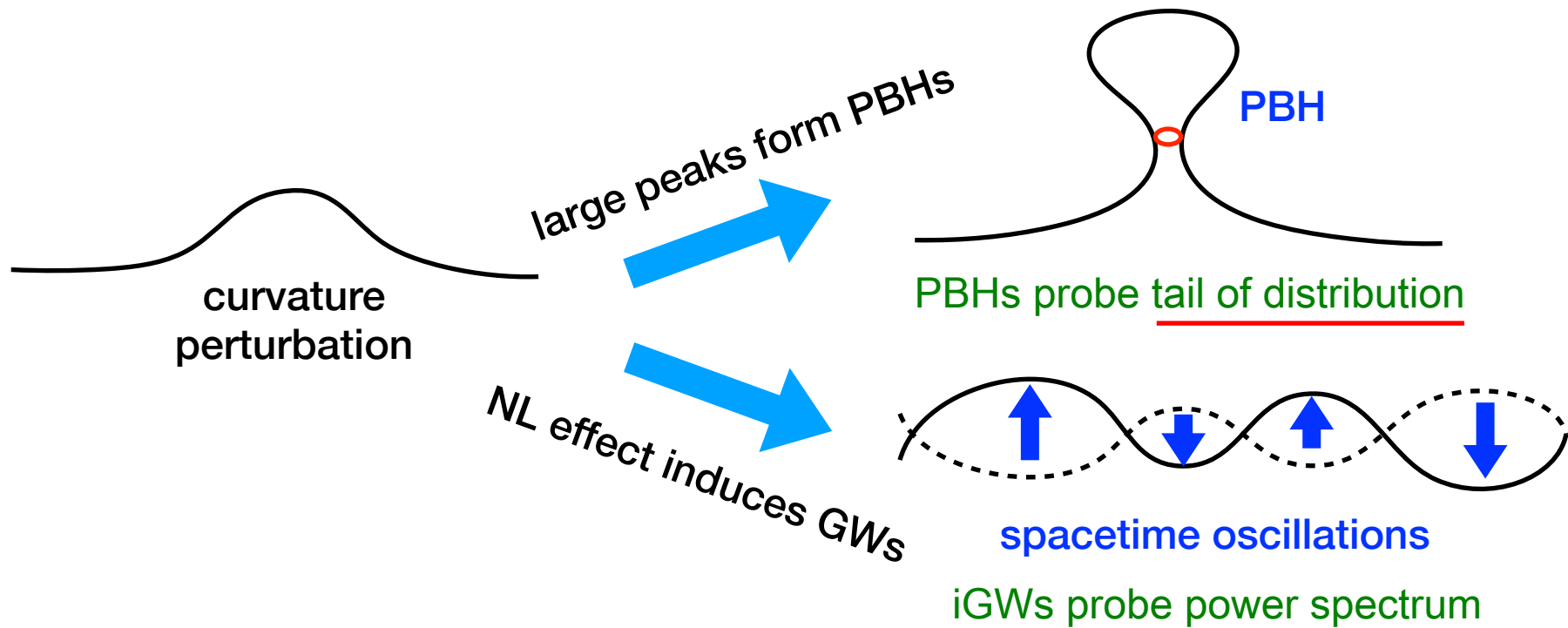
$M \simeq M_2 \cdots$ LIGO BHs

$$f(M) = \lambda \left(\frac{M_2}{M_1} \right)^{1/2} \left(\frac{M}{M_{\text{min}}} \right)^{-1/2} : M_2 < M$$

$M \gg M_2 \cdots$ SMBHs

induced GWs

GWs can capture PBHs!



PBHs = CDM with $M_{\text{PBH}} \sim 10^{21} \text{g}$
generates GWs with $f \sim 10^{-3} \text{ Hz}$



Background GWs
at LISA band

PBHs = LV BHs with $M_{\text{PBH}} \sim 10 M_{\odot}$
generates GWs with $f \sim 10^{-8} \text{ Hz}$

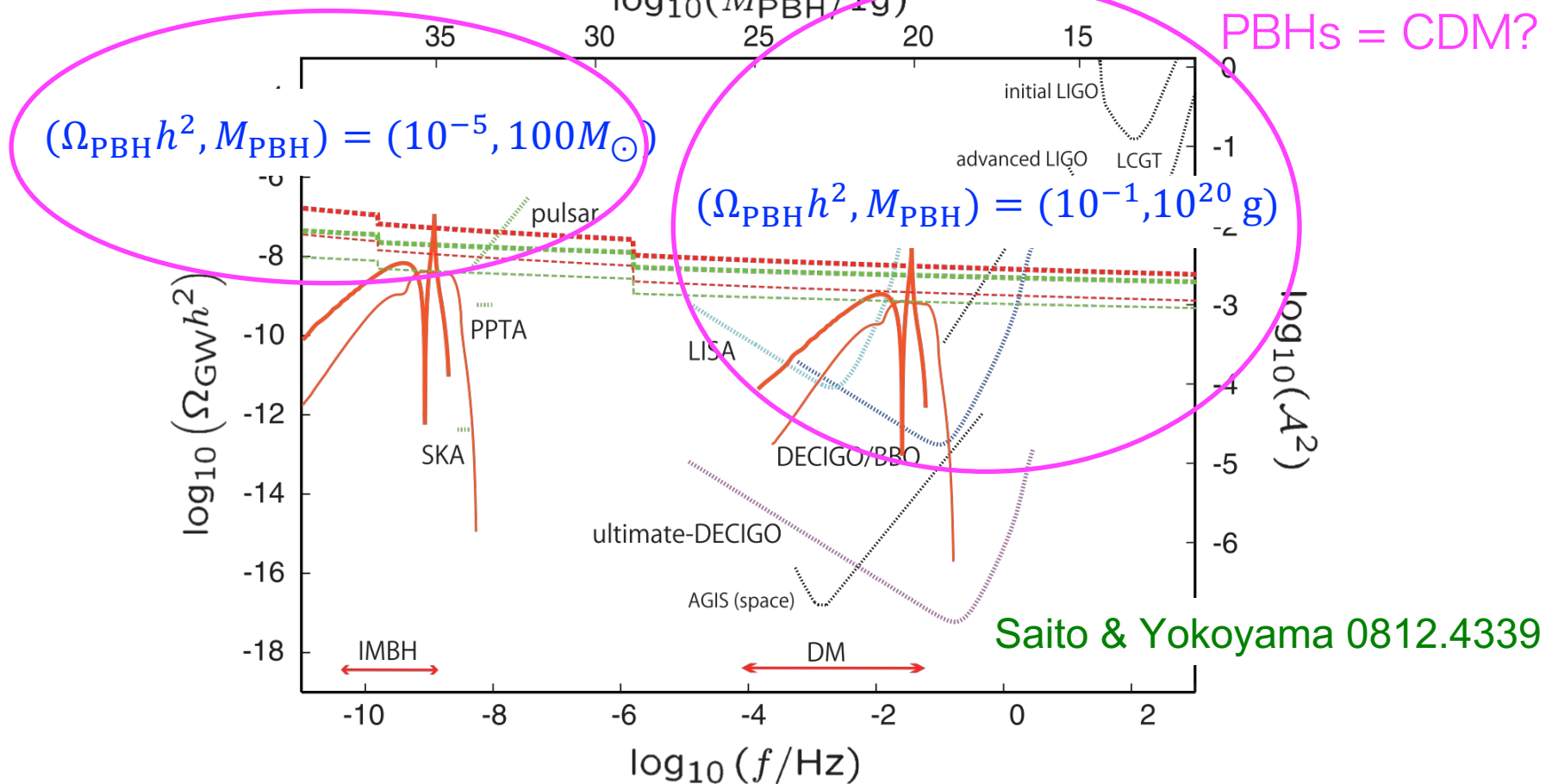


Background GWs
at PTA band

GWs can test PBH scenario!

PBHs = LV BHs?

PBHs = CDM?

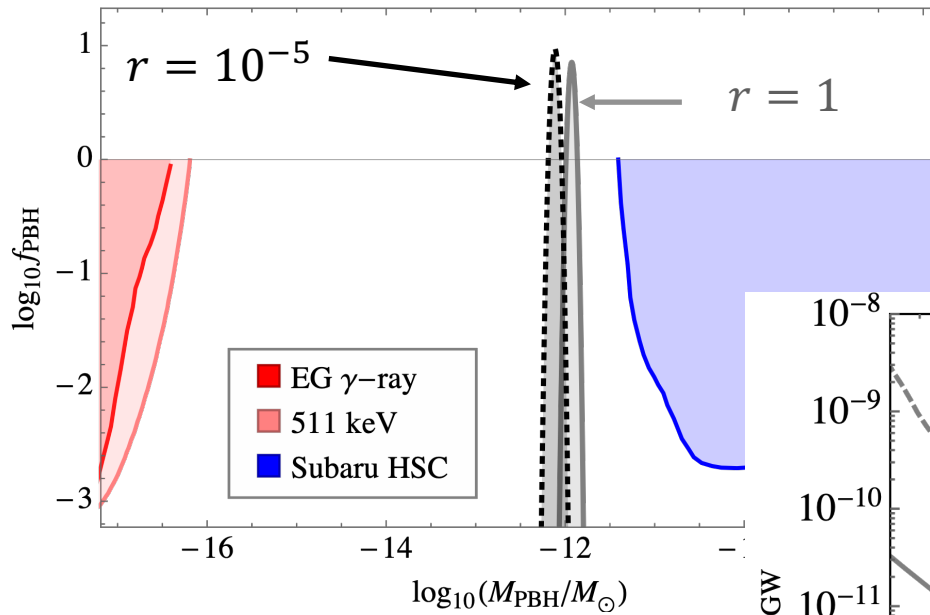


- PBHs = LV BHs scenario is already constrained by NANOGrav(PTA)

Cai, Pi, Wang & Yang 1907.06372

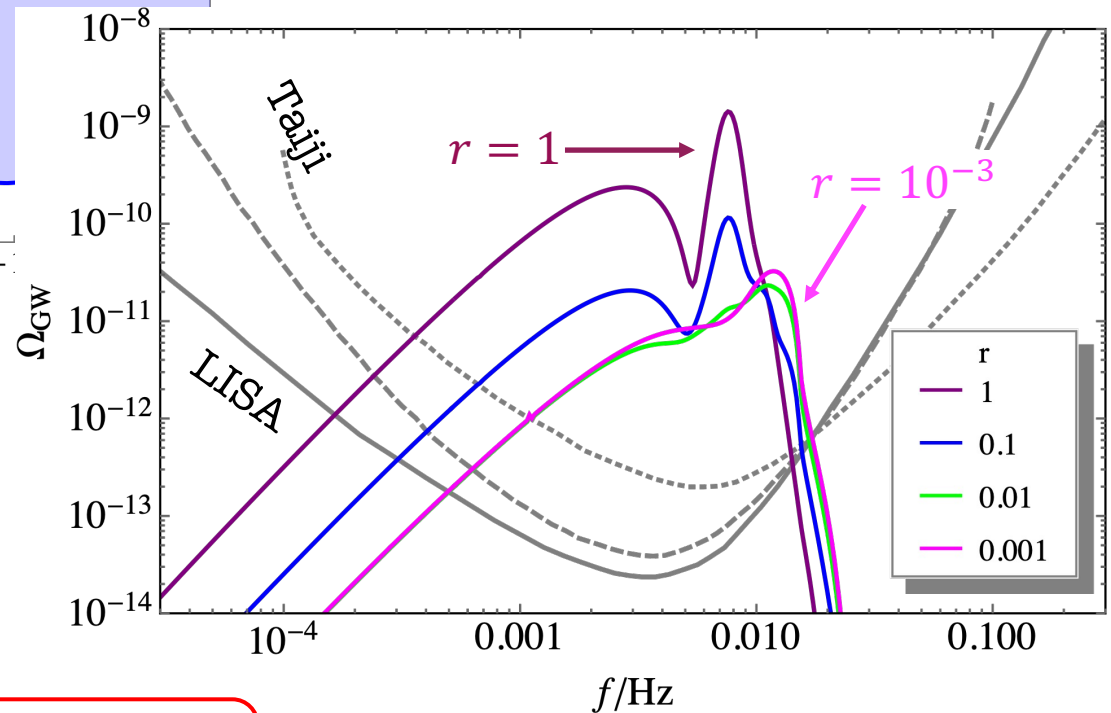
iGWs in non-minimal curvaton model

Pi & MS, 2112.12680



PBH mass function

width: $\frac{\Delta M}{M} \approx 0.1$



iGW spectrum

PBH=CDM scenario can be clearly verified or falsified by LISA!

summary

- various **inflation models** can lead to **PBH formation**
- **late stage** of inflation can be probed by **PBHs** and the associated **secondary/induced GWs**
- PBHs may be **formed during inflation**
- **PBHs** may play central roles in **GW cosmology**