

Black Holes as Dark Matter



Korean Institute for Advanced Studies (KIAS)

Thursday October 30, 2025



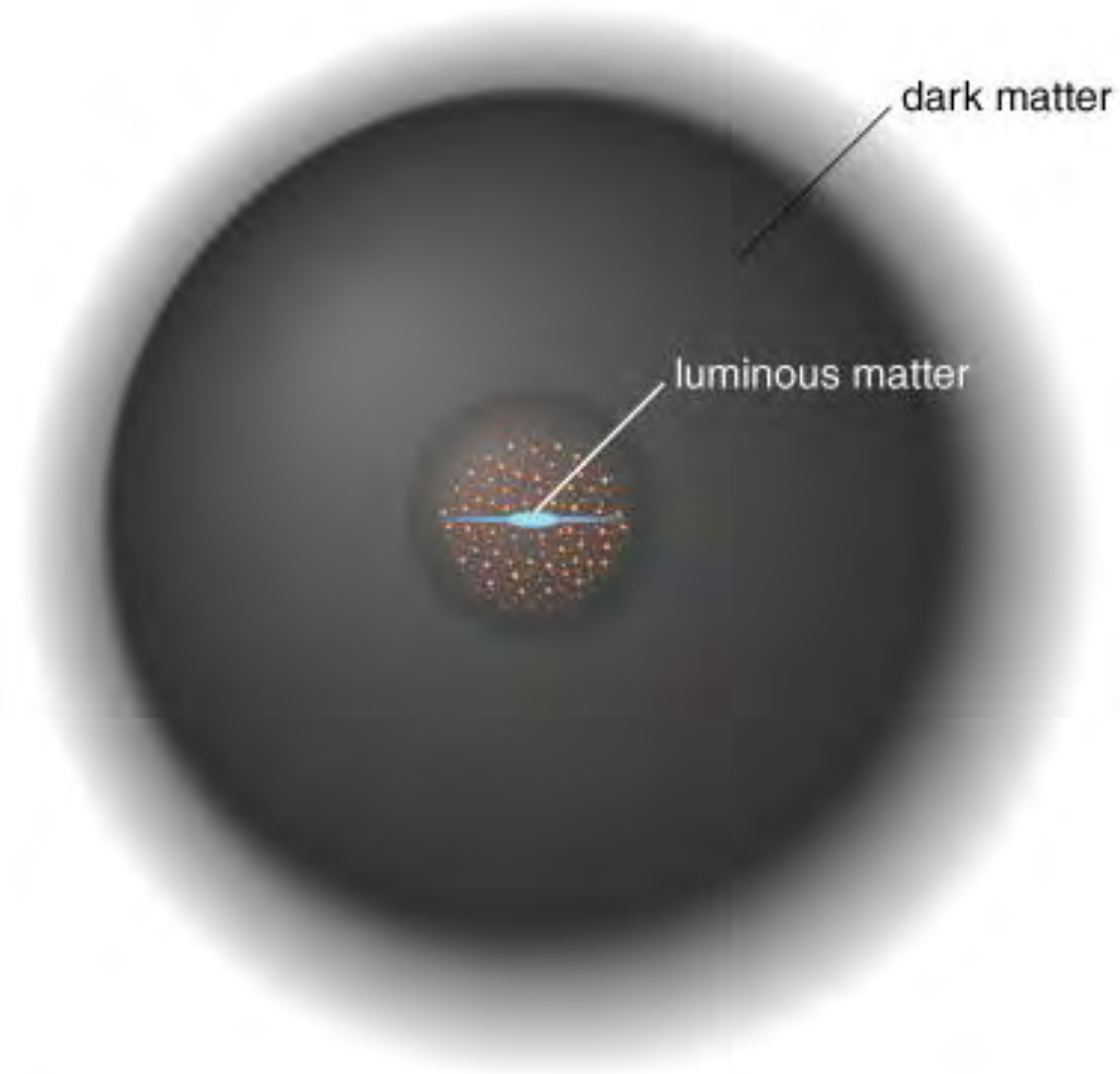
I wish I could be in **Korea!**

[next time, please!]

Contact me at profumo@ucsc.edu

If you are in **California**, reach out!





5/6

**a new
elementary particle**

what is an elementary particle?

**what is an
elementary particle?**

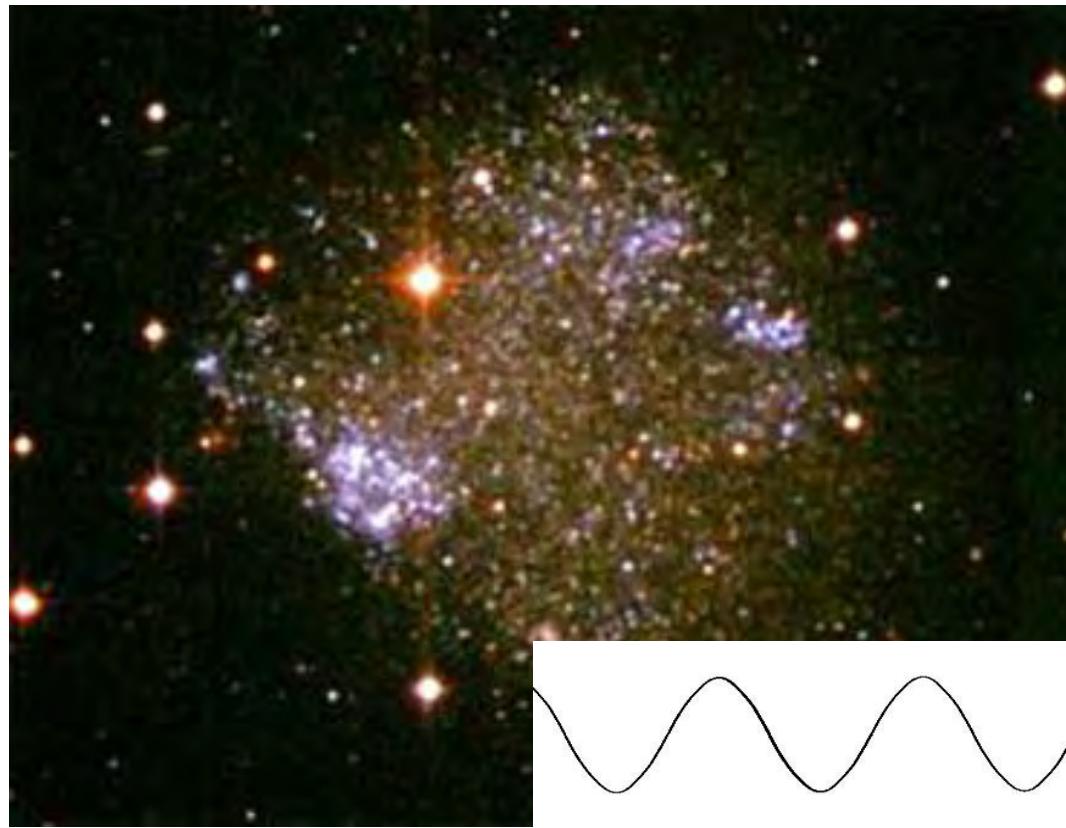
**an irreducible, unitary
representation of the
Poincaré Group**

(m, J)

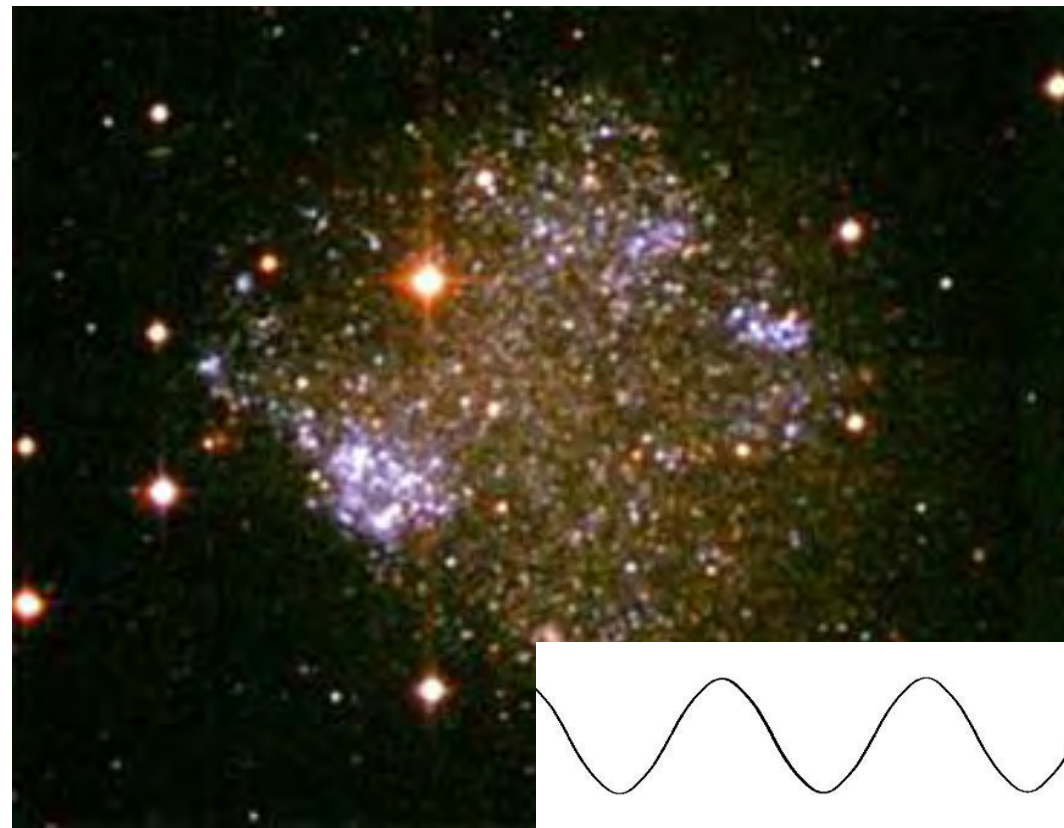
what do we know
about **m** and **J**?



Eugene Wigner (1902-95)

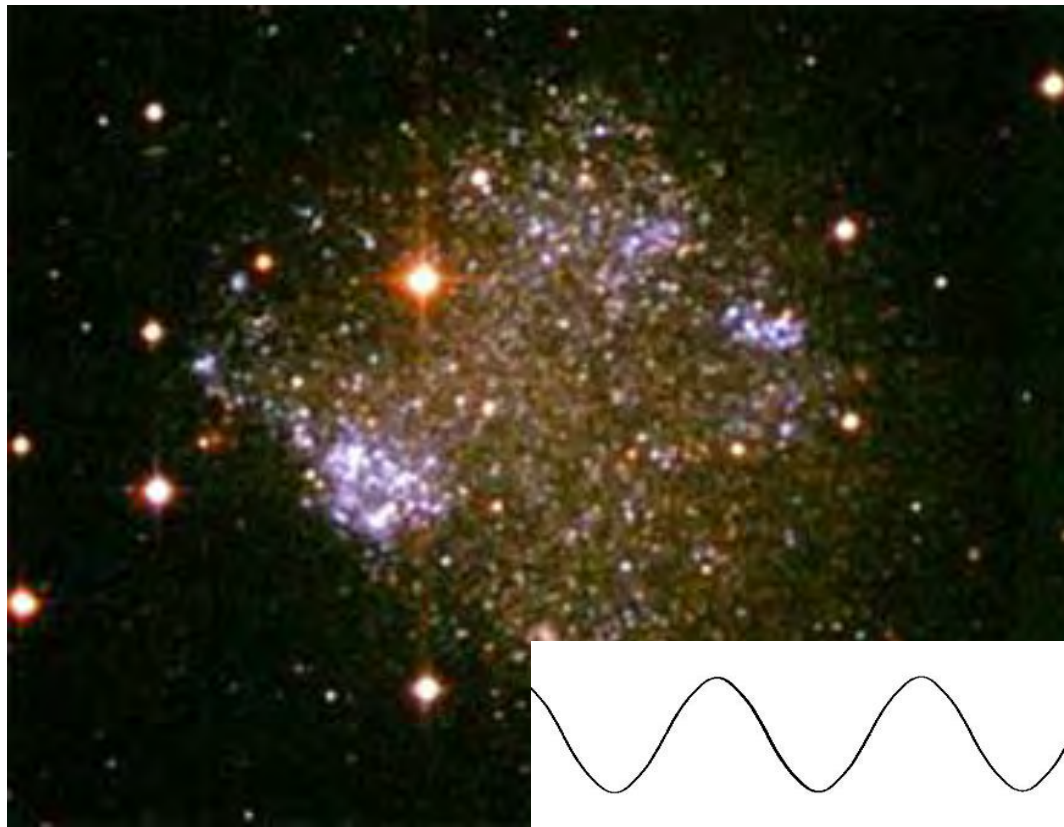


quantum effects must be
smaller than halos!



$$\lambda_{\text{DB}} = h/(mv) < 1 \text{ kpc}$$

$$\lambda_{\text{DB}} = 0.3 \text{ cm} (1 \text{ eV}/m) < 1 \text{ kpc} = 3 \times 10^{21} \text{ cm}$$



$m > 10^{-22} \text{ eV}$ [bosons!]

$$\lambda_{\text{DB}} = 0.3 \text{ cm} (1 \text{ eV}/m) < 1 \text{ kpc} = 3 \times 10^{21} \text{ cm}$$

State of the art: Koulen, Profumo & Smyth 2403.1901

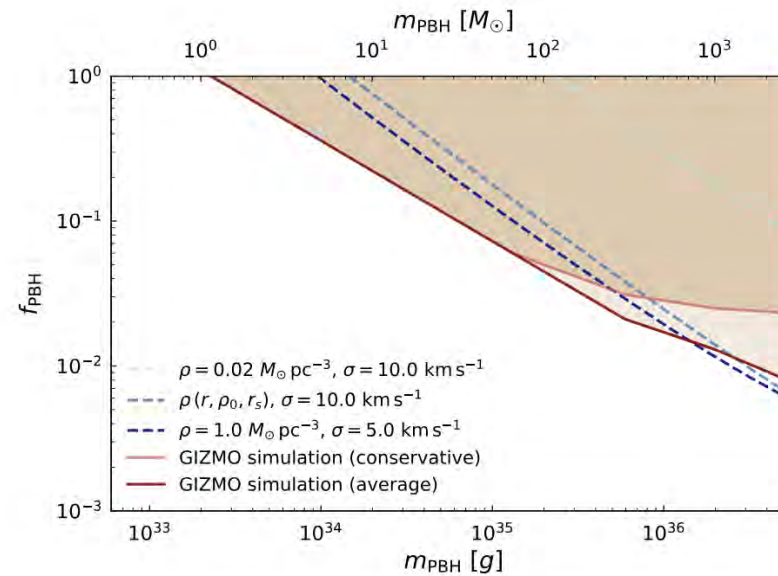
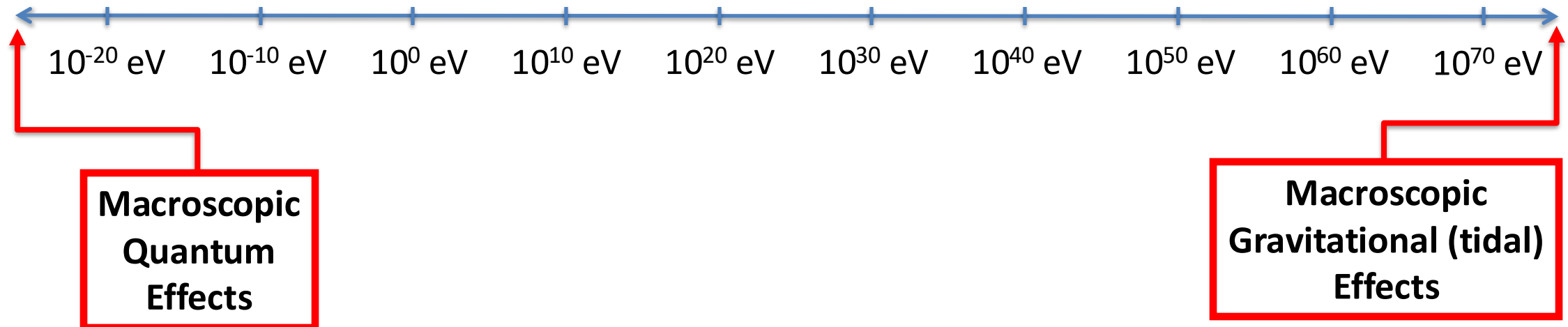
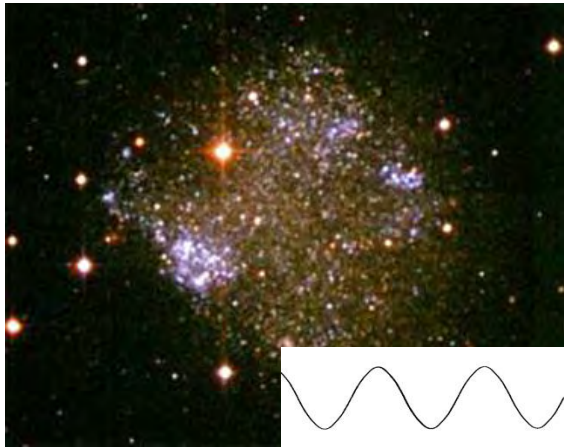


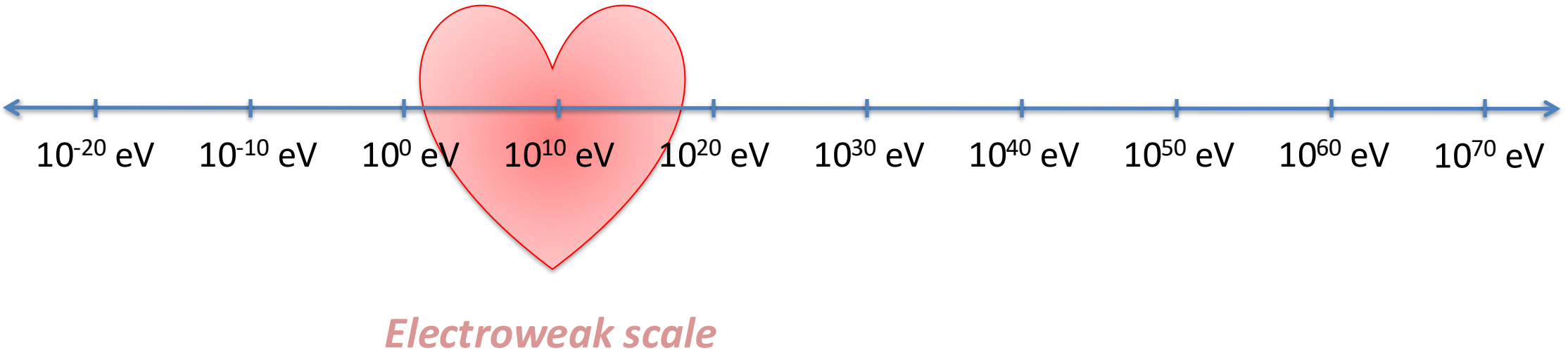
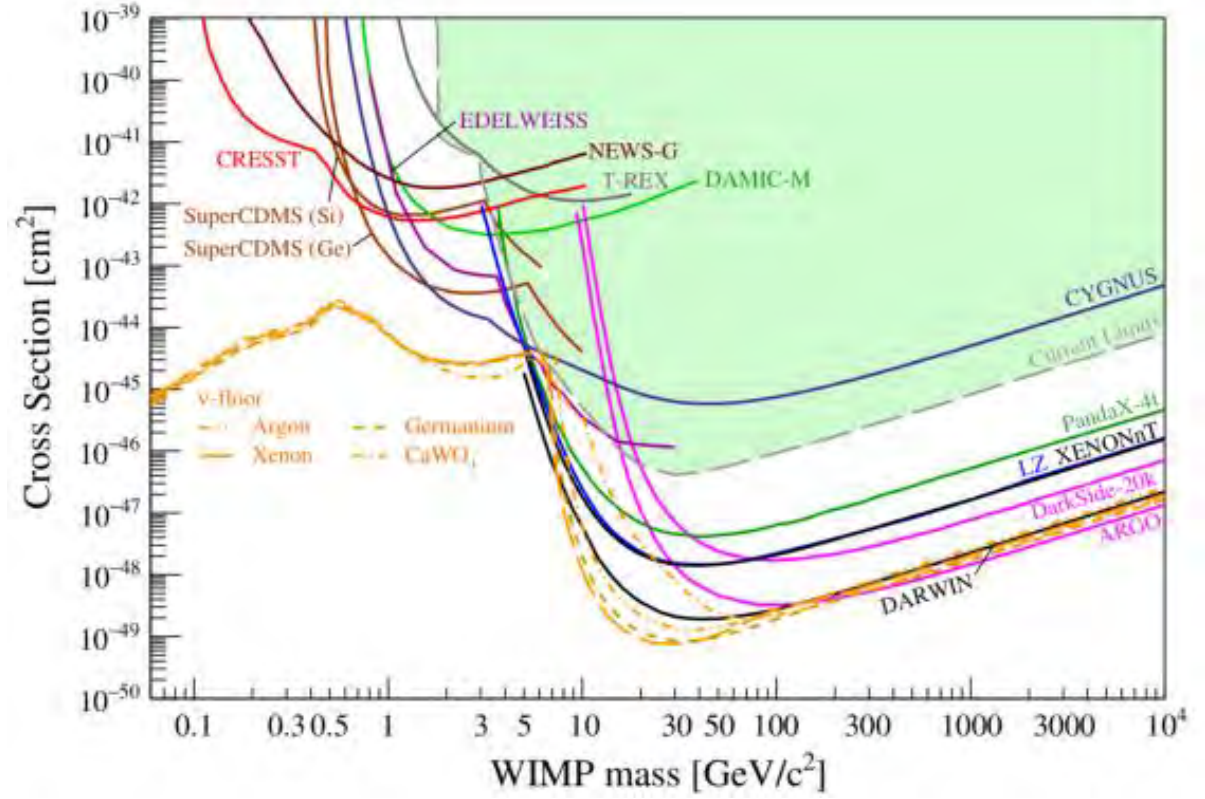
Figure 5. Constraints on f_{PBH} as a function of m_{PBH} . Limits are derived using the average r_h over 50 simulations for each mass. Also shown is a conservative case using r_h one standard deviation below the mean. The dashed lines are the semi-analytically derived constraints.

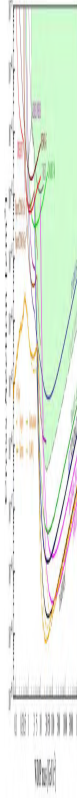
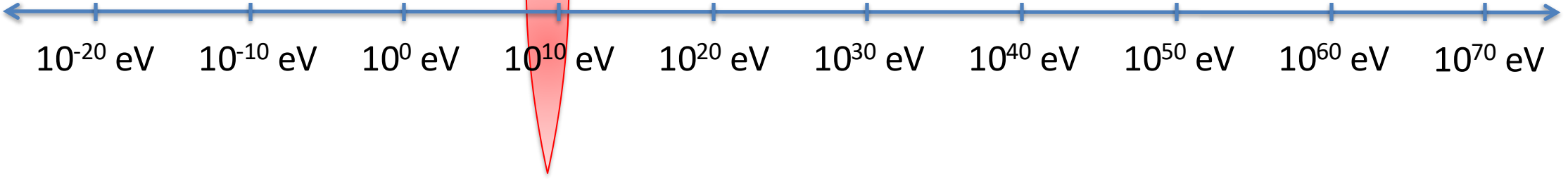


m < *few* solar masses $\sim 10^{67}$ eV

(m, J)







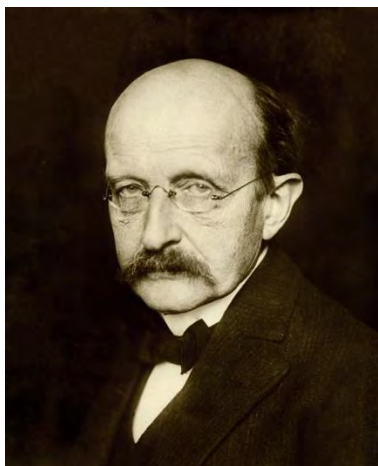
$$G = \frac{hc}{2\pi M_{\text{Pl}}^2}$$

$$R_s = \frac{2Gm}{c^2}$$

Schwarzschild radius

$$\lambda = \frac{h}{mc}$$

Compton wavelength



$$R_s(M_{\text{Pl}}) = \frac{2hcM_{\text{Pl}}}{2\pi c^2 M_{\text{Pl}}^2} = \frac{h}{\pi M_{\text{Pl}} c} \sim \lambda(M_{\text{Pl}})$$

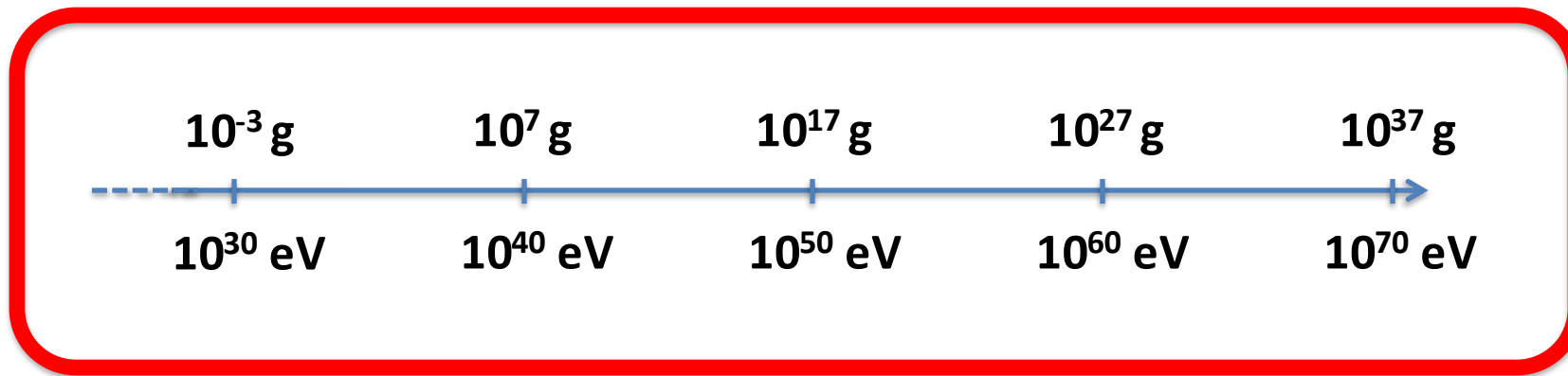
Elementary Particles?



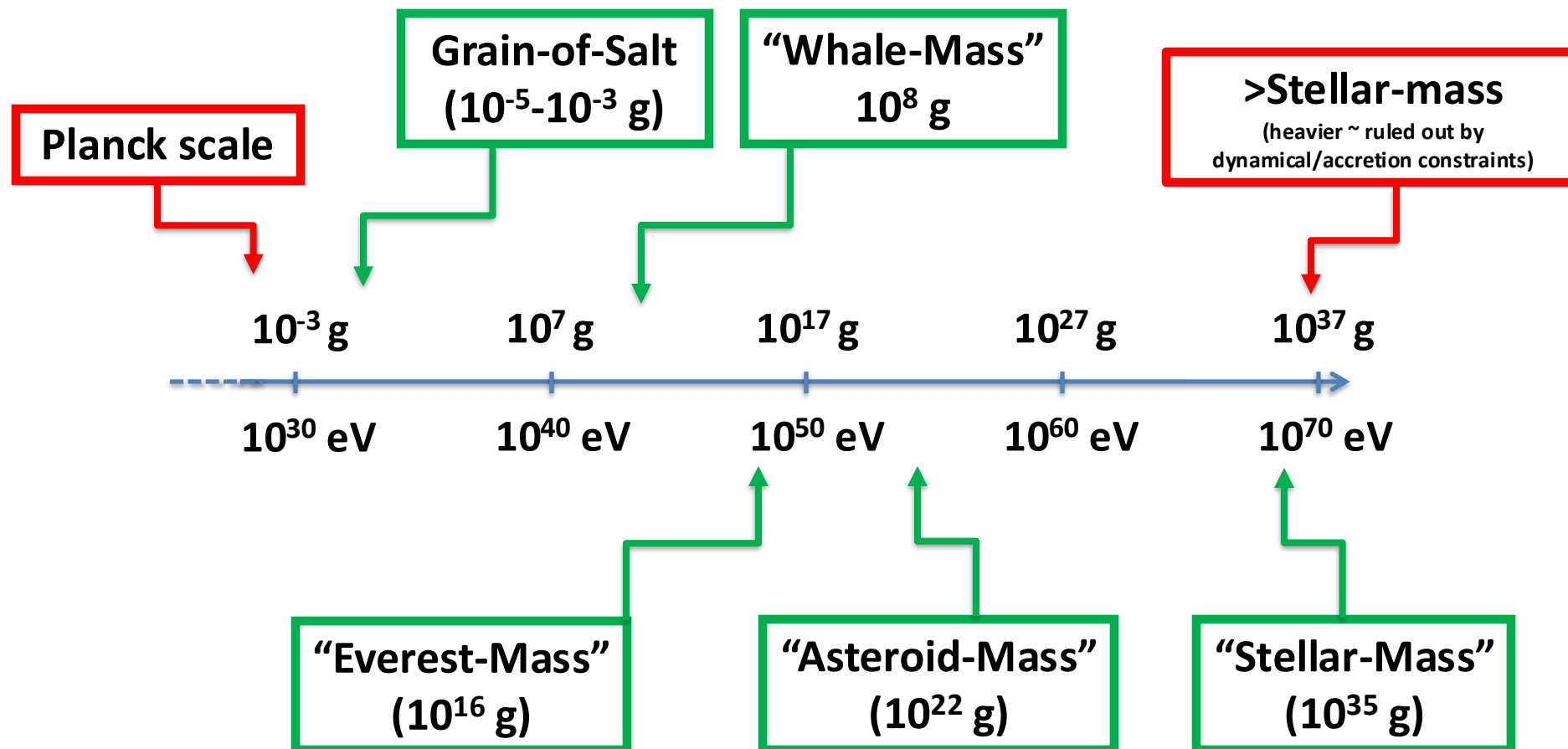
(m, J)

My talk today

Black Holes as Dark Matter

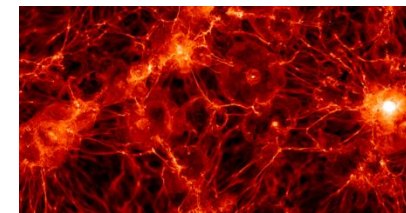


(m, J)

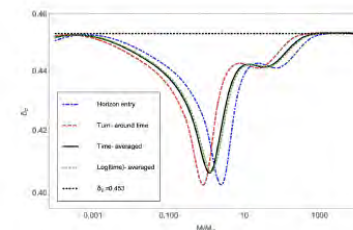


What I will **NOT** talk about today: how **Non-stellar** **Black Holes** form

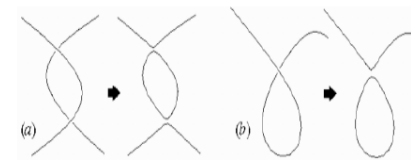
➤ Collapse of large density perturbations



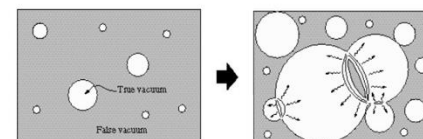
➤ Pressure Reduction



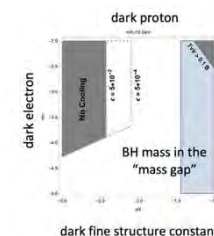
➤ Cosmic Strings Loops



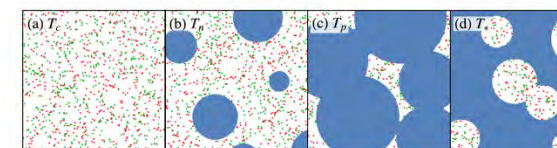
➤ Bubble Collisions



➤ Dark Stars Collapse*



➤ Collapse of trapped fermions



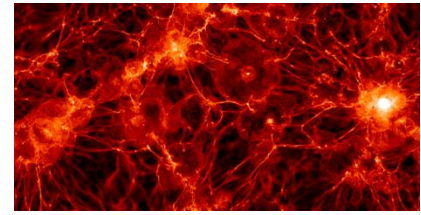
➤ Gravo-thermal collapse of self-interacting DM cores**



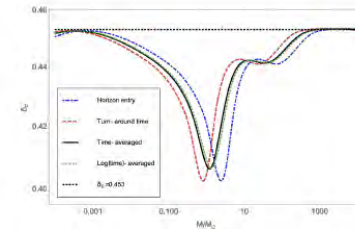
*Fernandez, Profumo+,
2208.08557

**Profumo et al, 2410.17480

Different formation mechanism lead to different mass functions!

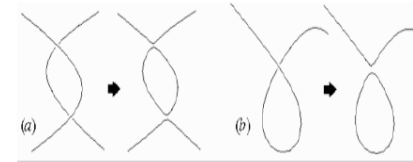


➤ Collapse of large density perturbations

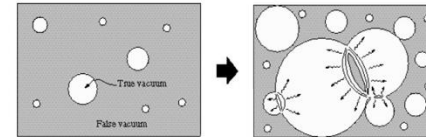


➤ Pressure Reduction

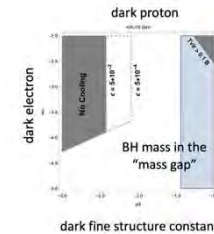
➤ Cosmic Strings Loops



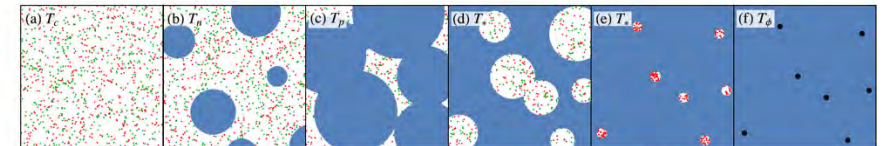
➤ Bubble Collisions



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➤ Gravo-thermal collapse of self-interacting DM cores**

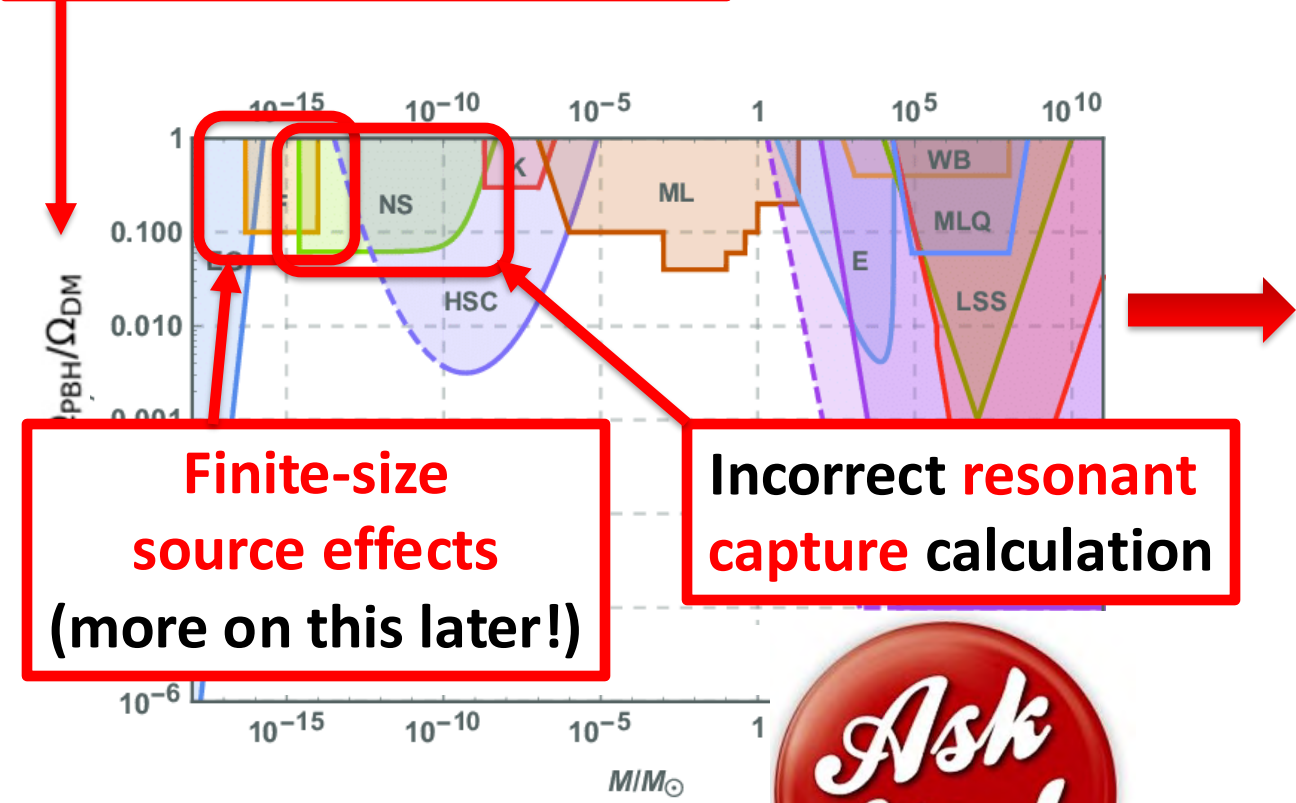
*Fernandez, Profumo+,
2208.08557

**Profumo et al, 2410.17480

f_{PBH} : Fraction of the dark matter consisting of non-stellar BHs of given mass

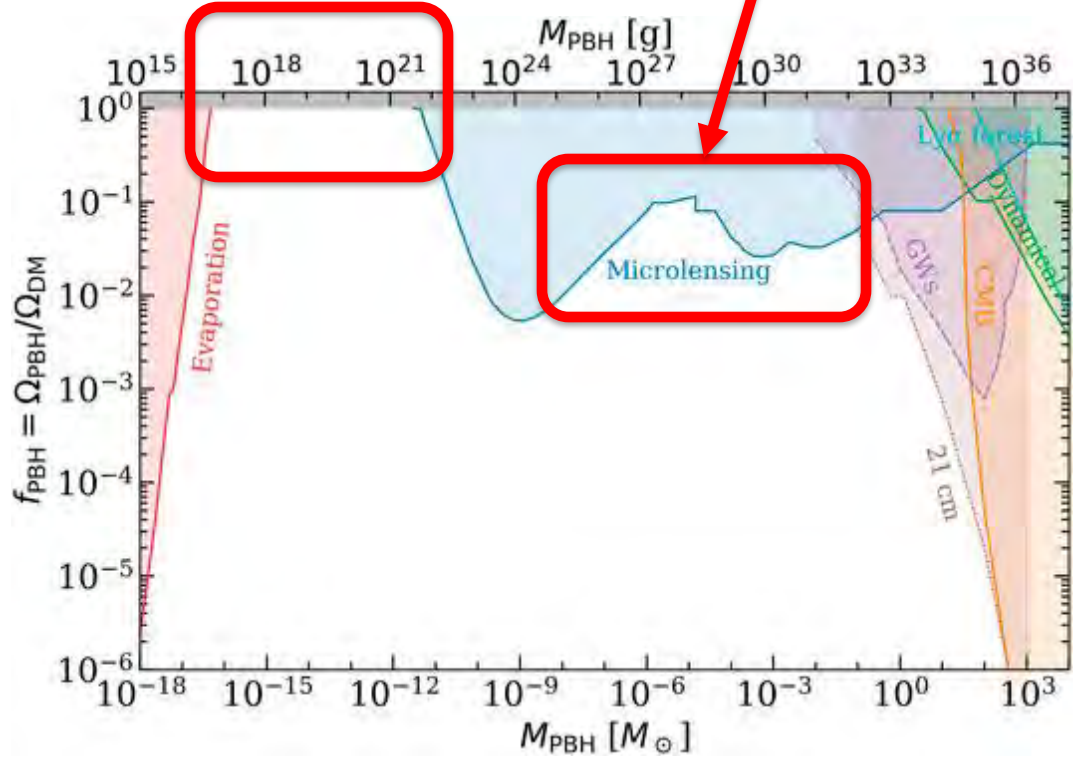
Asteroid-mass PBH can be **100%** of the dark matter

Earth mass and above PBH can be **1-10%** of the dark matter



Finite-size source effects (more on this later!)

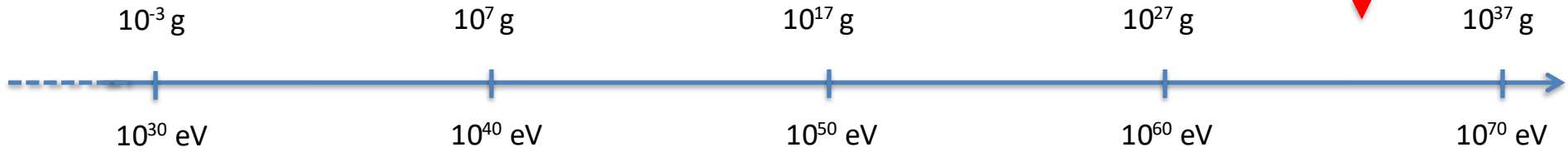
Incorrect resonant capture calculation



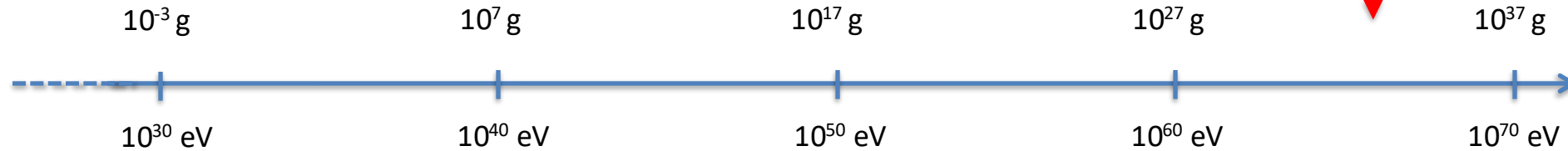
<2019

Now

**sub-“Stellar-Mass”
($<10^{33}$ g)
Black Holes**



sub-“Stellar-Mass”
($<10^{33}$ g)
Black Holes



✓ Is there an **unmistakable signature** for PBH as DM?

Smoking gun: **BH merger** with a **sub-TOV** limit ($\sim 2xM_{\text{sun}}$) [or a **Sub-Solar-Mass** merger]

(possible “new physics” **imposters** such as **strange quark stars**, anomalously **light** or dark sector **white dwarfs** or **neutron stars** – but with different GW signals from **tidal effects** [Begnoni and Profumo*,])



Primordial black holes versus their impersonators at gravitational wave observatories

Andrea Begnoni^{a,b} and Stefano Profumo^{c,d}

^aDipartimento di Fisica Galileo Galilei, Università di Padova, I-35131 Padova, Italy

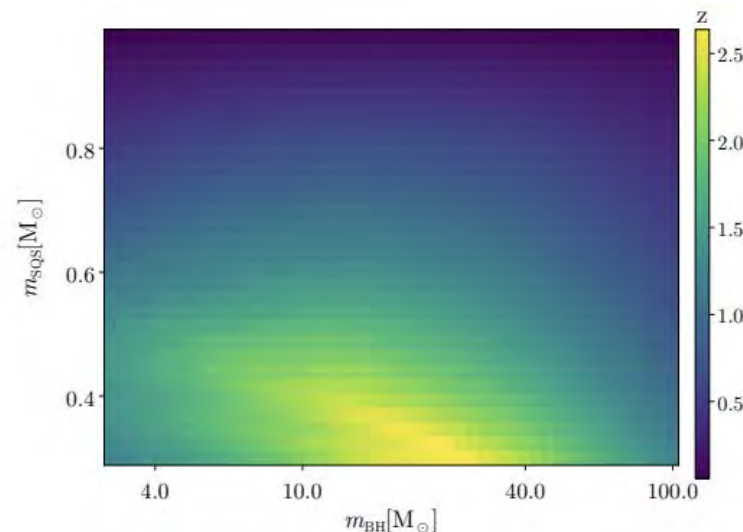
^bINFN Sezione di Padova, I-35131 Padova, Italy

^cSanta Cruz Institute for Particle Physics,
Santa Cruz, CA, 95064, USA

^dDepartment of Physics, University of California, Santa Cruz
Santa Cruz, CA, 95064, USA

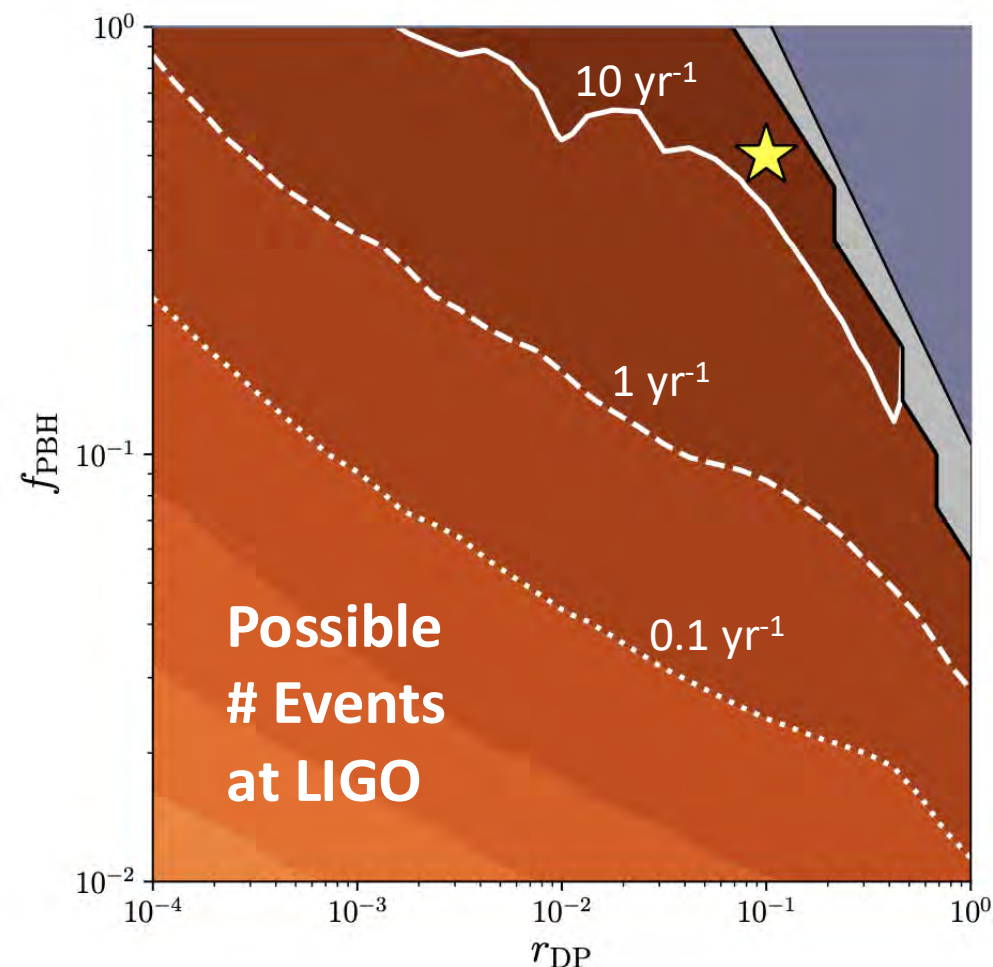
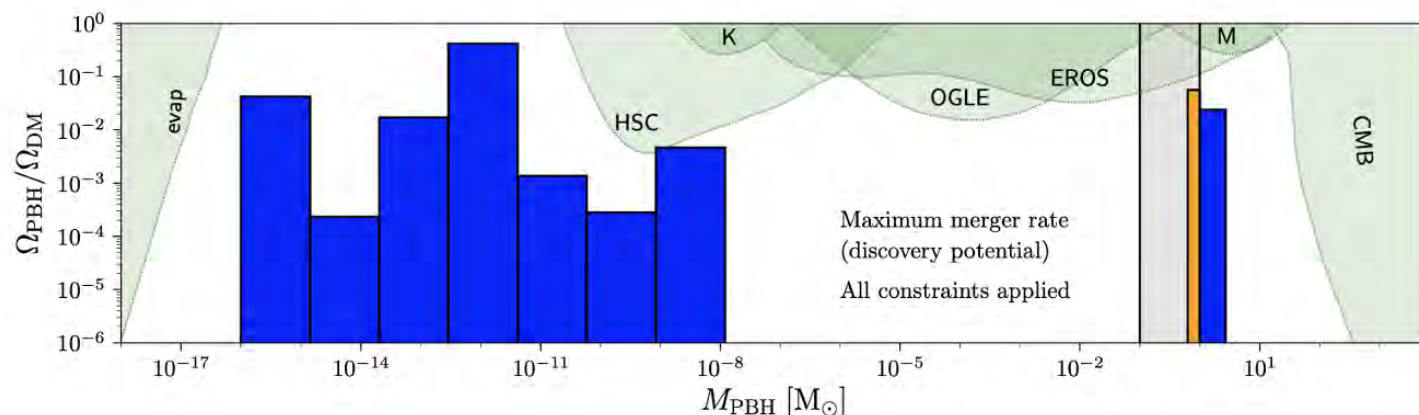
E-mail: andrea.begnoni@phd.unipd.it, profumo@ucsc.edu

Abstract. The detection of primordial black holes (PBHs) would mark a major breakthrough, with far-reaching implications for early universe cosmology, fundamental physics, and the nature of dark matter. Gravitational wave observations have recently emerged as a powerful tool to test the existence and properties of PBHs, as these objects leave distinctive imprints on the



Smoking gun: **BH merger** with a **sub-TOV** limit ($\sim 2xM_{\text{sun}}$) [or a **Sub-Solar-Mass** merger]

- We calculated the max and min **event rate** for a given f_{PBH} and fraction of “**goldilocks**” (light+detectable) PBH, and it can be **sizable**!*



(fraction of light+detectable PBHs)

Smoking gun: **BH merger** with a **sub-TOV** limit ($\sim 2xM_{\text{sun}}$) [or a **Sub-Solar-Mass** merger]

- We calculated the max and min **event rate** for a given f_{PBH} and fraction of “**goldilocks**” (light+detectable) PBH, and it can be **sizable**!*
- LIGO-VIRGO-KAGRA **searches** are ongoing, **no confirmed events** in O3/O4
- ...**two candidate events**, controversial and low-confidence, reported by independent groups!

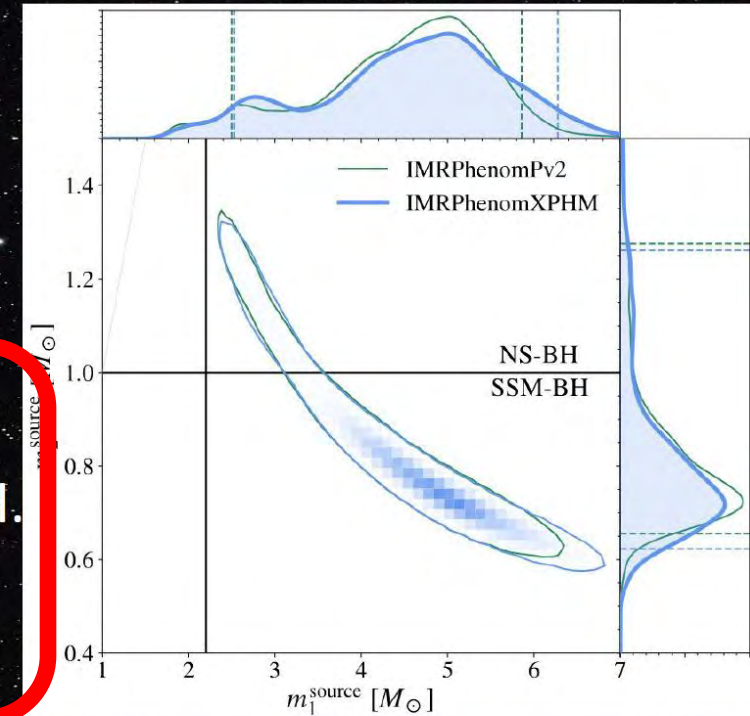
SubSolarMass (SSB) 170401

[given time stamp, hopefully
not April's fool joke]

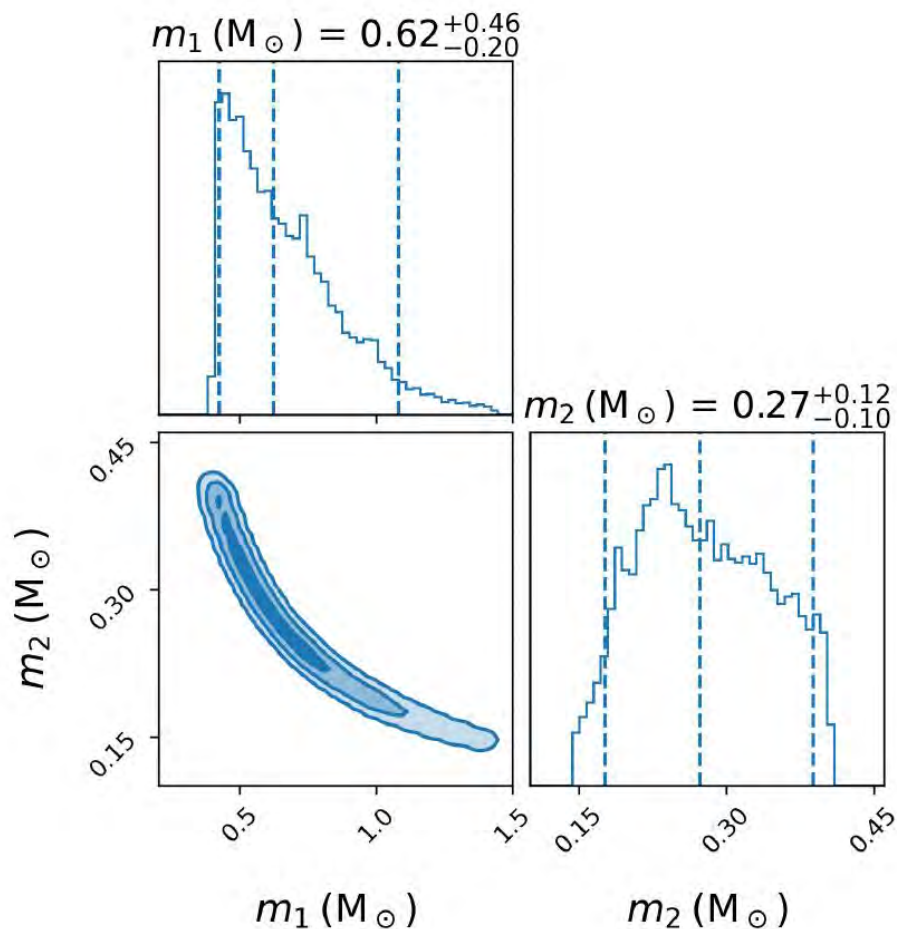
- For 16% of the posterior we have $m_1 > 2.2M_\odot$ and $m_2 > 1.0M_\odot$, i.e. masses compatible with an NS-BH.
 - The primary would be a BH in the mass gap
 - The secondary could be a NS or a BH

- For 84% of the posterior we have $m_1 > 2.2M_\odot$ and $m_2 < 1.0M_\odot$, i.e. masses compatible with an SSM-BH.
 - The primary can be in the mass gap or above
 - The secondary is an SSM black hole.

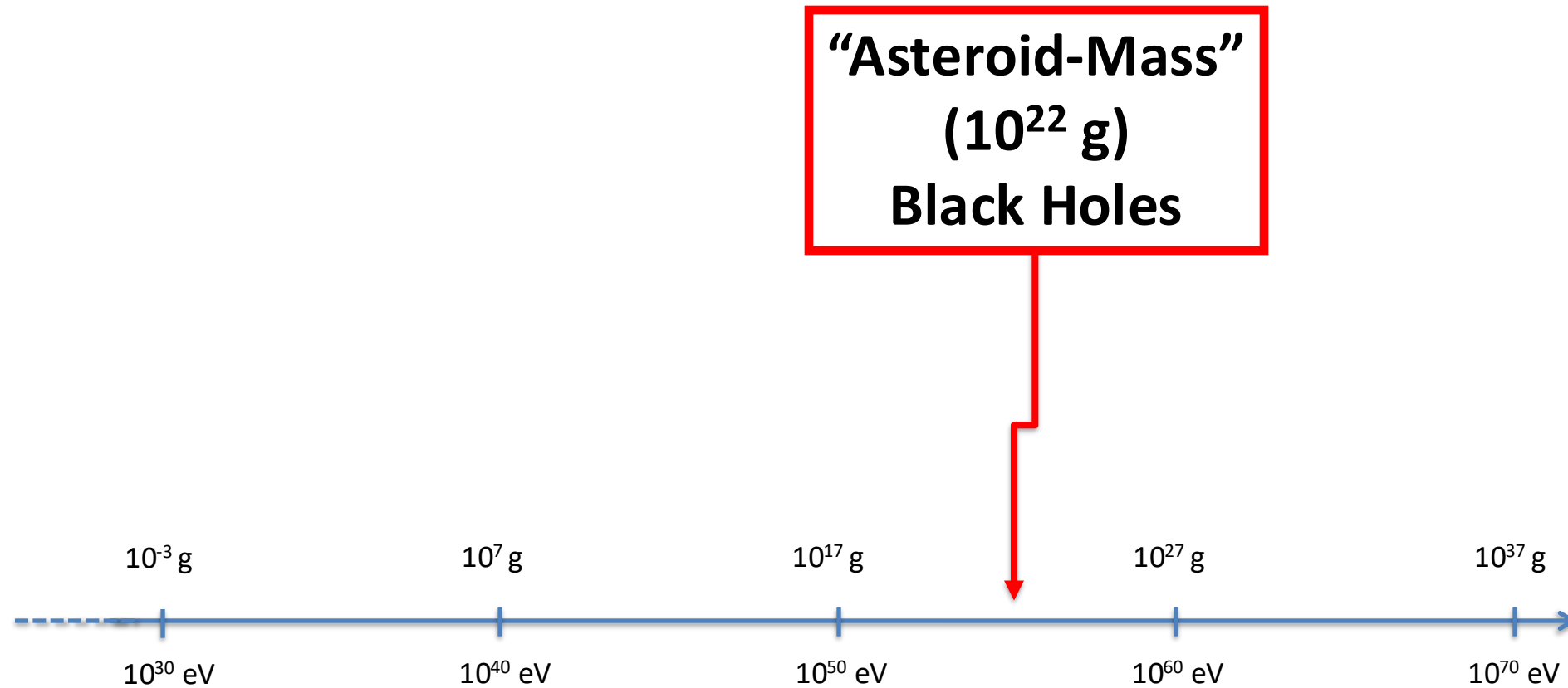
- In all cases, astrophysical models have problems generating a system with these masses



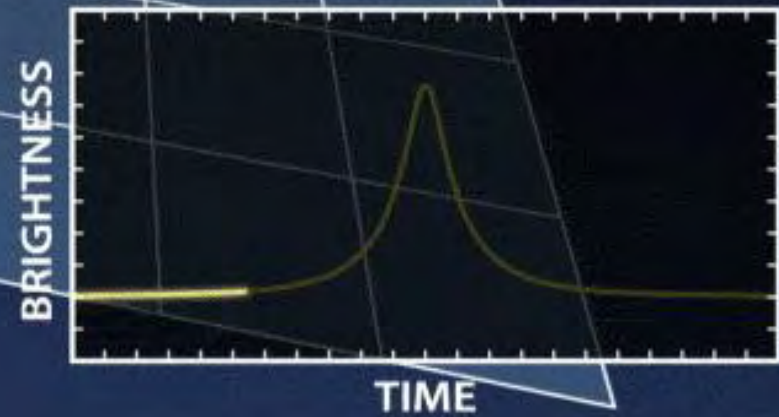
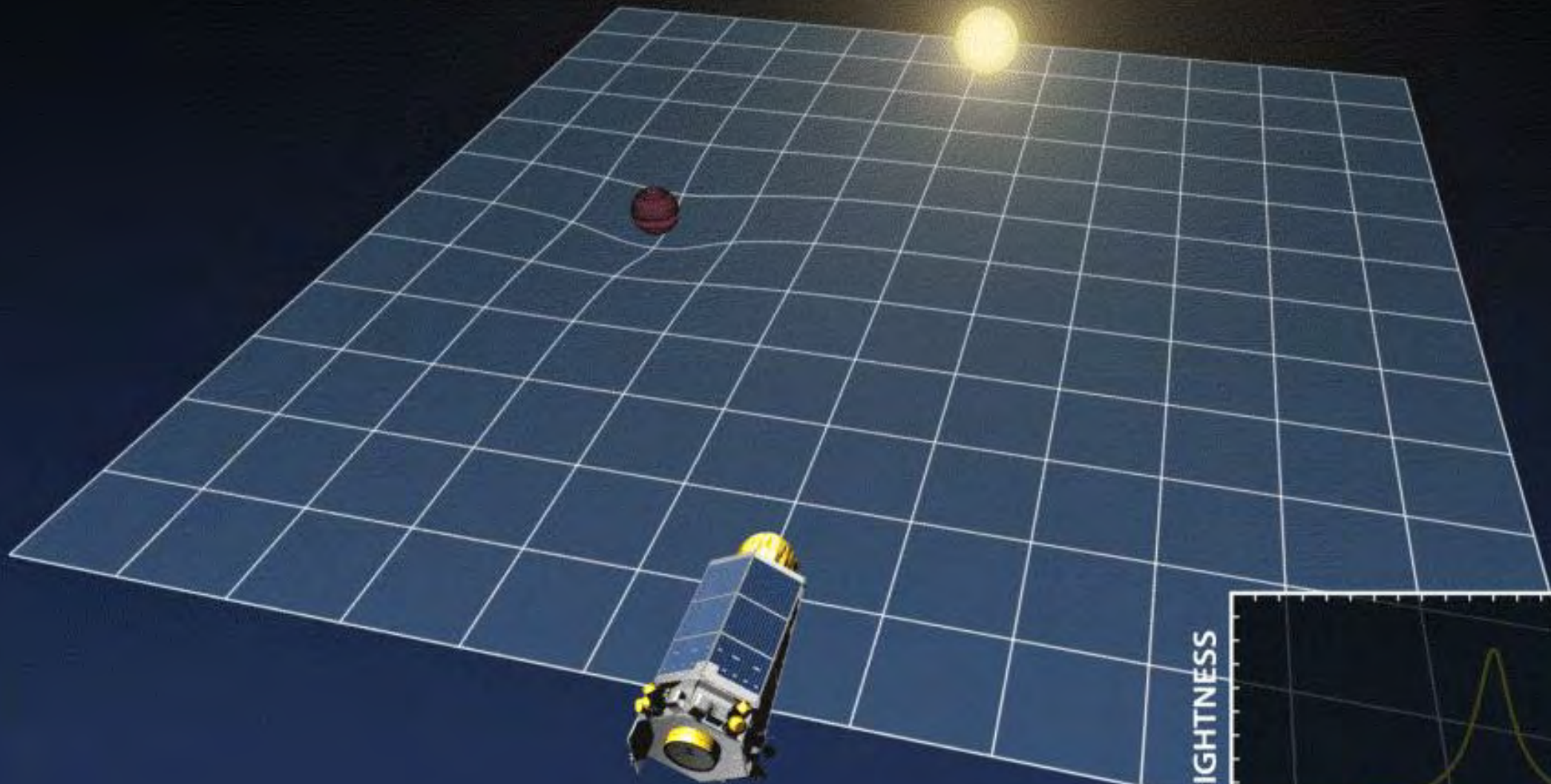
SubSolarMass (SSB) 200308



an in-depth analysis of one of the most significant candidates reported in the O3b search for SSB black hole binaries [12] with SNR = 8.90 and FAR = 0.20 yr⁻¹. Even if the candidate does not show enough significance to claim the firm detection of a gravitational wave event, it is of great interest to study and characterize the candidate. We also demonstrate that the ROQ method can be efficiently used to reduce the computational cost of the PE for such long signals. The inferred masses show that SSM200308, if coming from a GW event, is consistent with a binary of two SSB black holes; $m_1 = 0.62^{+0.46}_{-0.20} M_\odot$ and $m_2 = 0.27^{+0.12}_{-0.10} M_\odot$ (90% credible intervals). Given the very low masses of the candidate's components, their neutron star nature seems disfavoured [47, 64]. The question of the nature of the source of SSM200308 therefore remains open. The unusual characteristics of SSM200308 could be explained by the two components being black holes of primordial origin. If SSM200308 is a real signal, we can expect that improved detector sensitivities and longer observing time would within a few years allow for the firm detection of an SSB black hole.



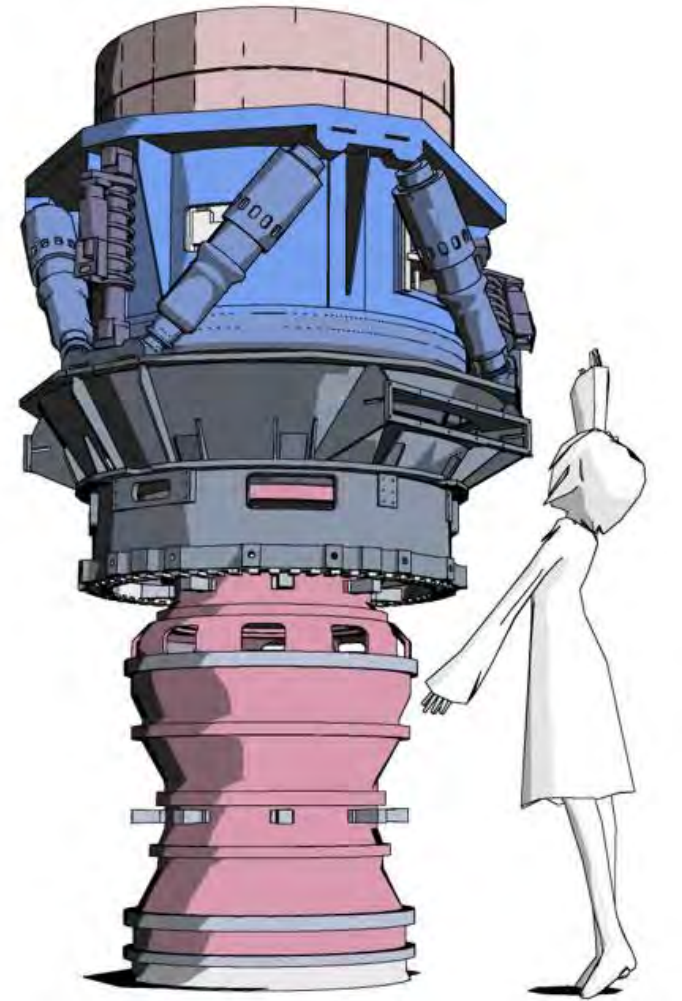
- ✓ Can be **100%** of the **Dark Matter!**
- ✓ **Microlensing** searches limitations
- ✓ How can we **detect** them?



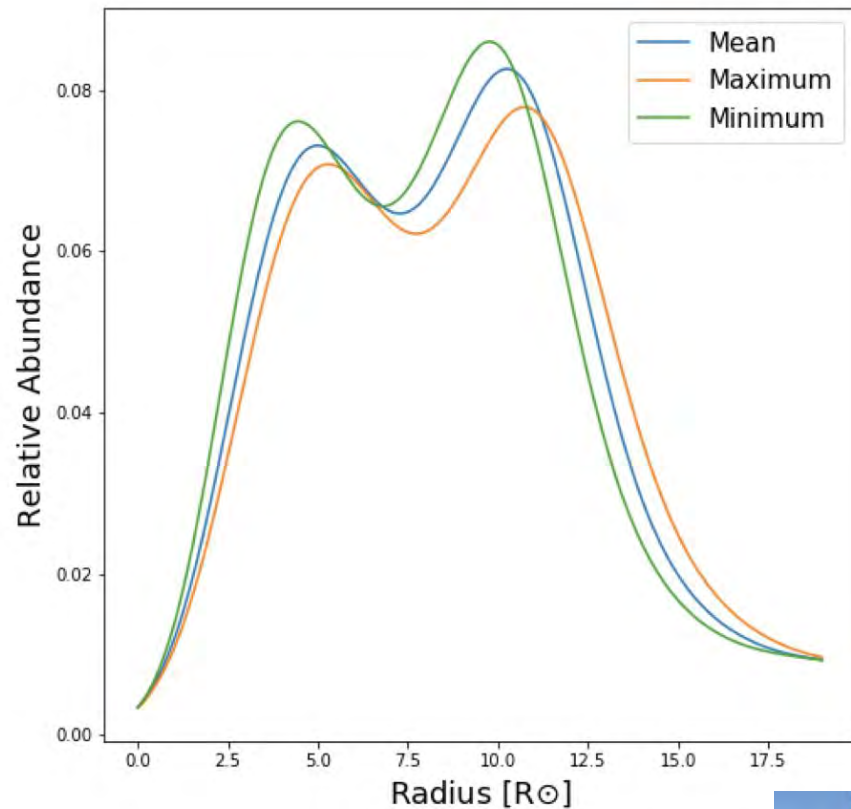
Strongest **Microlensing** Constraints on Light **Black Holes**



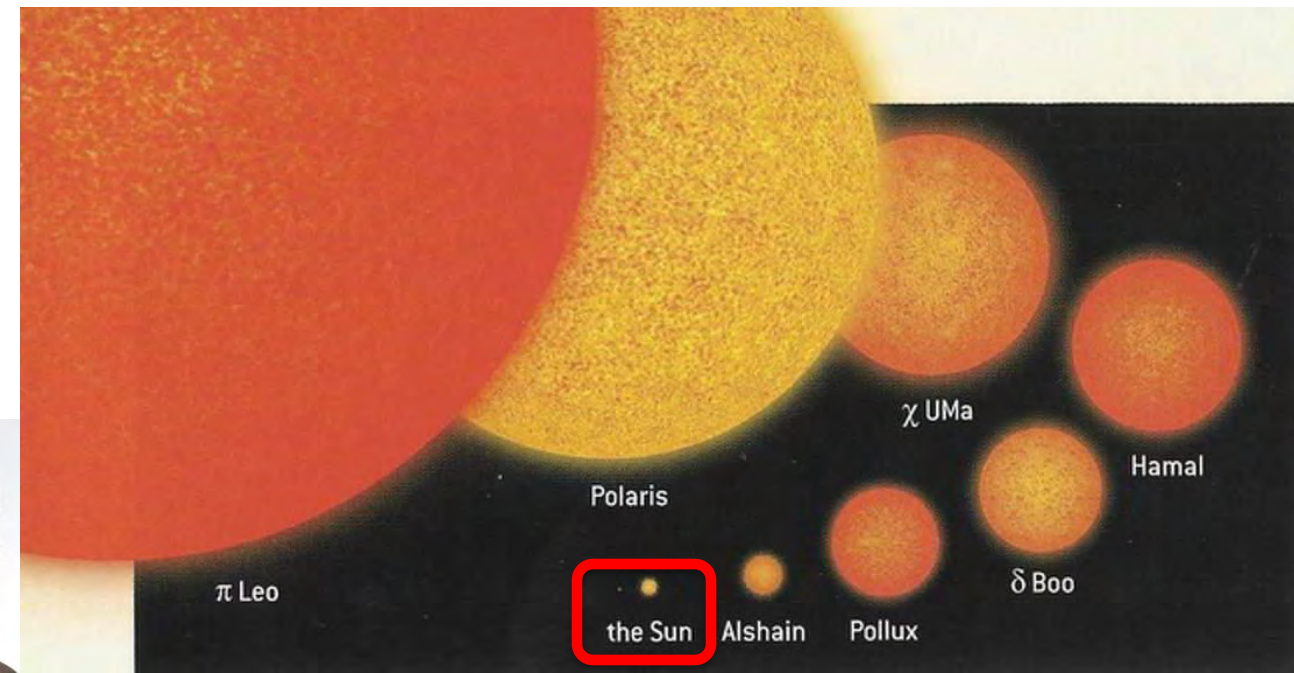
~ 7h
observations
of $O(10^6)$ stars
in **M31** with
Subaru HSC



HSC study assumes **all stars** in M31 are Sun-like...
but Sun-like stars are **too dim** for HSC!



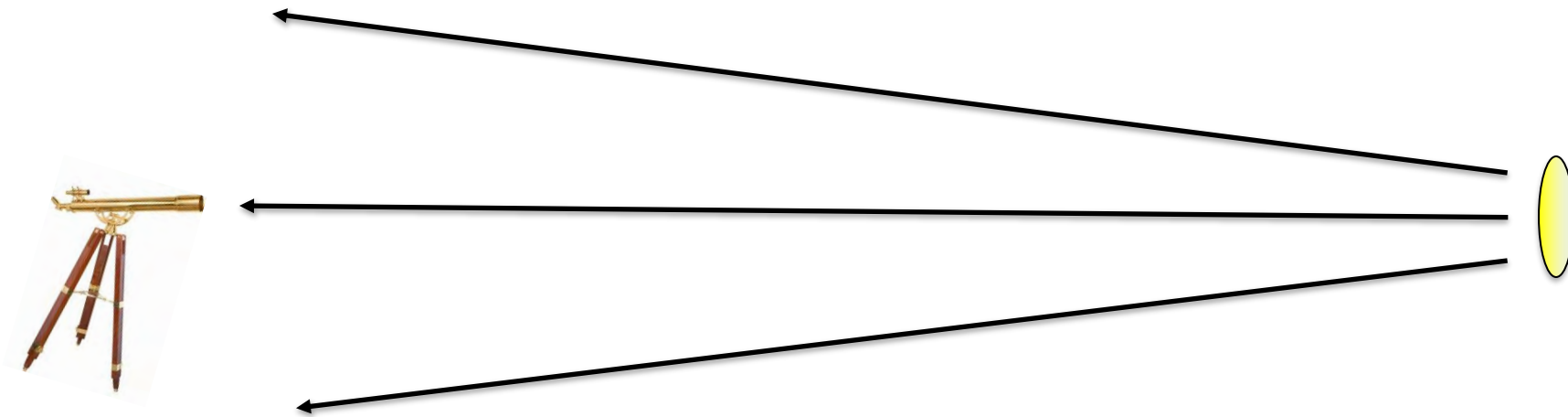
Stars that contribute to the
microlensing constraints are
 $\sim 100\times$ larger in the sky than the Sun!



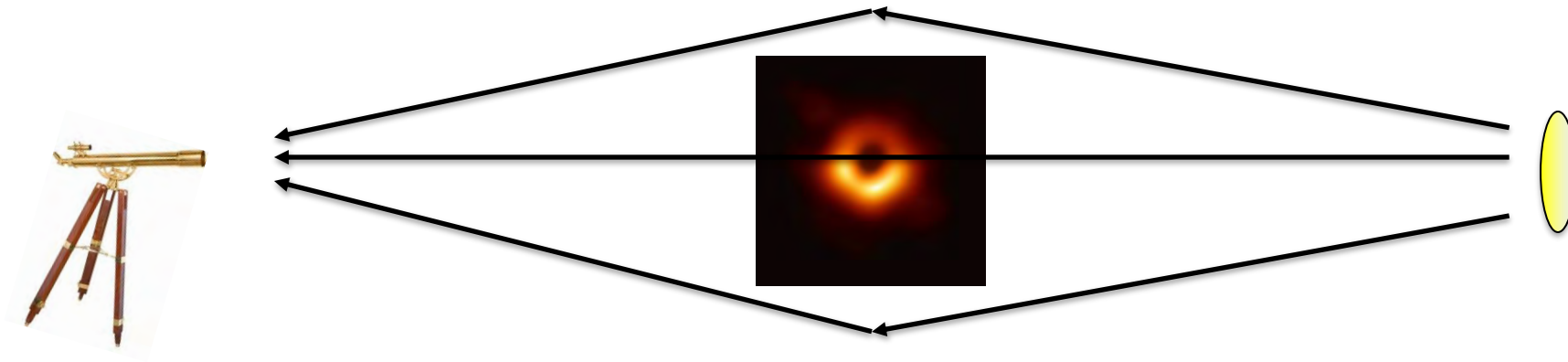
* Profumo, Smyth+ PRD 2020



Key effect previously neglected; finite-**source-size** effects*!

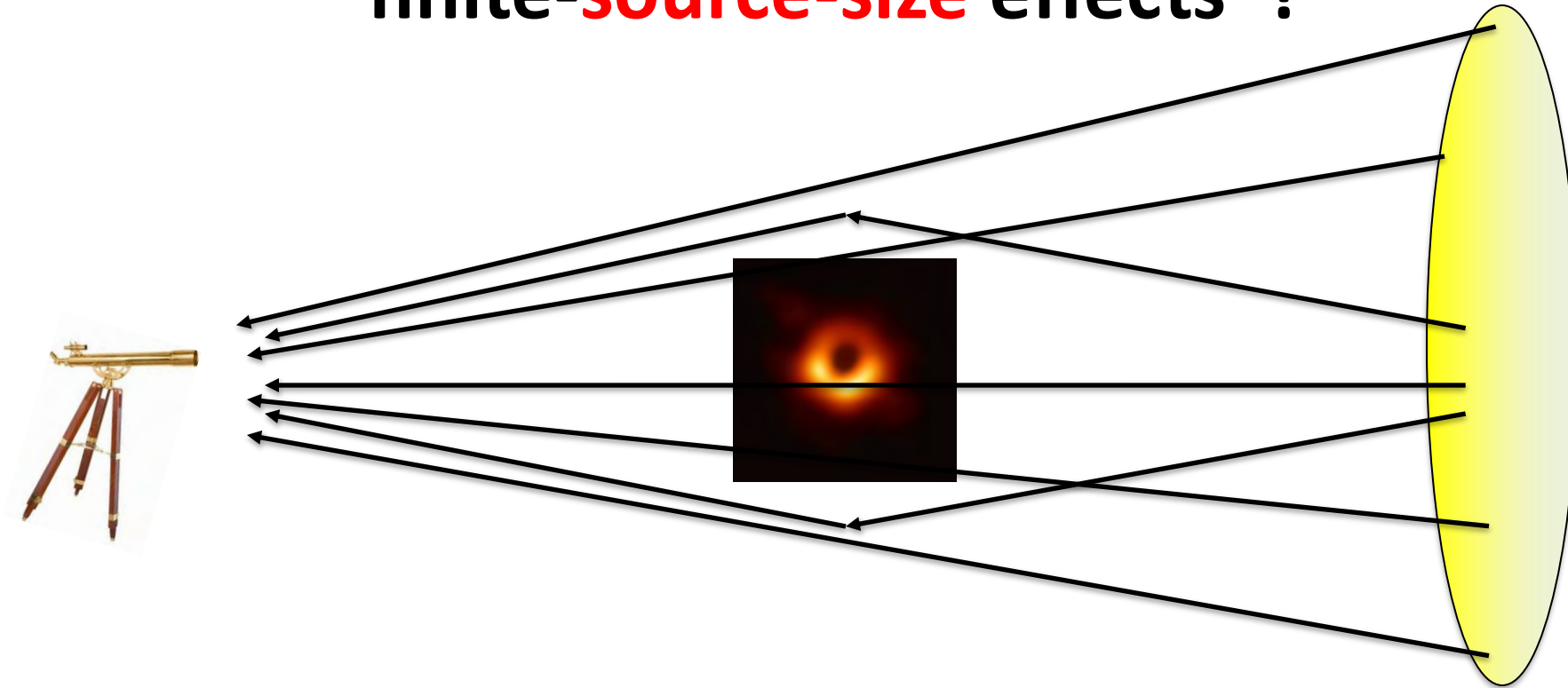


Key effect previously neglected; finite-**source-size** effects*!



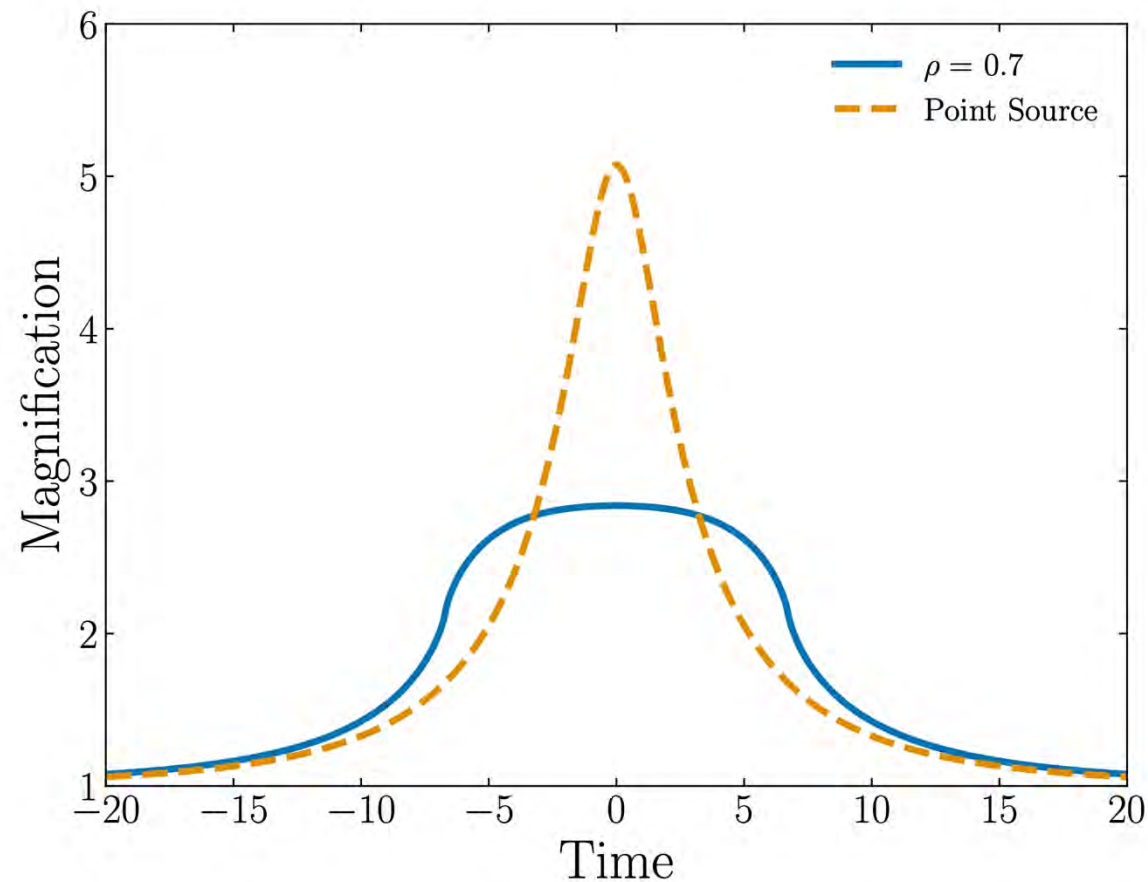
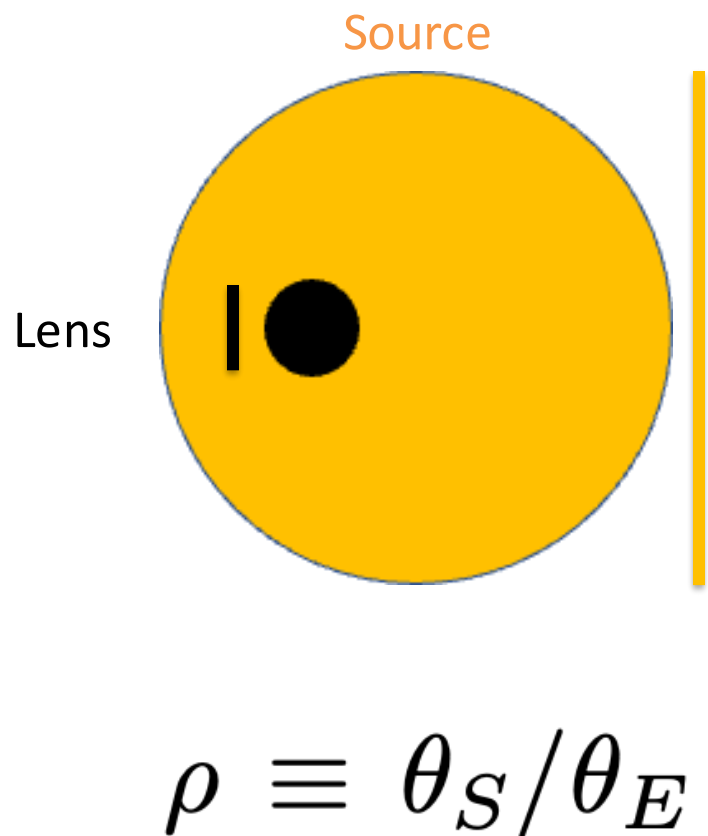
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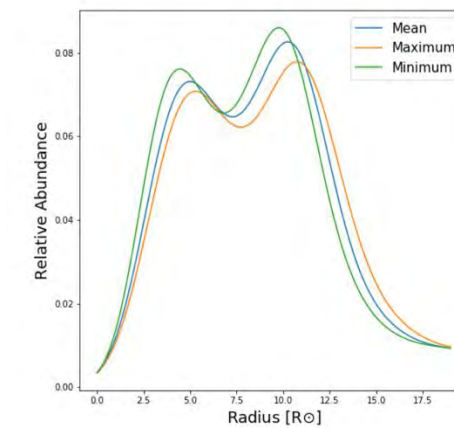
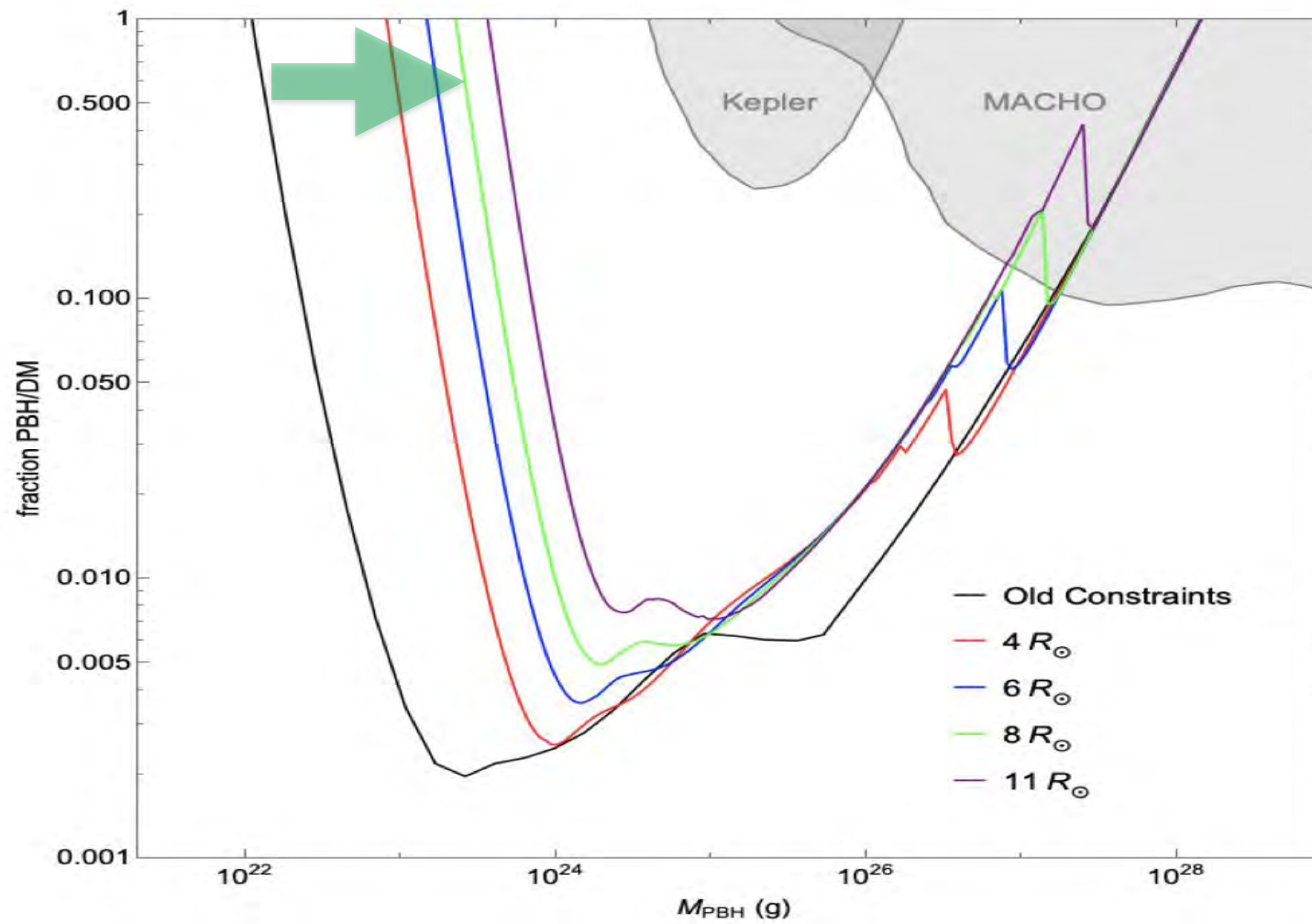
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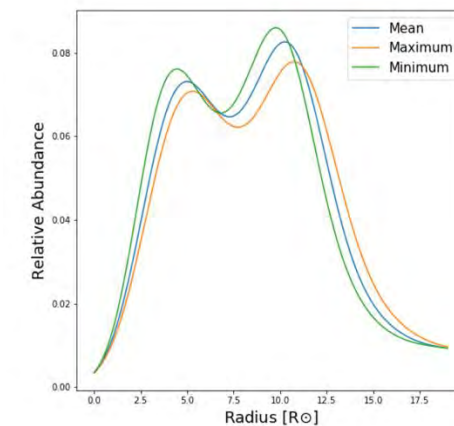
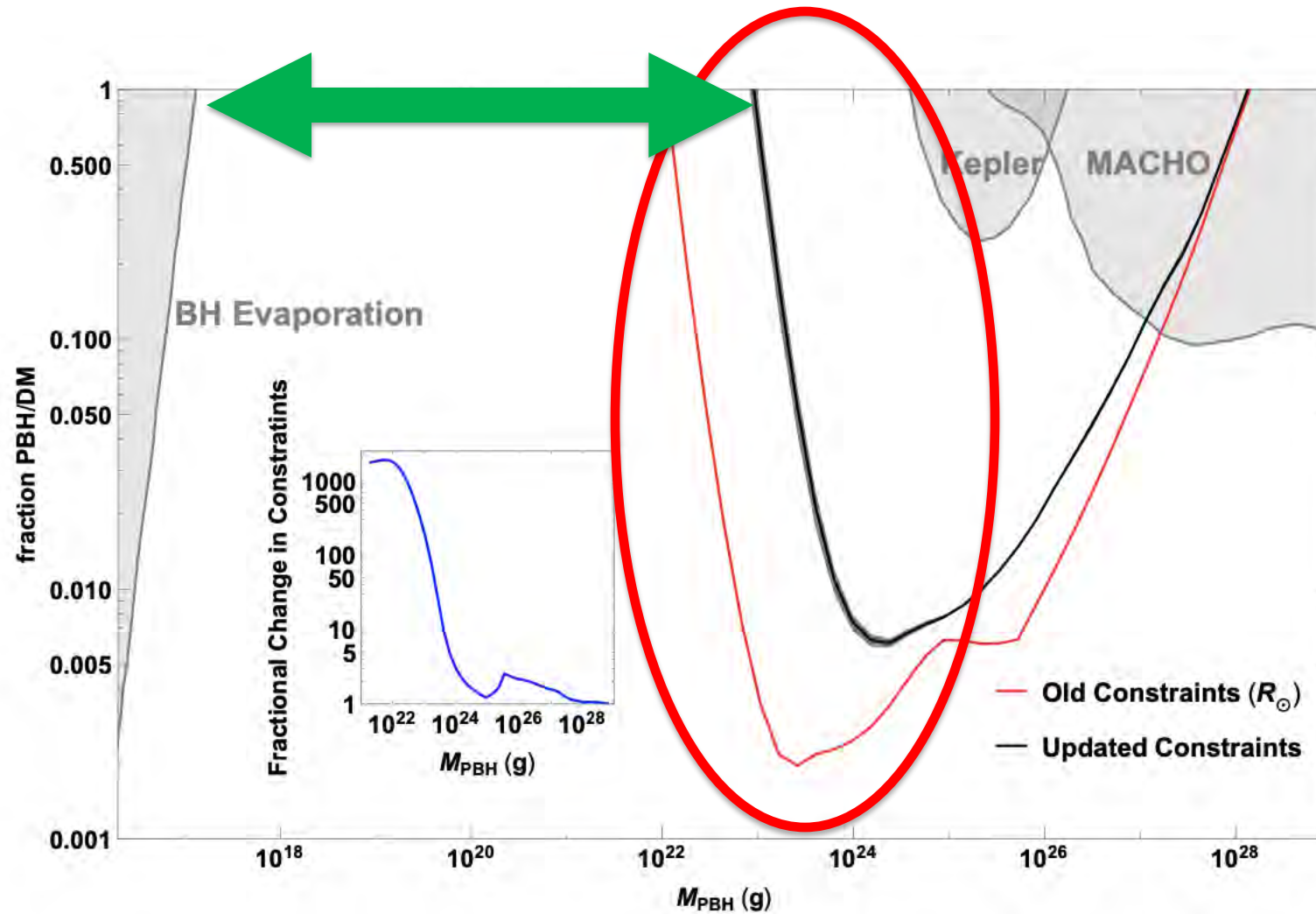
* Profumo, Smyth+ PRD 2020

The bigger the star, the more important finite-source-size effects!

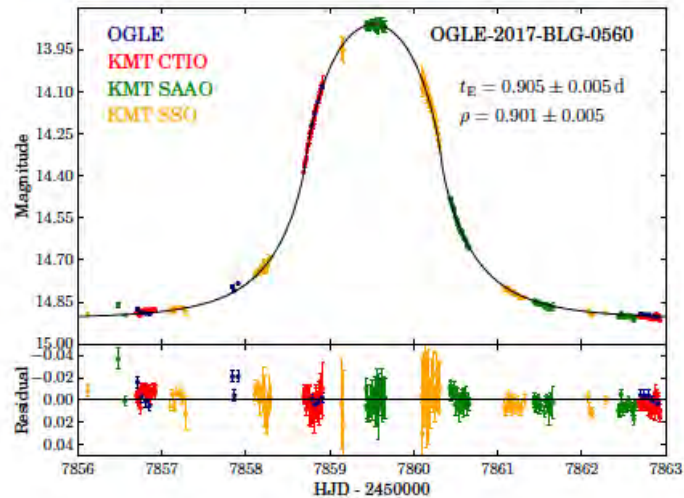
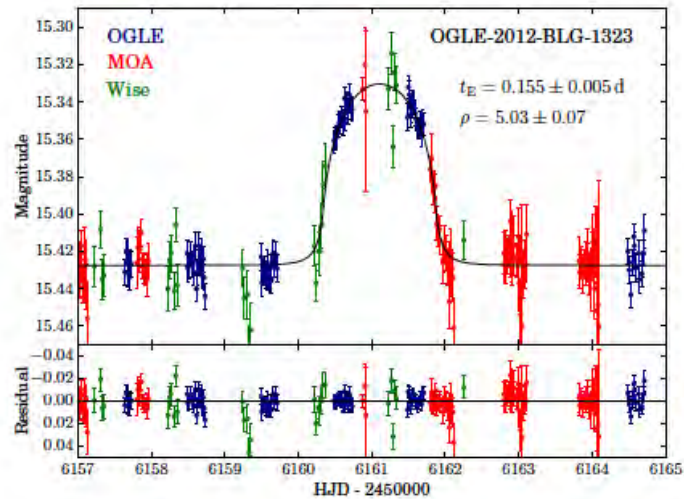




Black holes can be **100% of Dark Matter here!**



Microlensing events have been **detected**... most recently by **OGLE**



...but have we discovered PBH, or
Freely-Floating Planets*?



***Freely-Floating Planet**: “rogue” planet not orbiting a star but, rather, floating in the MW



Nancy Grace Roman Space Telescope (May 2027)

FFPs are an irreducible
background to searches for PBH
...and **VICEVERSA!**

HOW ROGUE WORLDS ARE CREATED

Free-floating planets don't orbit stars the way Earth does. Here are four theories of how these untethered worlds come to be.



PLANETARY ENCOUNTER

When a massive, Jupiter-like planet has a close approach with a much smaller planet, the little guy can get ejected from its orbit.



STAR-LIKE FORMATION

Planets can form independently, alone, from gas and dust in the same stellar nurseries where stars are born.



ENCOUNTER WITH ANOTHER STAR

When two stars approach closely, they can destabilize each others' planets; this may happen often in dense clusters where stars form.



SUPERNOVA EXPLOSION

The intense explosion that comes with the death of a giant star disperses much of the star's mass and gravity. Orbiting planets could be cut loose to wander the cosmos.

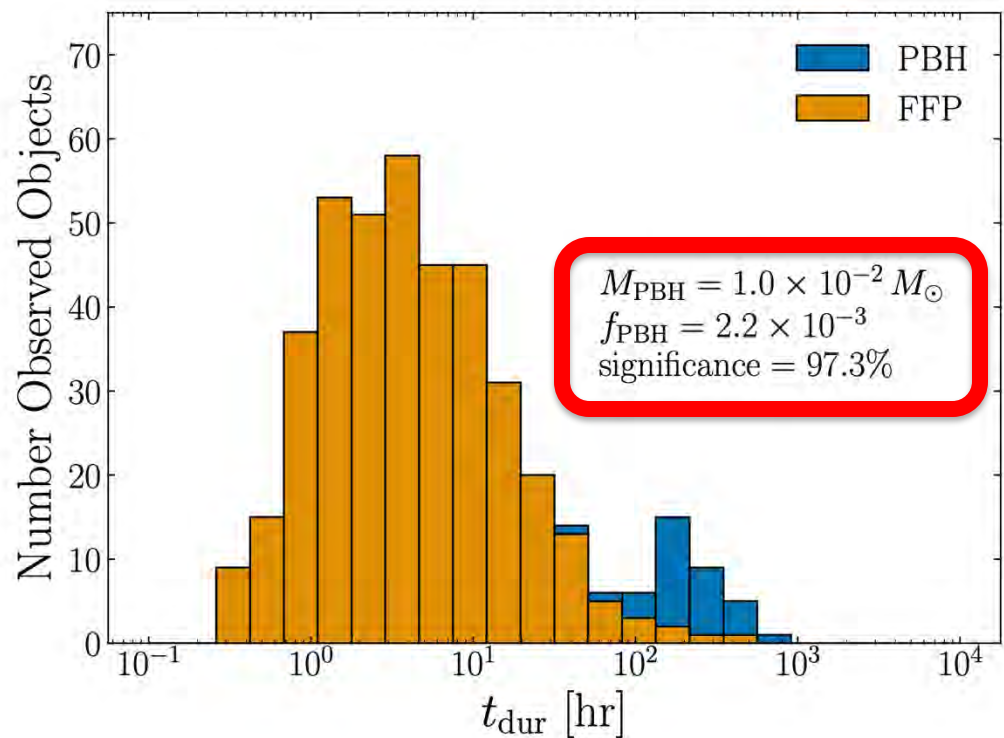
Will **Roman** enable **disentangling** FFP versus **FFP+PBH**?

We studied the **statistics** of anticipated event **duration distribution**

Bonus: we released a code that computes statistics of microlensing events highly efficiently, **LensCalcPy****

Will Roman enable disentangling FFP versus FFP+PBH?

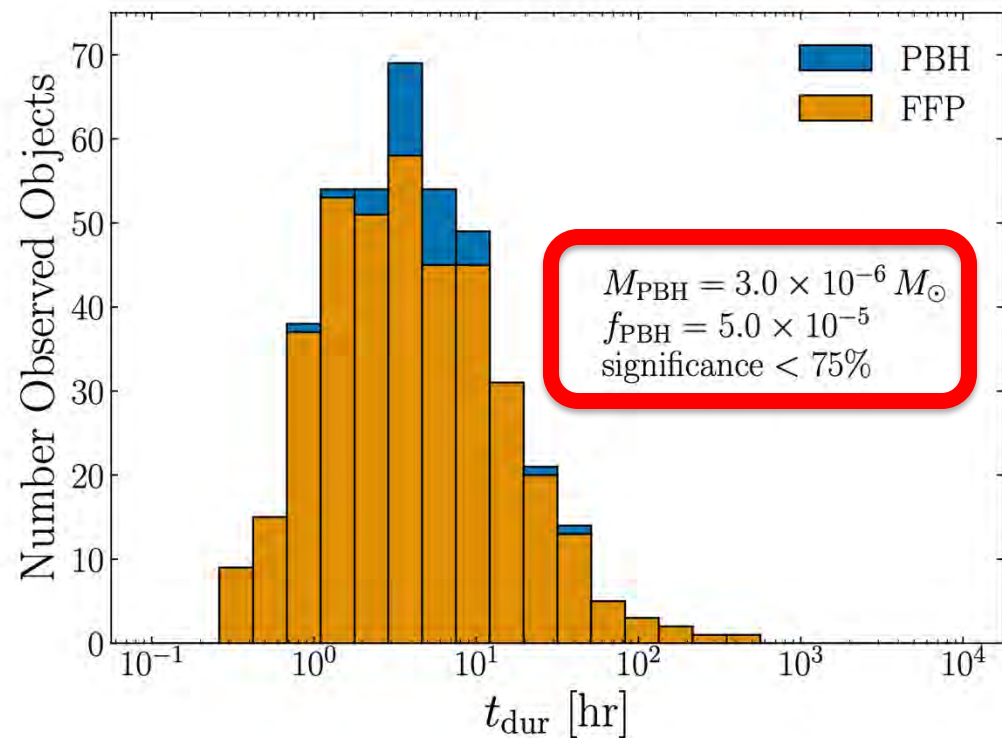
Yes



Null Hypothesis:
Only FFPs

Alternate Hypothesis: FFPs
+ PBHS

No

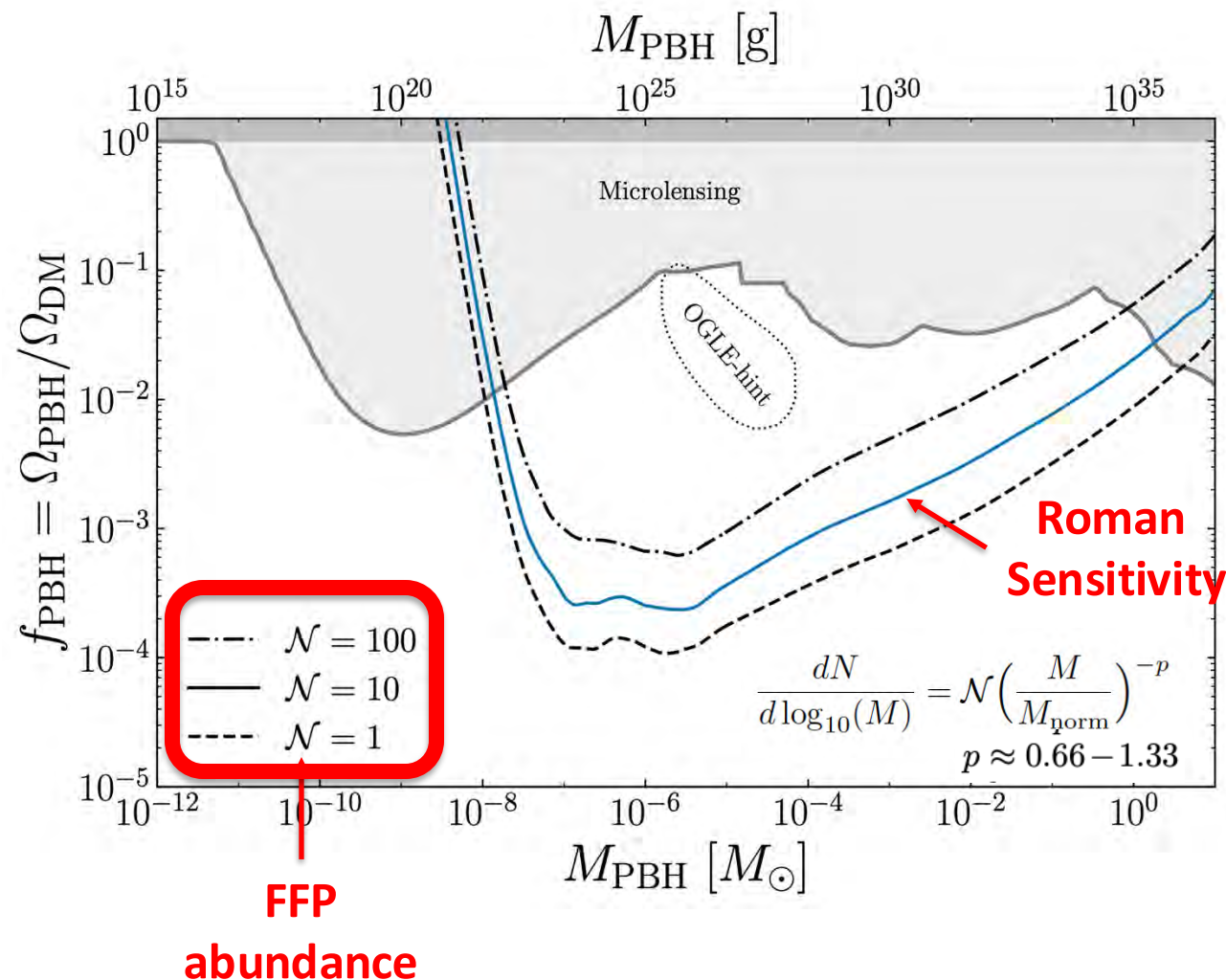


Statistical test: two-sample
Anderson-Darling test

Can we find PBHs **hidden** in **FFP** background?

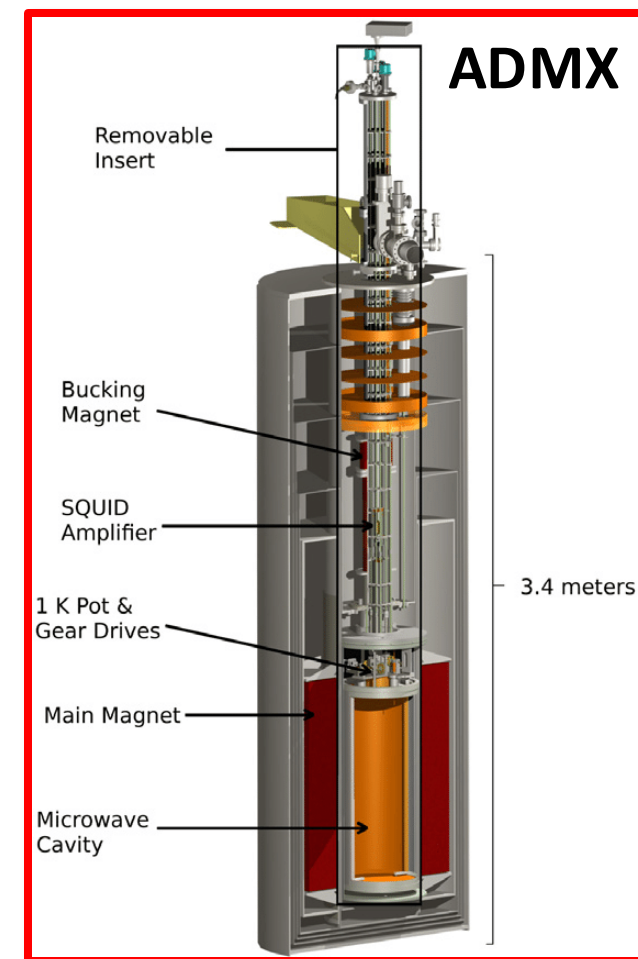
Roman: strongest bounds by far on PBHs as dark matter, even with substantial **FFP background**

Will be able to completely **probe** the preferred OGLE PBH population!



How can we access the “**Asteroid Window**”?

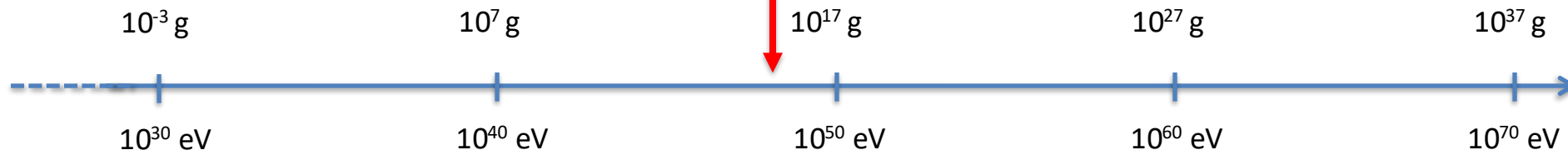
- Production mechanism can source e.g. Stochastic **Gravitational Wave** Background (but model dependent!)
- Lensing with **higher-frequency** observations, e.g. X-ray (but issue of statistics, observation length), or **femtolensing** (but again finite-size source effects!)
- LIGO-style BH-BH **mergers**?*
- Perturbations to **planetary orbits**?**



* Profumo et al, 2410.15400

** Geller, Lehmann et al, 2312.17217

“Exploding” Black Holes (10^{16} g)



A field theory defined on a black-hole background is in a **thermal** state whose temperature at infinity is

$$T = M_P^2 / M_{BH}$$

Black holes radiate (~)like any **black body**, and, as such, shed their mass at a rate

$$\frac{dM}{dt} \propto A(T) T^4 \propto \frac{M^2}{M^4} \propto M^{-2}$$

[Stefan-Boltzmann]

The resulting runaway evaporation process gives a lifetime

“Black Hole Explosion”*

$$\tau \approx 407 \left(\frac{f(M)}{15.35} \right)^{-1} M_{10}^3 \text{ s.}$$

[10¹⁰ grams units]

Black holes formed in the early universe, with a mass **$M_U \sim 5 \times 10^{14}$ grams, $T \sim 100$ MeV** are exploding **today****

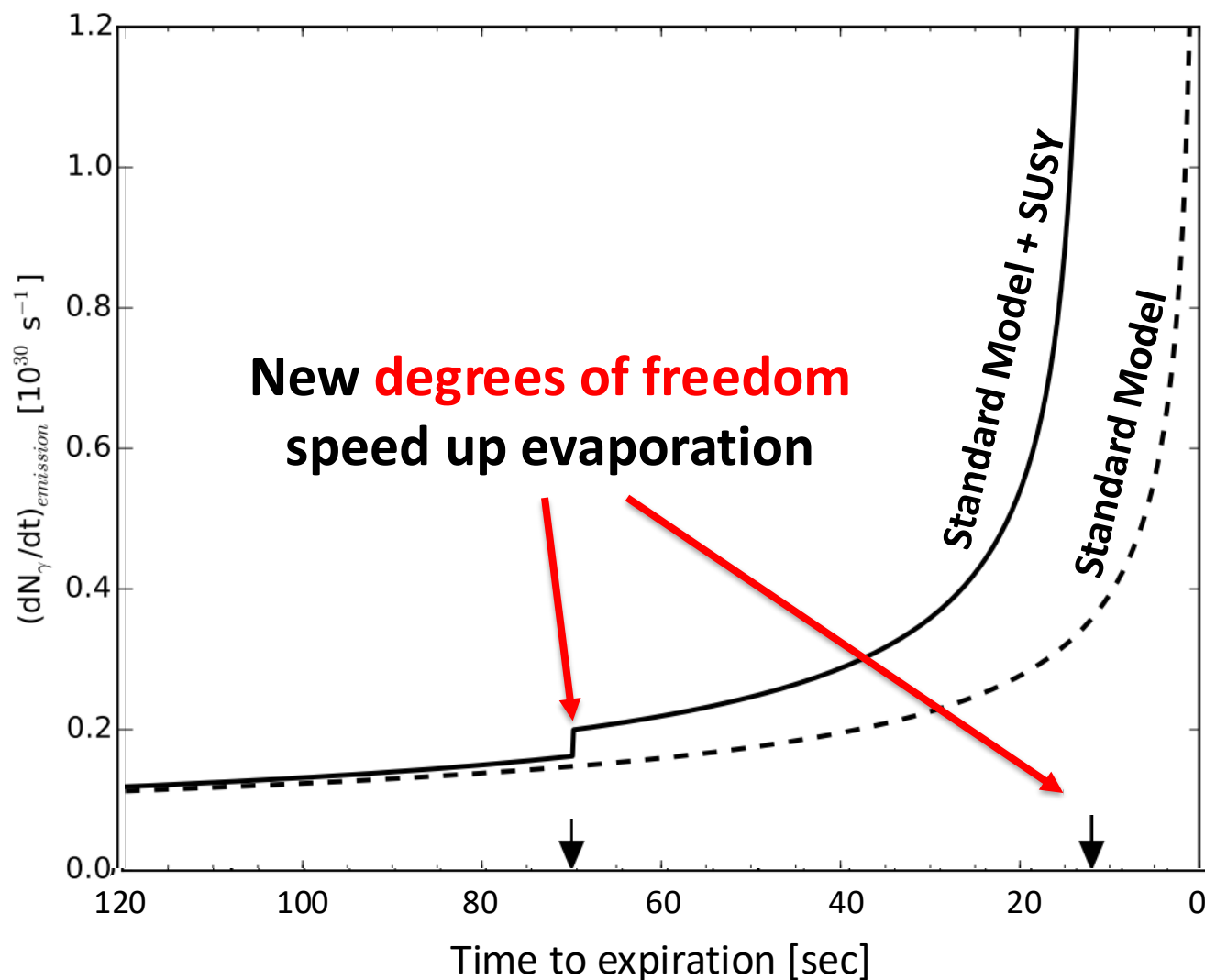
*Hawking, 1974

****or lighter, and formed later!**

The (universal*) Black Hole Light-curve

Universal **Luminosity**
→ perfect
“standard candle”

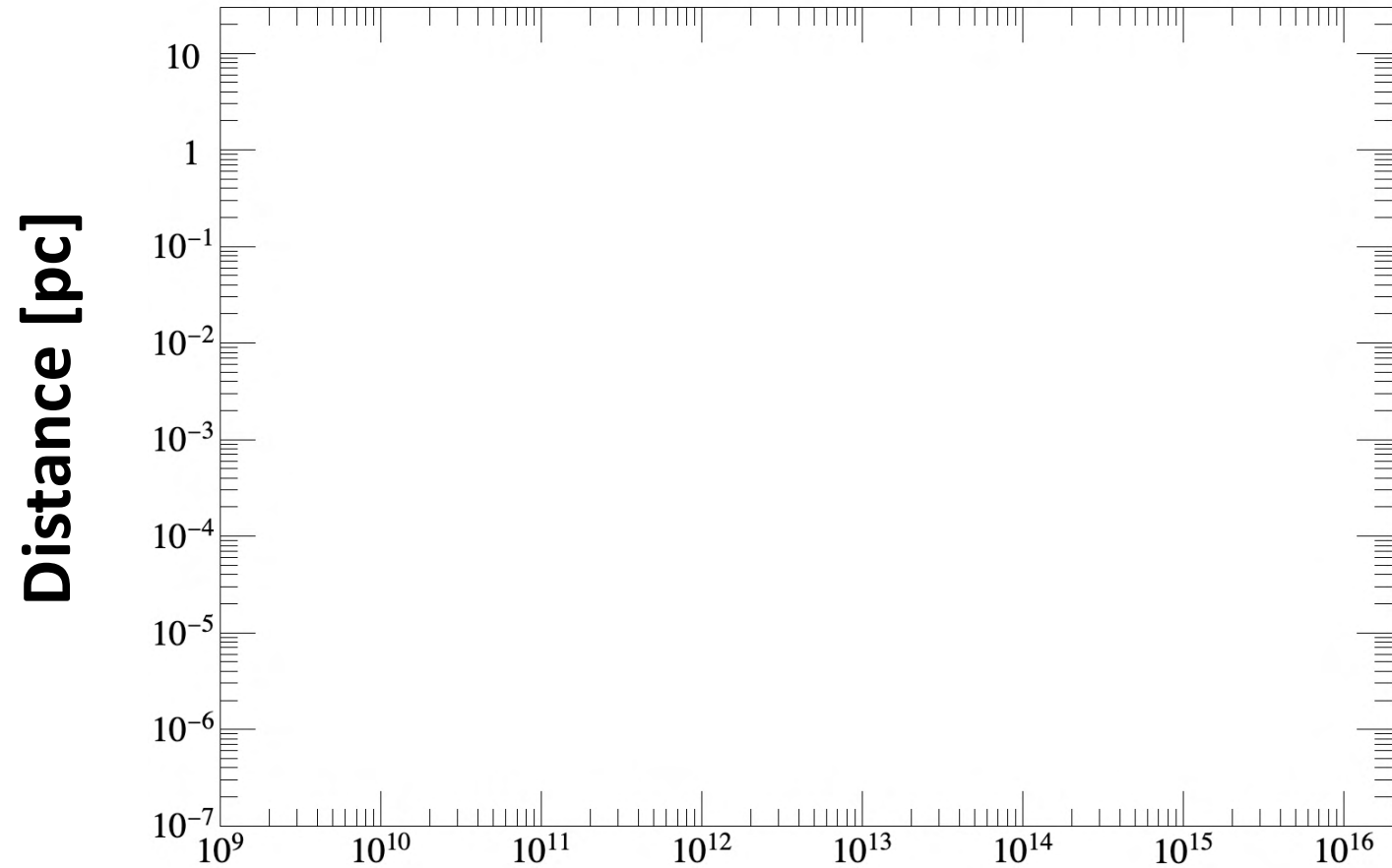
“Backwards”
Gamma-Ray Burst



No **afterglow**
(BHs don't form a
photosphere or
chromosphere)

*for a given spin/charge, and evaporating d.o.f

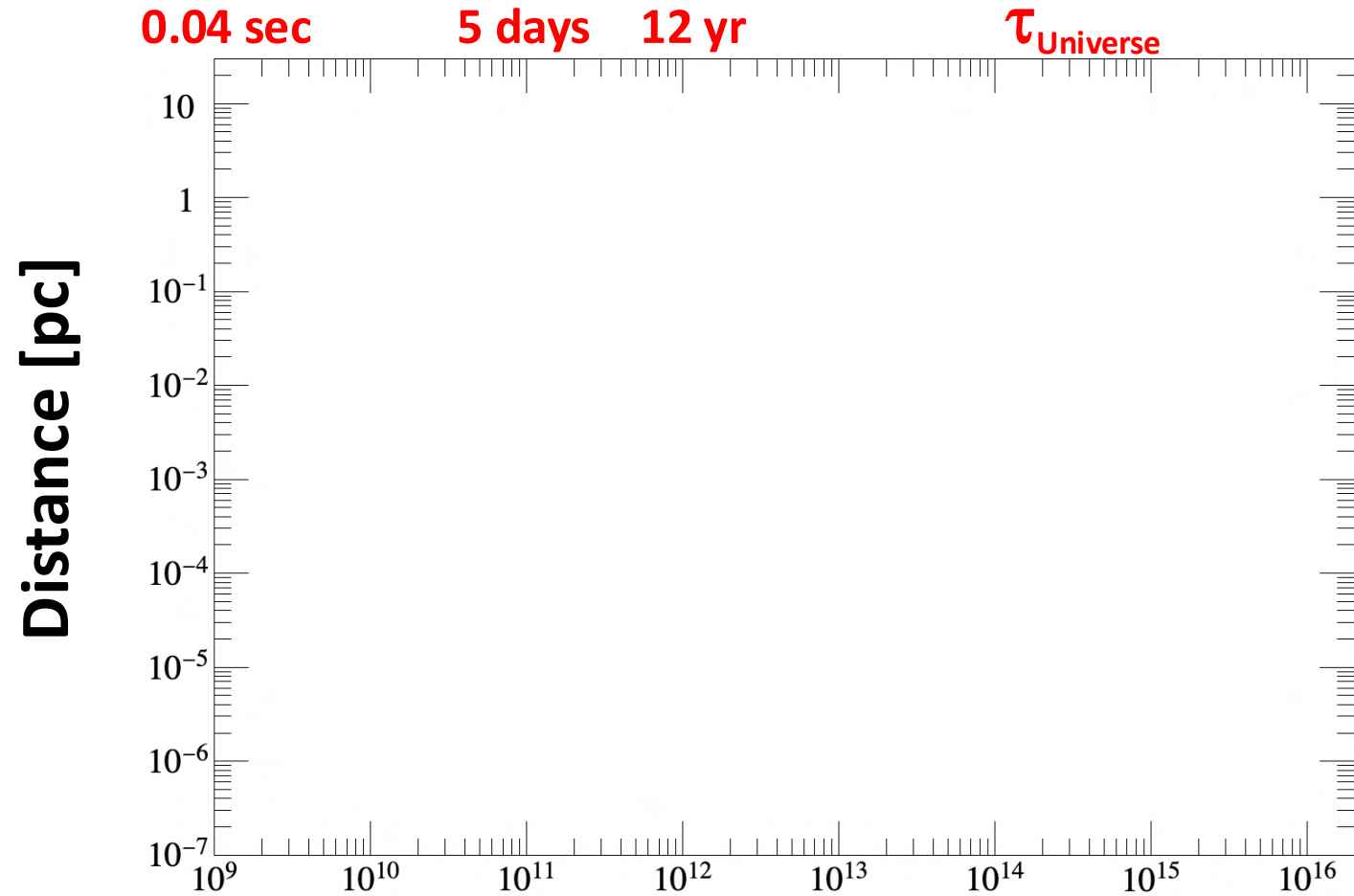
How can we **search** for PBH explosions **today**?



Mass/Lifetime [today, at detection, grams]

Remaining lifetime

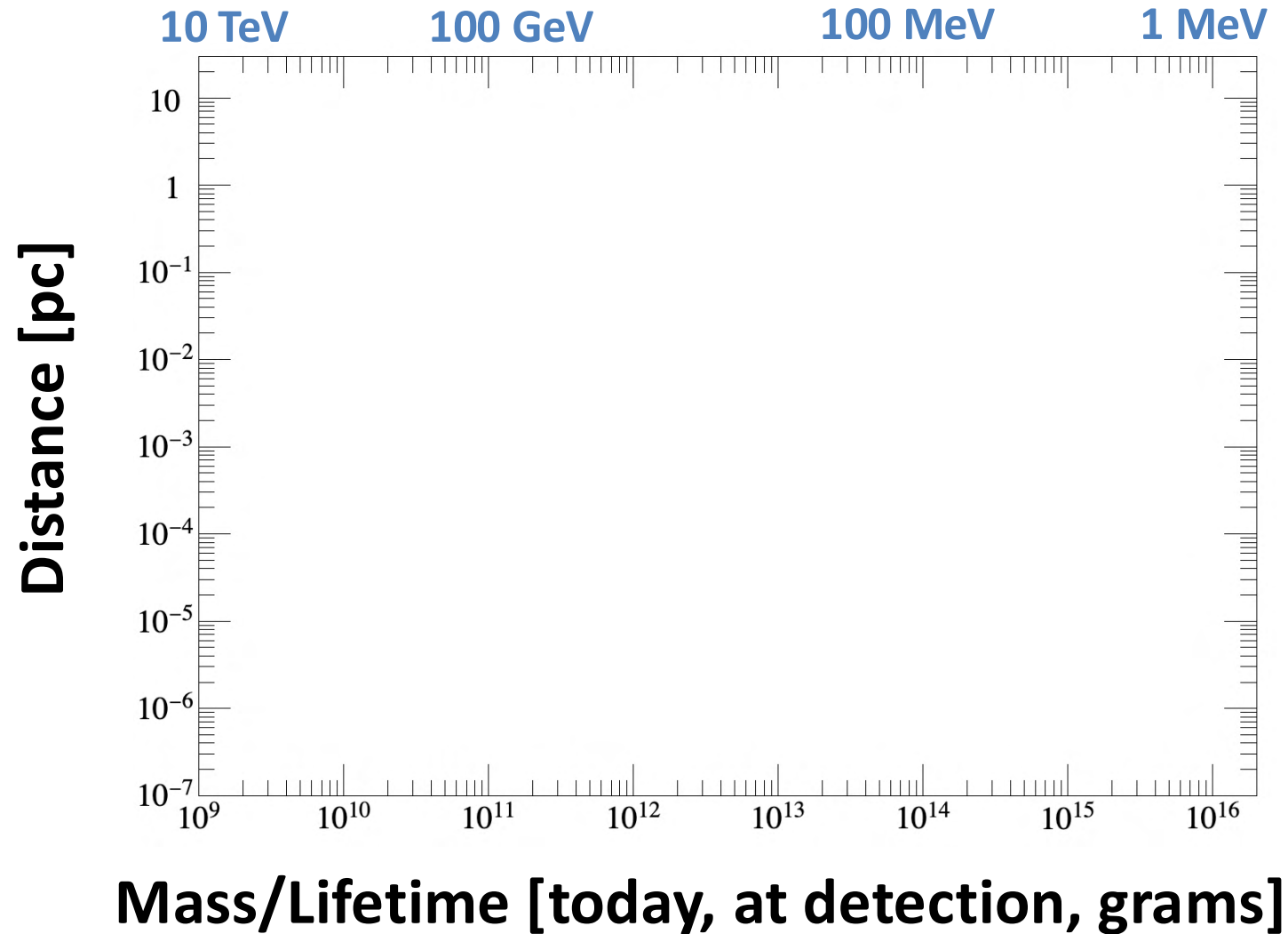
$$\tau \simeq 400 \text{ sec} \left(\frac{M_{\text{BH}}}{10^{10} \text{ g}} \right)^3$$



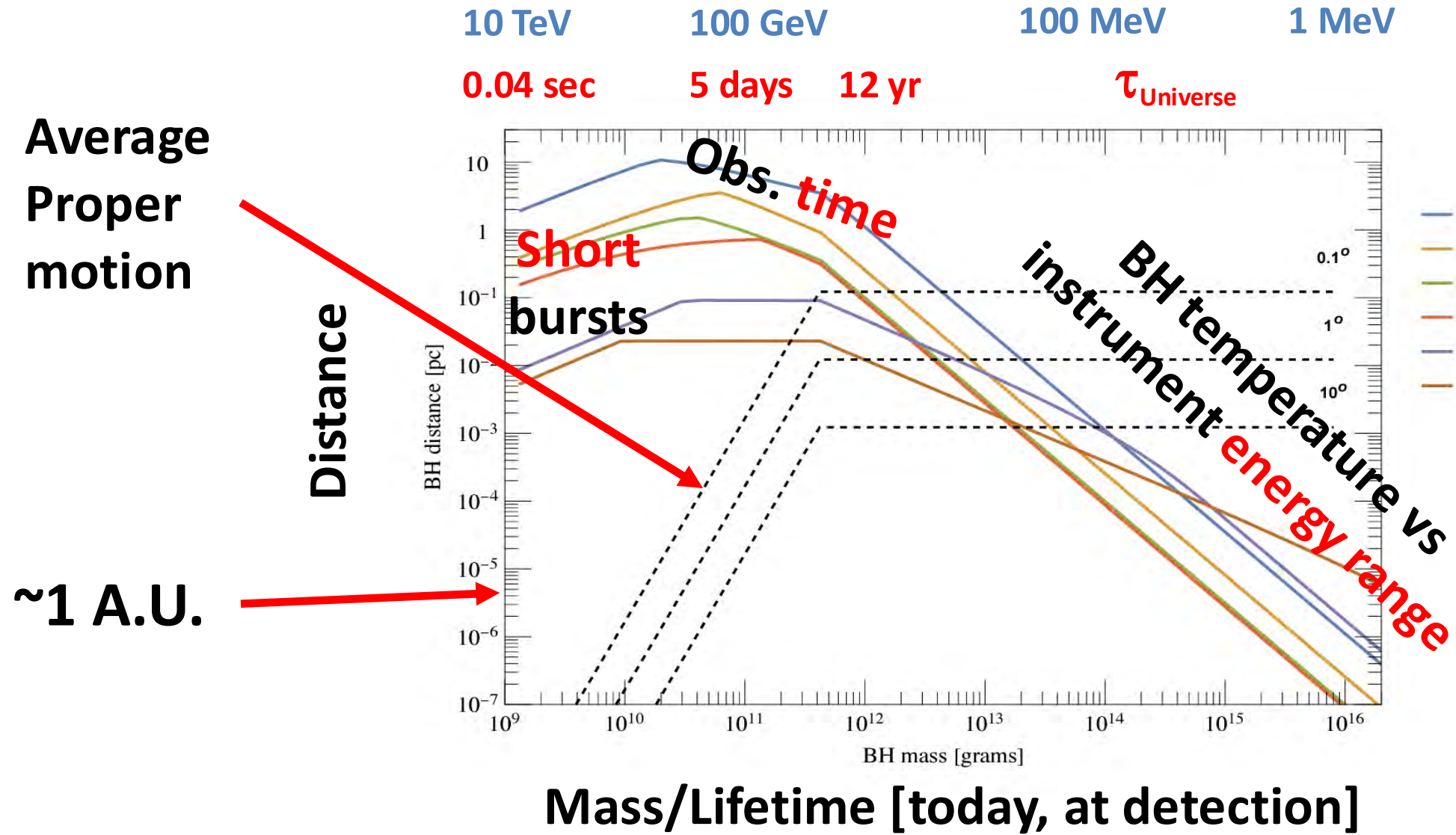
Mass/Lifetime [today, at detection, grams]

Temperature

$$T \simeq 1 \text{ TeV} \left(\frac{10^{10} \text{ g}}{M_{\text{BH}}} \right)$$



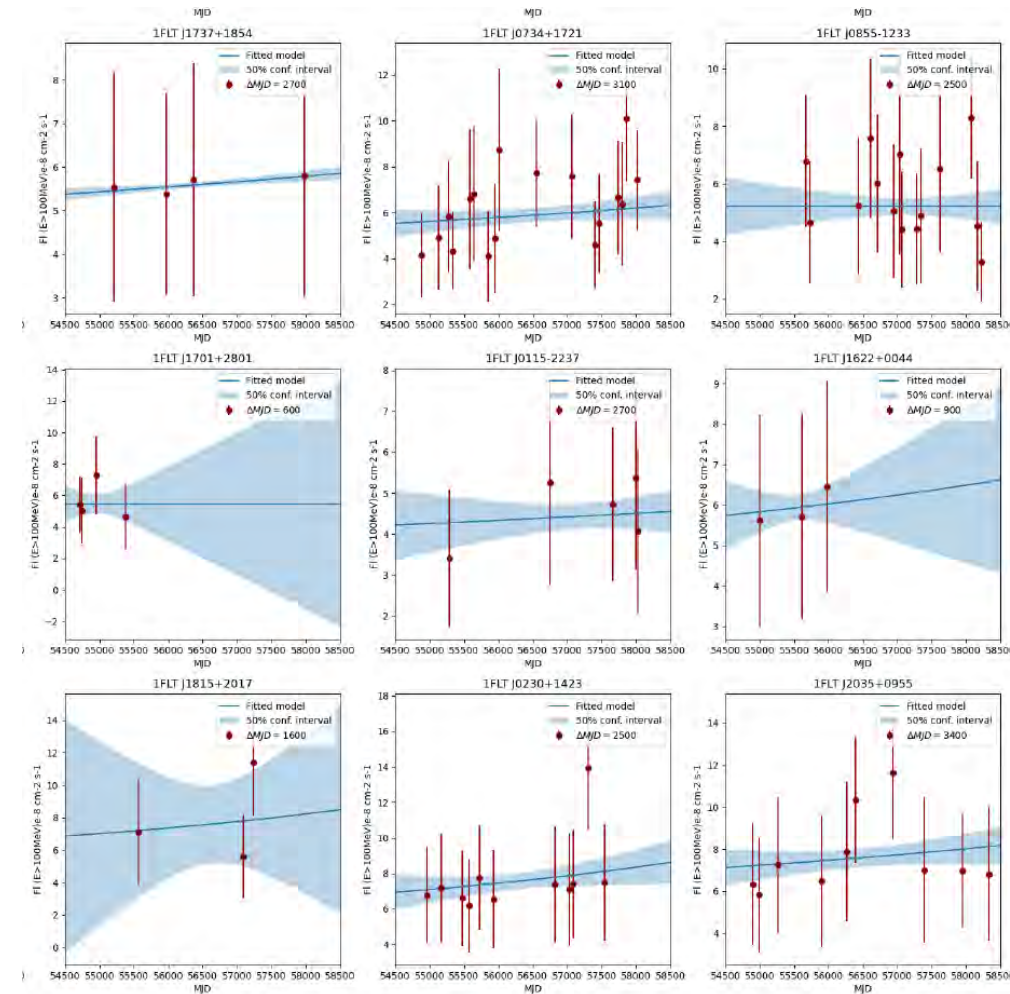
Gamma-Ray Observatory



*Profumo, Boluna, Ble, Hennings, 2023 (JCAP)

LAT and GBM Gamma-ray burst catalog

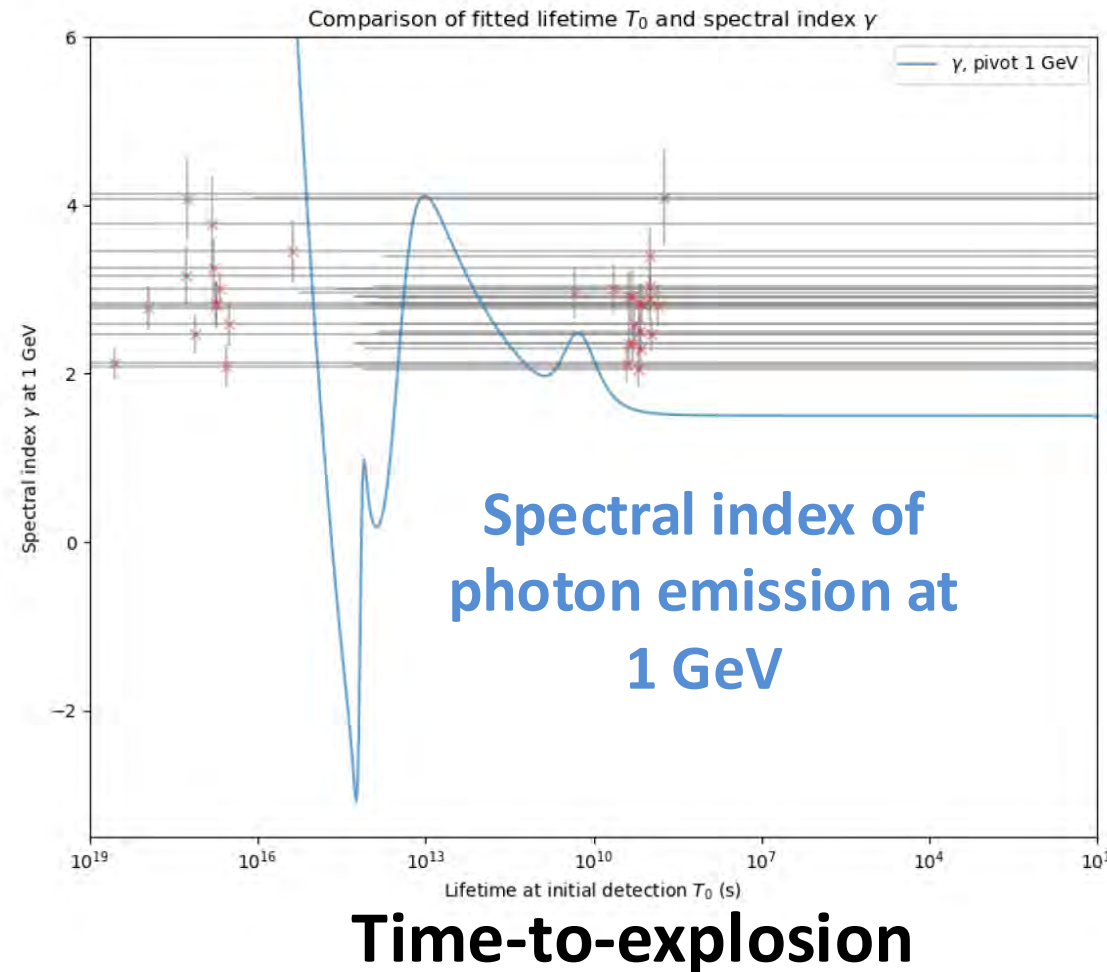
(1) LAT **variable** sources



time

LAT and GBM Gamma-ray burst catalog

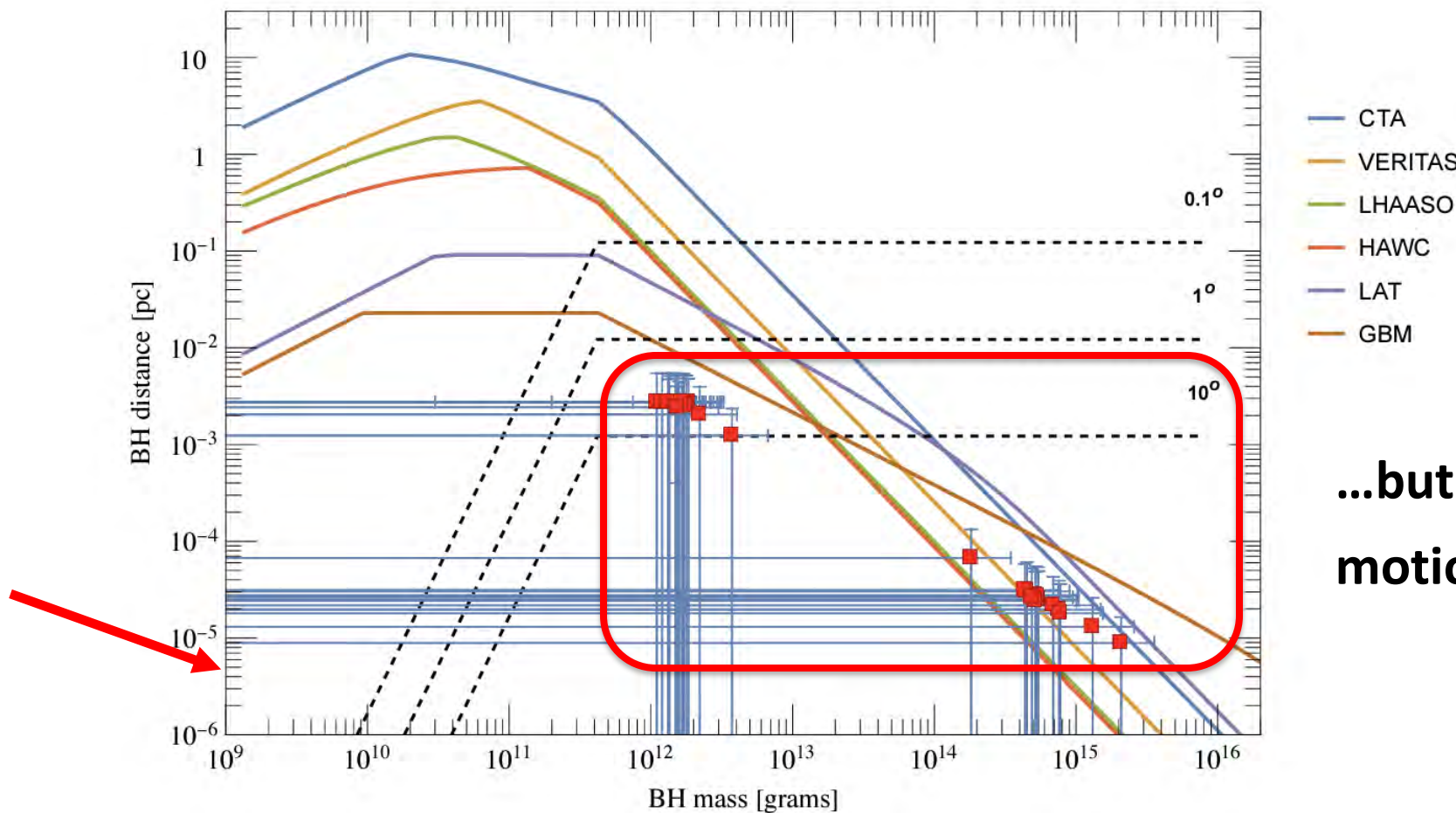
(2) LAT variable sources: **spectral** fit



LAT and GBM Gamma-ray burst catalog

(3) LAT variable sources: distance-age fit

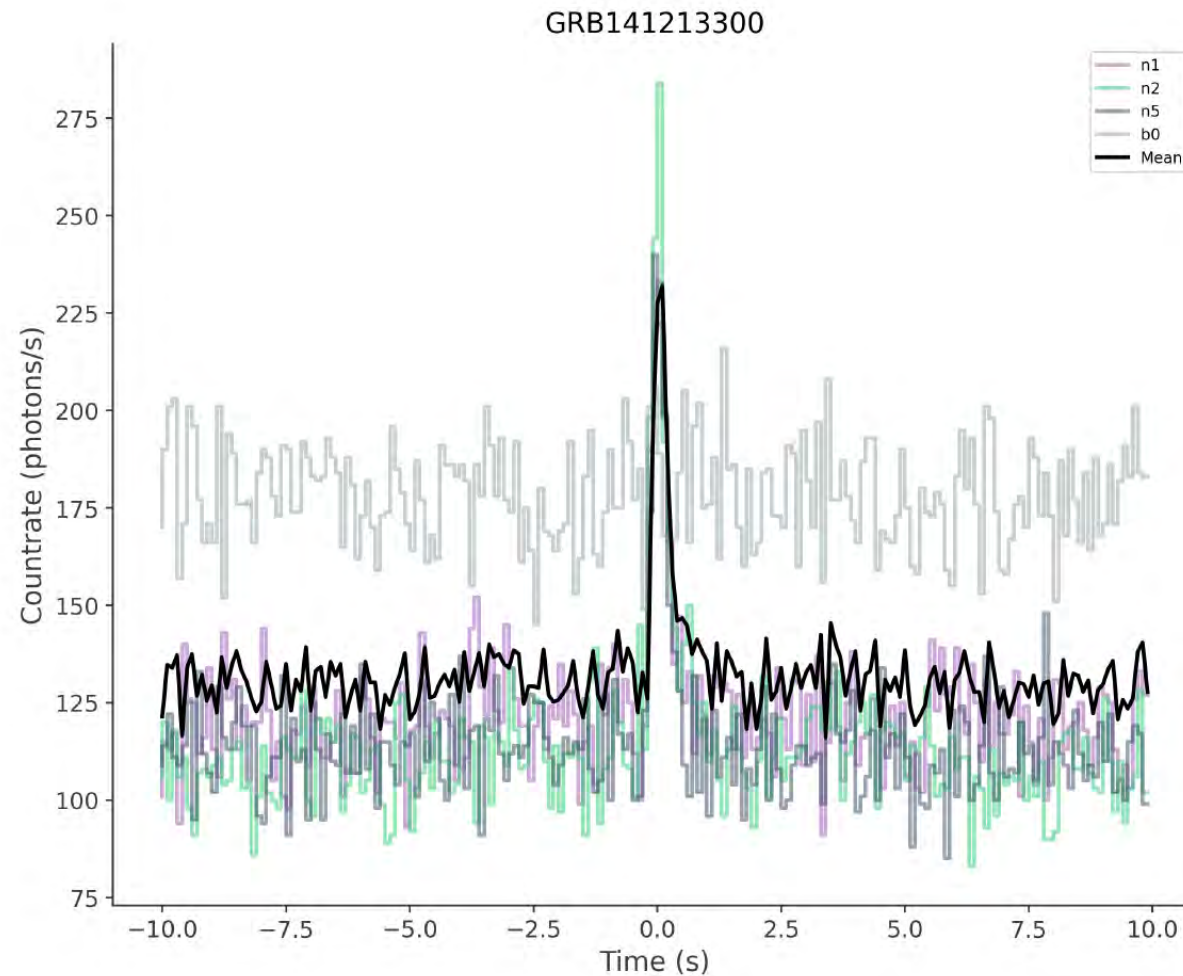
~1 A.U.



...but no proper
motion detected 🙄

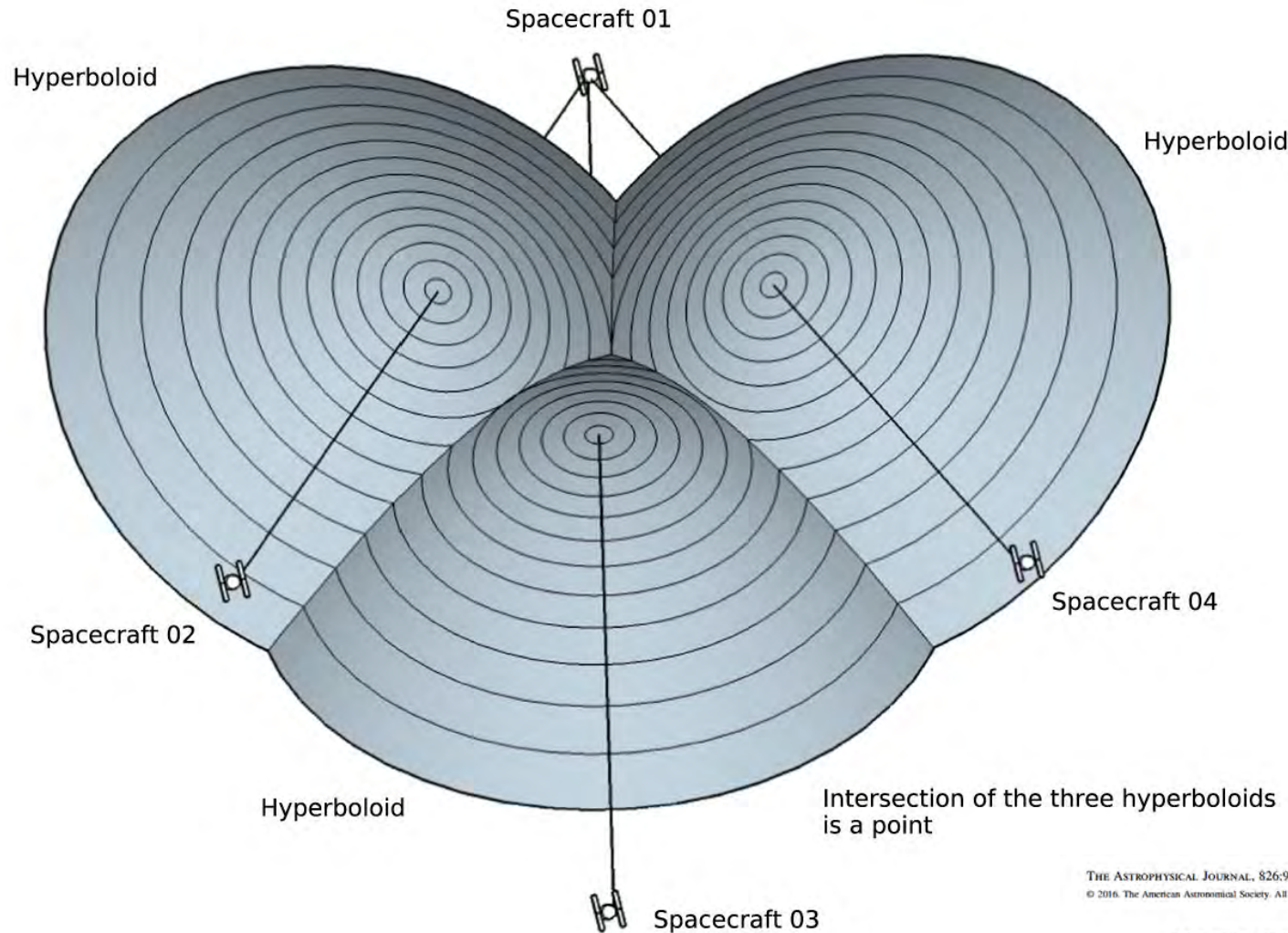
LAT and GBM Gamma-ray burst catalog

(4) Short duration: **light curve**



time

(5) Interplanetary GRB monitors network (IPN)



Timing
resolution
distinguishes
cosmological vs
local events!

THE ASTROPHYSICAL JOURNAL, 826:98 (14pp), 2016 July 20
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doi:10.3847/0004-637X/826/1/98
CrossMark

INVESTIGATION OF PRIMORDIAL BLACK HOLE BURSTS USING INTERPLANETARY
NETWORK GAMMA-RAY BURSTS

T. N. UKWATTA¹, K. HURLEY², J. H. MACGIBBON^{3,17}, D. S. SVINKIN⁴, R. L. APTEKAR⁴, S. V. GOLENETSKI⁴, D. D. FREDERIKS⁴,
M. D. BARNARD⁴, I. GARDNER⁵, W. HARRISON⁶, A. G. KRAUSS⁷, A. PIER⁸, J. R. KRAUSS⁸, N. ZUCKER⁸, V. G. ZUCKER⁹

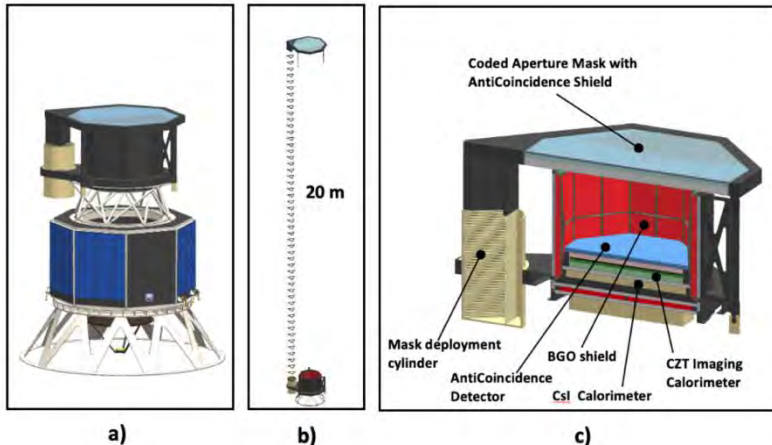
Evaporation products (gamma rays, cosmic-ray positrons) are **detectable**, constraining the **fraction** of light PBH that can be the DM



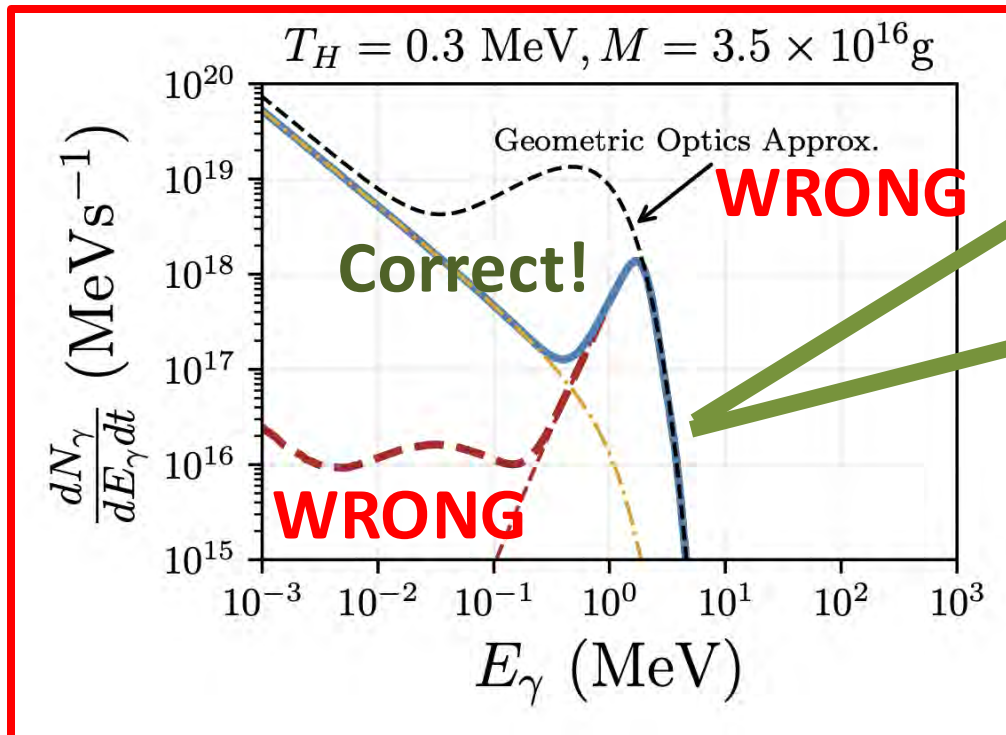
Issue: relevant temperature $T \sim T_{QCD}$

- Chiral Perturbation Theory
- Vector Meson Dominance
- Hadronic Form Factors

Bonus: can recycle for **MeV Dark Matter** pheno!



Evaporation products (gamma rays, cosmic-ray positrons) are **detectable**, constraining the **fraction** of light PBH that can be the DM



Direct Detection of Hawking Radiation from Asteroid-Mass Primordial Black Holes

Adam Coogan,^{1,*} Logan Morrison,^{2,†} and Stefano Profumo^{2,‡}

¹GRAPPA, Institute of Physics, University of Amsterdam, 1098 XH Amsterdam, The Netherlands

²Department of Physics, University of California, Santa Cruz, CA 95064, USA

(Dated: October 13, 2020)

Light, asteroid-mass primordial black holes, with lifetimes in the range between hundreds to several millions times the age of the universe, are well-motivated candidates for the cosmological

- Home
- Description
- Manual
- Download
- Contact

BlackHawk

By **Alexandre Arbey** and **Jérémy Auffinger**

Calculation of the Hawking evaporation spectra of any black hole distribution

BlackHawk is a public C program for calculating the Hawking evaporation spectra of any black hole distribution. This program enables the users to compute the primary and secondary spectra of stable or long-lived particles generated by Hawking radiation of the distribution of black holes, and to study their evolution in time.

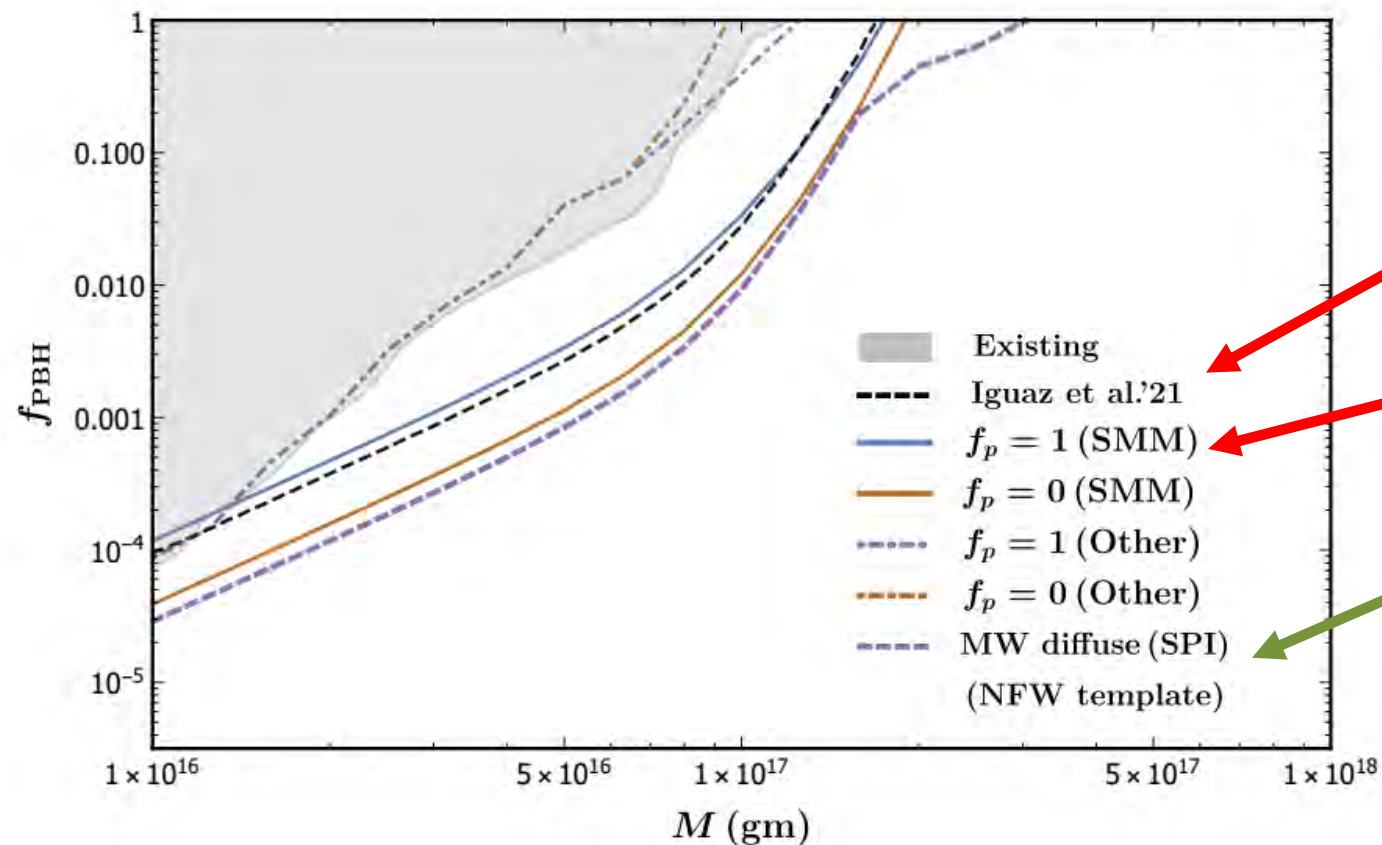
If you use BlackHawk to publish a paper, please cite:

- A. Arbey and J. Auffinger, *Eur. Phys. J. C* **79** (2019) 693, [arXiv:1905.04268 \[gr-qc\]](#)
- A. Arbey and J. Auffinger, *Eur. Phys. J. C* **81** (2021) 910, [arXiv:2108.02737 \[gr-qc\]](#)

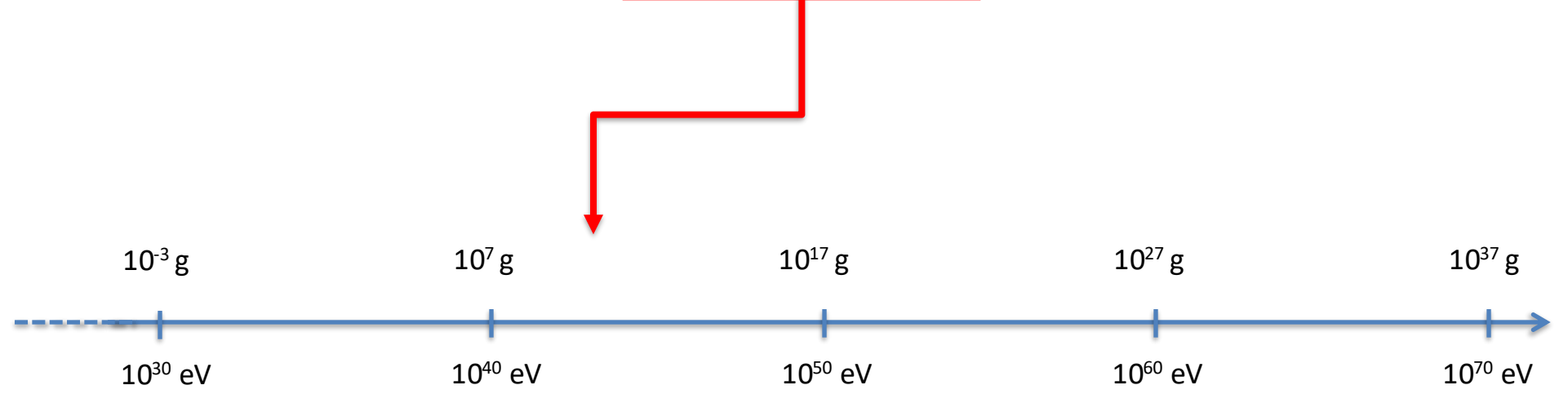
If you use the hadronized spectra of BlackHawk, we also advise that you cite the corresponding particle physics code:

- PYTHIA spectra (hadronization_choice = 0 or 2): T. Sjöstrand *et al.*, *Comput. Phys. Commun.* **191** (2015) 159-177, [arXiv:1410.3012 \[hep-ph\]](#)
- HERWIG spectra (hadronization_choice = 1): J. Bellm *et al.*, *Eur. Phys. J. C* **76** (2016) 4, 196, [arXiv:1512.01178 \[hep-ph\]](#)
- HAZMA spectra (hadronization_choice = 3): A. Coogan, L. Morrison, S. Profumo, *JCAP* **01** (2020) 056, [arXiv:1907.11846 \[hep-ph\]](#)
- HDMSpectra spectra (hadronization_choice = 4): C. W. Bauer, N. L. Rodd, B. R. Webber, *JHEP* **06** (2021) 121, [arXiv:2007.15001 \[hep-ph\]](#)

Strongest constraints to date:
MW diffuse gamma-ray emission from Integral-SPI,
including the 511 keV line



**Whale-Mass
“Exploded”
Black Holes**



Exploded Black Holes

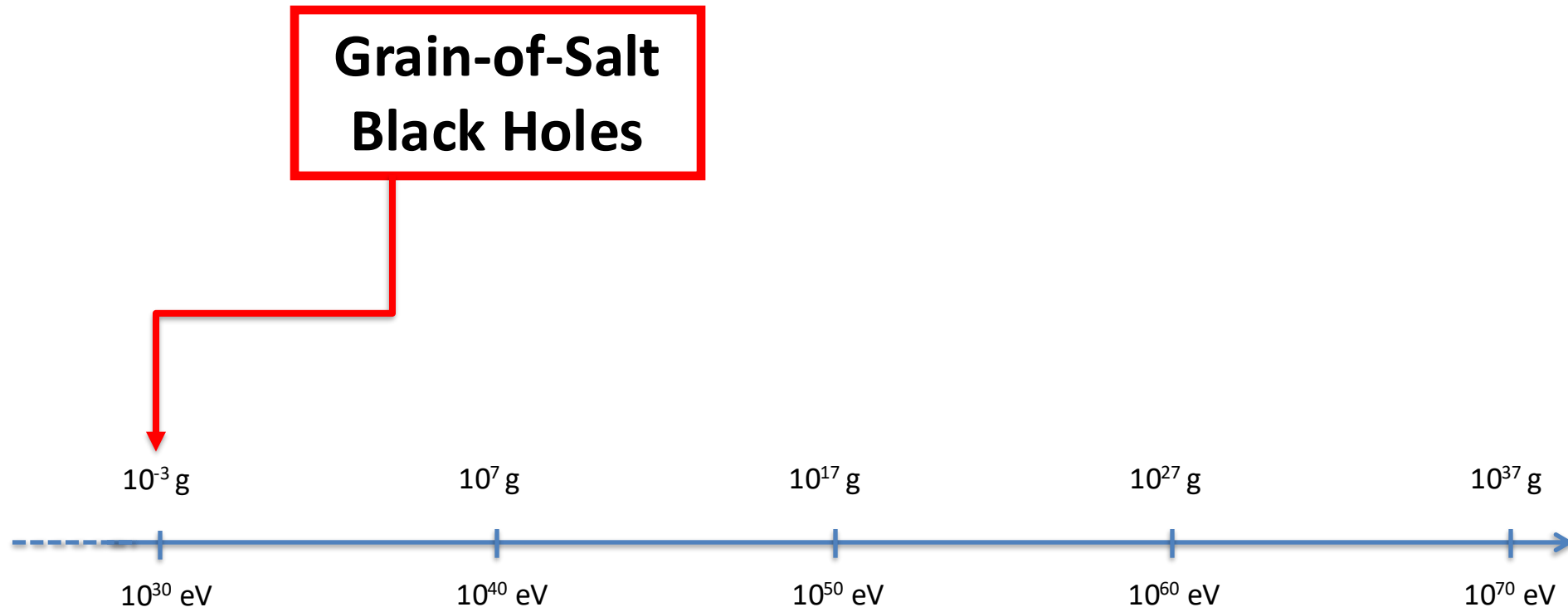
- Can have sources both the **dark matter** and the **baryon asymmetry** via **evaporation**
[Morrison, Profumo, Yu, 2022]

- Can have sourced both the dark matter and the baryon asymmetry via coupling of **Kretschmann scalar to B-L current**
[Profumo, Smyth, Santos-Olmsted, 2023]

$$\mathcal{S} \supset \frac{1}{M_*^4} \partial_\alpha (\mathcal{R}_{\mu\nu\rho\sigma} \mathcal{R}^{\mu\nu\rho\sigma}) J_{\text{B-L}}^\alpha$$

- Can seed high-frequency **gravitational waves**, especially if...
 - ✓ ...**modified cosmology** shifts the spectrum to reasonably low frequencies
 - ✓ ...**extradimensions** lower the Planck scale and the frequency of GW at the end of evaporation
[Ireland, Profumo, Sharnhorst, 2023, 2024]





If **evaporation stops** around the Planck scale, the relic PBHs can acquire a significant **relic** electric **charge**

(under simple **assumptions**) $P(Q) \sim \exp(-4\pi\alpha(Q/e)^2)$
the relic charge is
approximately **Gaussian*** $(8\pi\alpha)^{-1/2} \approx 2.34$

If (1) **evaporation stops** around the Planck scale (because of **extremality** (RN BH) or because of **quantum gravity**)
AND (2) if Schwinger-pair discharge also **stops**
we are **left with** a population of **charged, Planck-scale relics!**

* Page, 1977

** Lehmann, Johnson, Profumo and Schwemberger, 1906.06348 (JCAP10(2019)046)

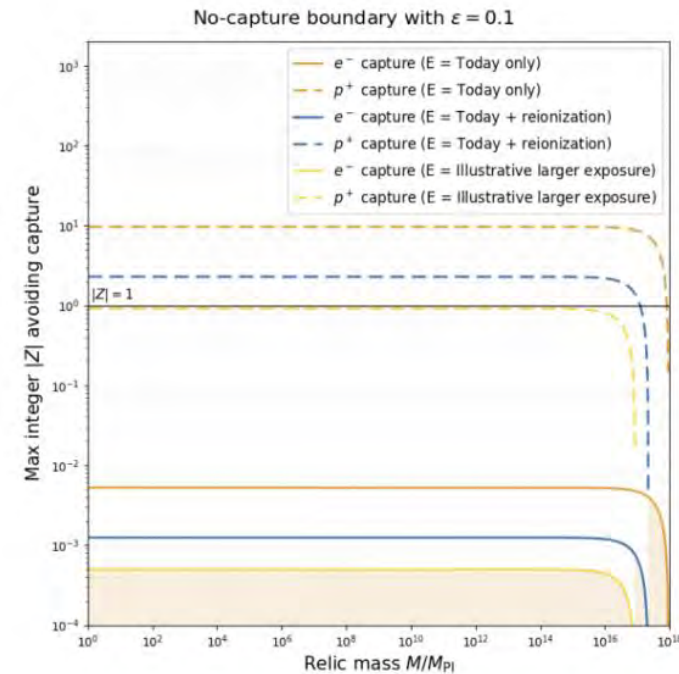
On the (Im)possibility of Electrically Charged Planck Relics

Stefano Profumo^a

^aDepartment of Physics and Santa Cruz Institute for Particle Physics (SCIPP),
University of California, Santa Cruz, CA 95064, USA

E-mail: profumo@ucsc.edu

Abstract. I revisit whether black-hole remnants, from sub-Planckian compact objects to Planck relics and up to (super)massive black holes, can preserve Standard-Model (SM) electric charge. Two exterior-field mechanisms—Coulomb-focused capture from ambient media and QED Schwinger pair production—robustly neutralize such objects across cosmic history. I



Stefano Profumo
University of California
Santa Cruz



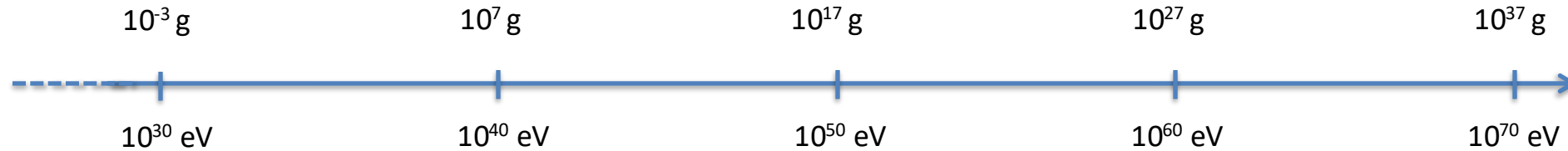
If (1) **evaporation stops** around the Planck scale (because of **extremality** (RN BH) or because of **quantum gravity**)
AND (2) if Schwinger-pair discharge also **stops**
we are **left with** a population of **charged, Planck-scale relics!**

...or are we?

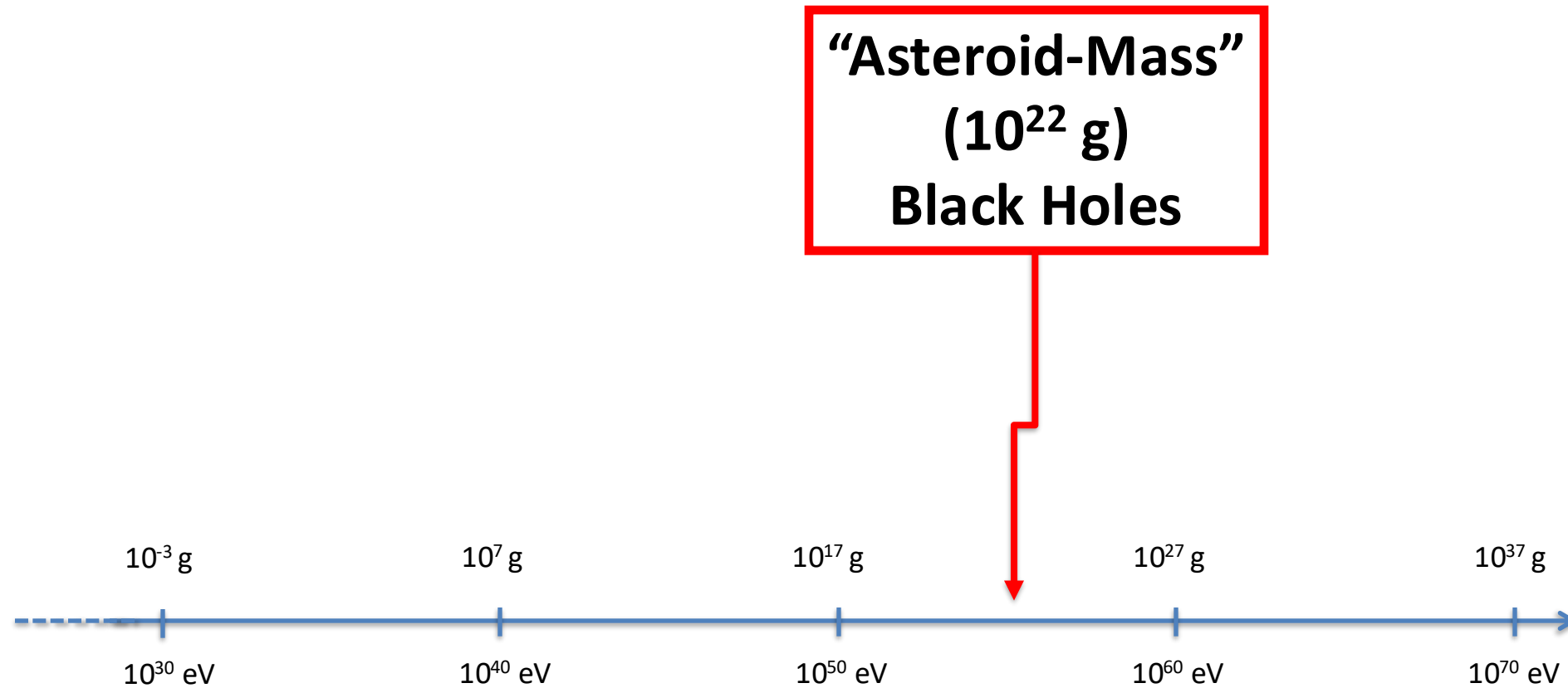


*** Lehmann, Johnson, Profumo and Schwemberger, 1906.06348 (JCAP10(2019)046)**

**“Stellar-Mass”
(10^{35} g)
Black Holes**

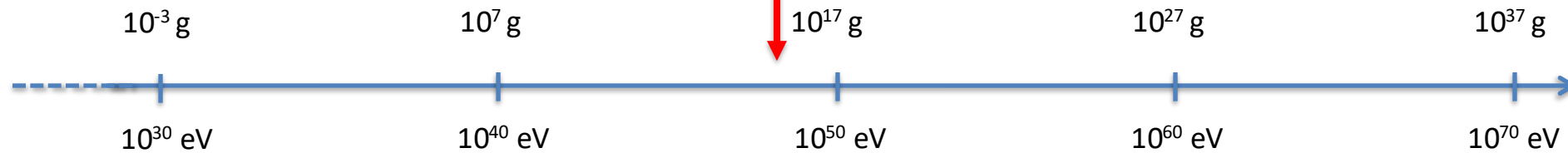


- ✓ Sub-TOV **goldilocks** events! (beware of **impostors**!)
- ✓ Searches **ongoing**, perhaps already **there**!



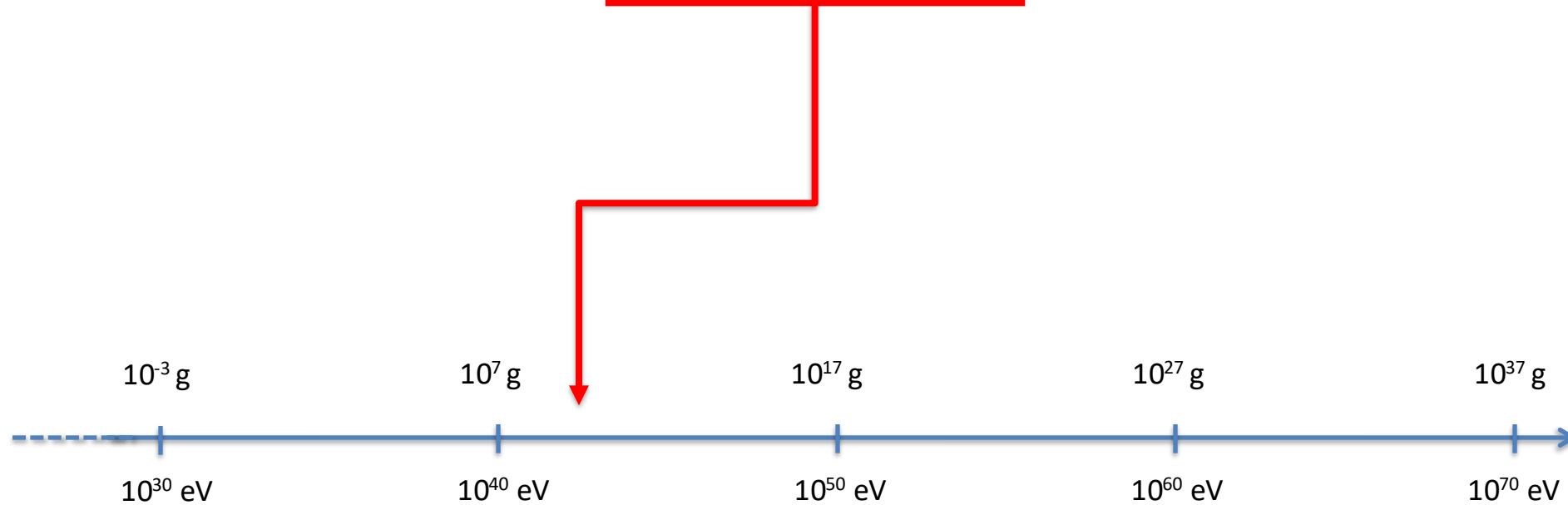
- ✓ Can be 100% of the **Dark Matter**!
- ✓ Microlensing **trickier** than previously thought (**finite size** source)!
- ✓ If detected, how do we distinguish PBH from **rogue planets**? Statistics!

**“Everest-Mass”
(10^{16} g)
“Exploding” Black Holes**

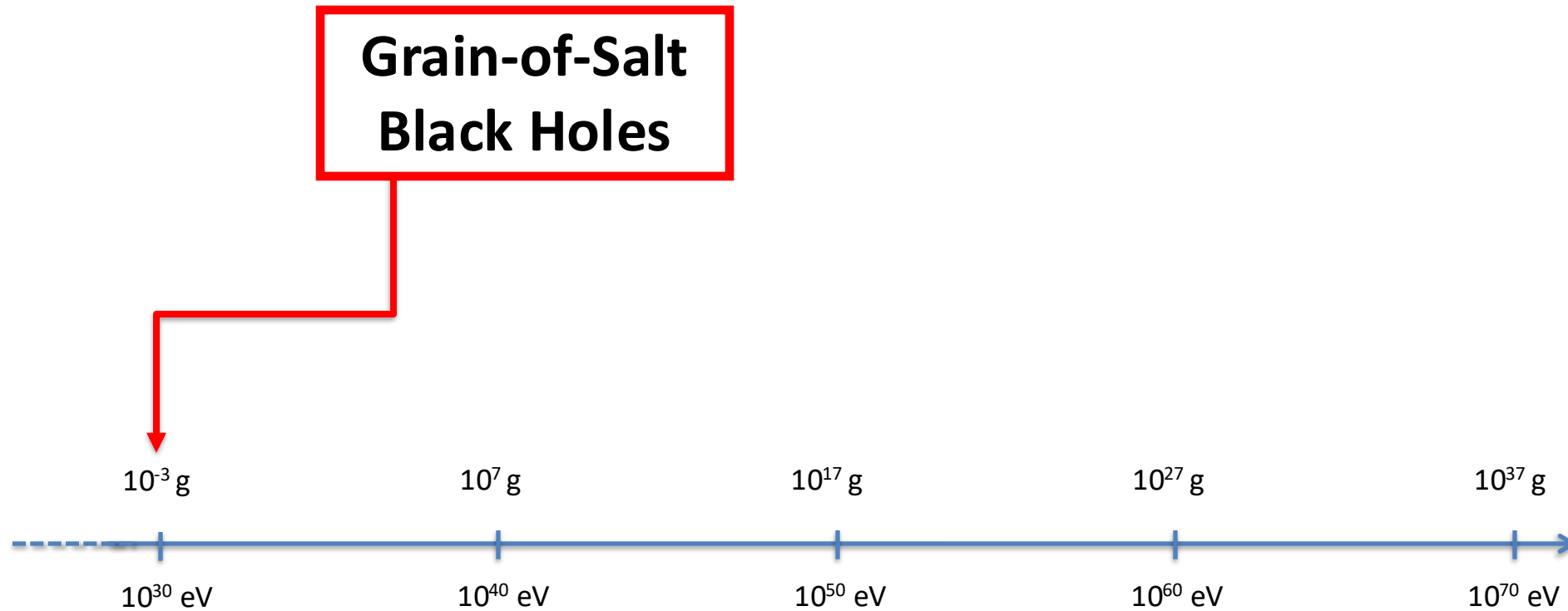


- ✓ BH explosions as “**reverse GRBs**”
- ✓ A window on **Dark Sectors**
- ✓ Best constraints: **diffuse** gamma-rays

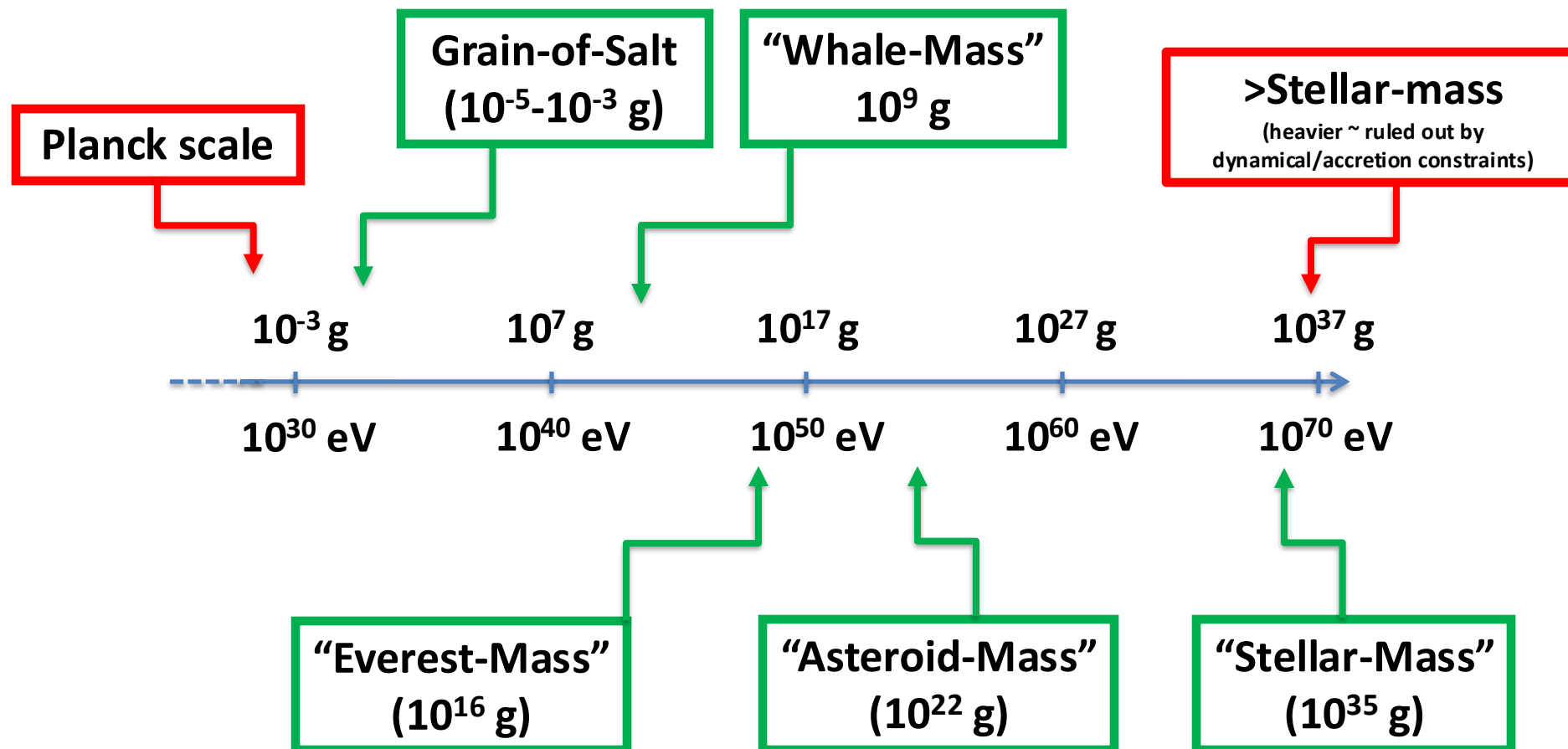
“Whale-Mass” Black Holes



✓ Decays can produce **DM, BAU, Gravitational Waves**



- ✓ Likely (partly) **charged**
- ✓ Detectable! ...best with **Paleo-Detectors**



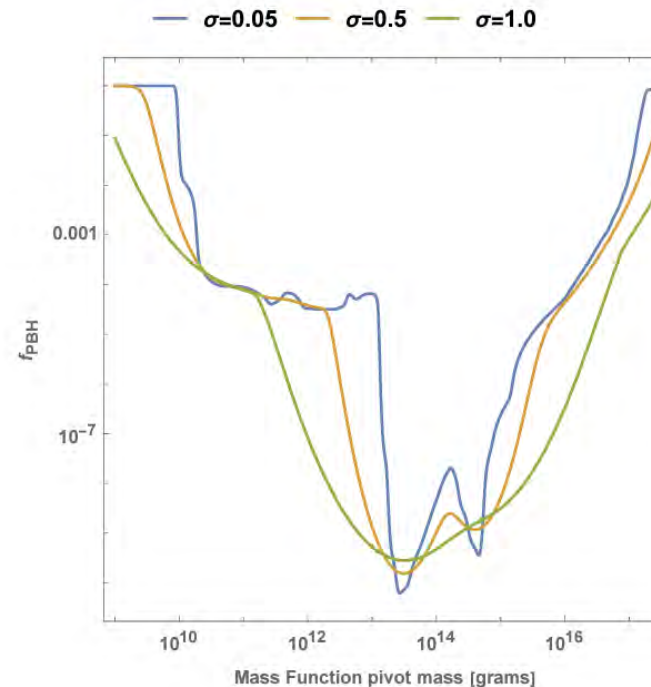
BACKUP SLIDES

- Detection of **Hawking radiation** would be momentous!
- Unique opportunity to discover **dark sectors**!
- Number of handles to **discriminate** against **GRBs**
 - ✓ Light-curve
 - ✓ Spectrum
 - ✓ Distance
 - ✓ Afterglow

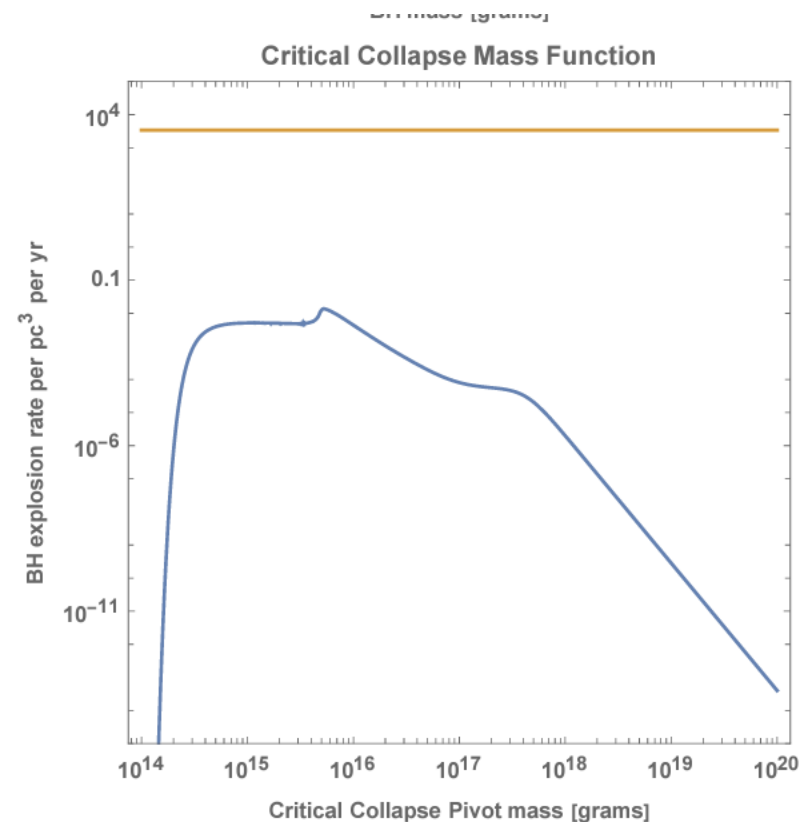
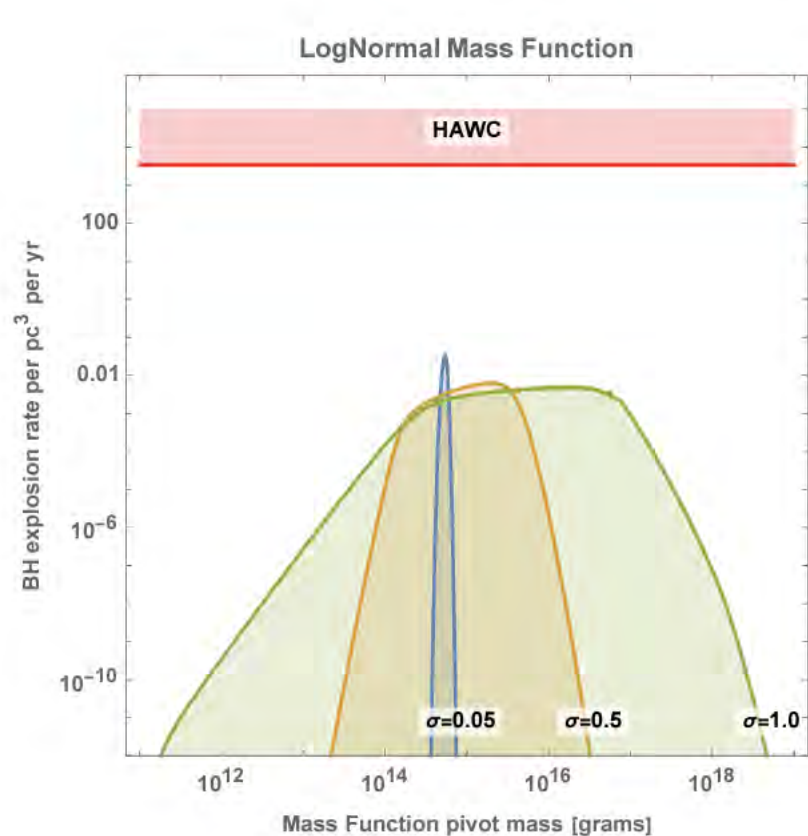
Given a mass function, constraints are calculated via $\int dM \frac{\psi(M)}{f_{\max}(M)} \leq 1.$

For instance, for a lognormal mass function
(typical of a smooth, symmetric power
spectrum density perturbation)

$$\psi(M) = \frac{f_{\text{PBH}}}{\sqrt{2\pi}\sigma M} \exp\left(-\frac{\log^2(M/M_c)}{2\sigma^2}\right)$$



A given width σ and pivot mass M_* produce a different rate of PBH explosions today



HAWC

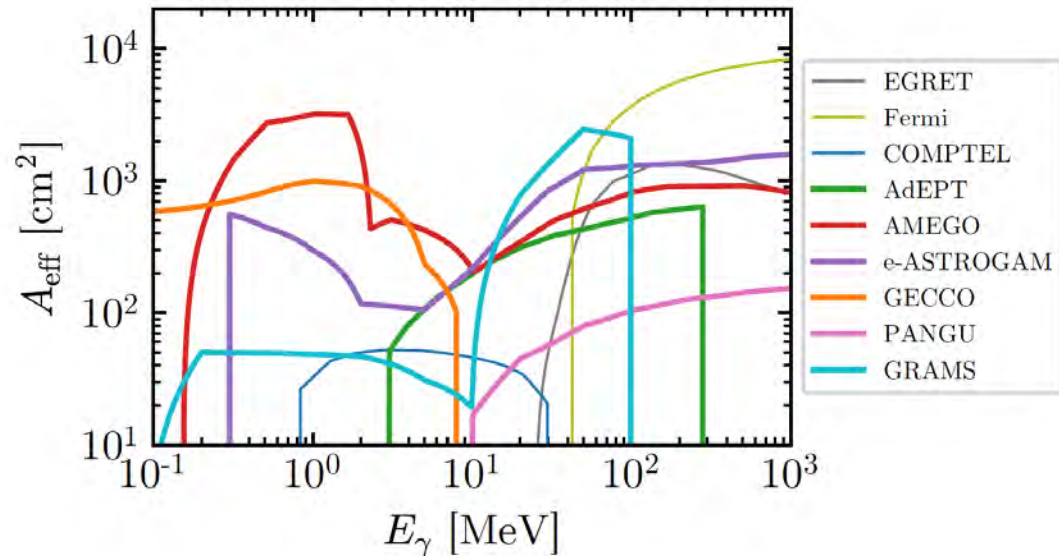


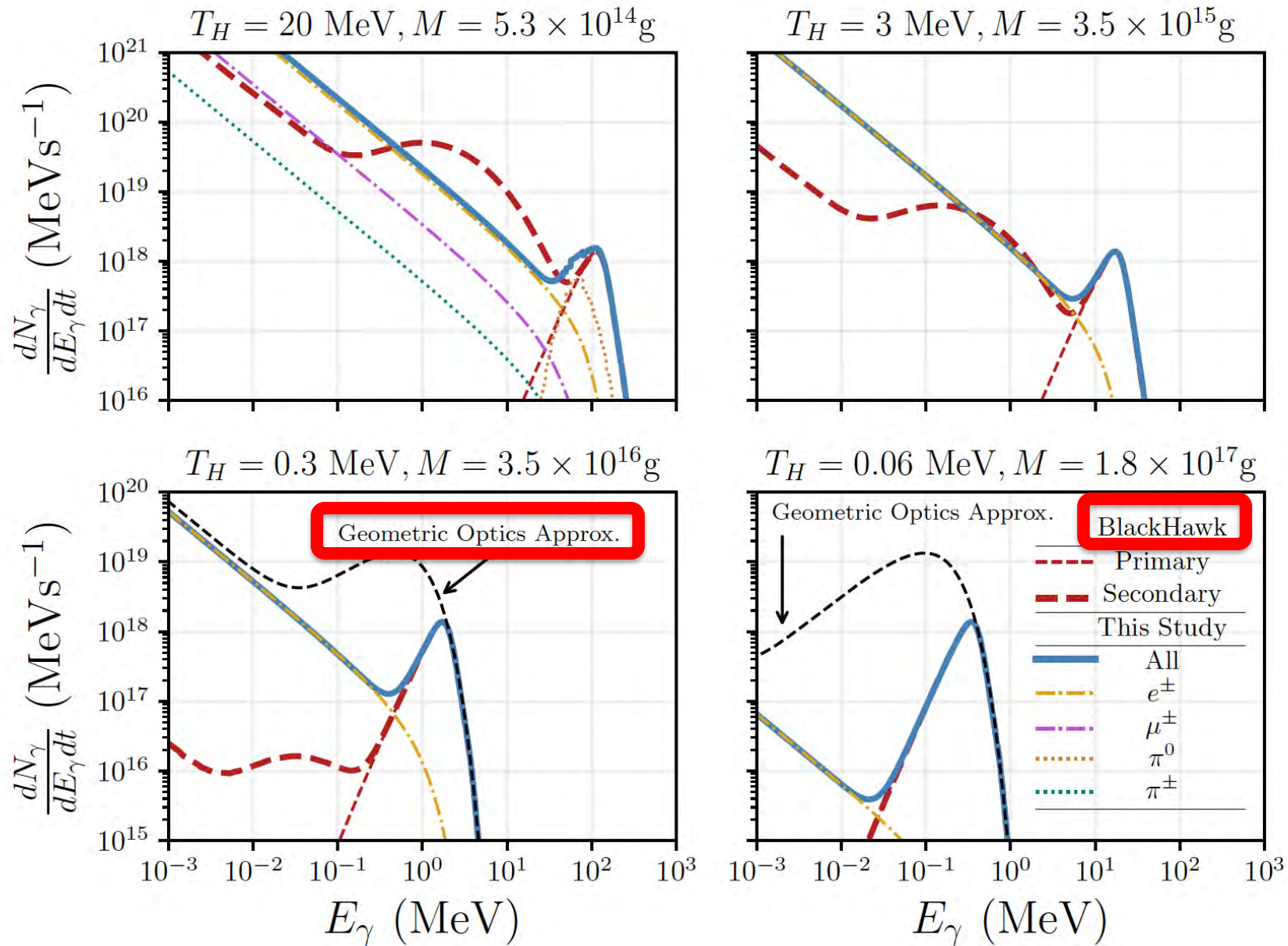
$$\psi(M, M_c) \propto M^{2.85} \exp\left\{-(M/M_c)^{2.85}\right\}$$

Lightest PBH that can be dark matter...

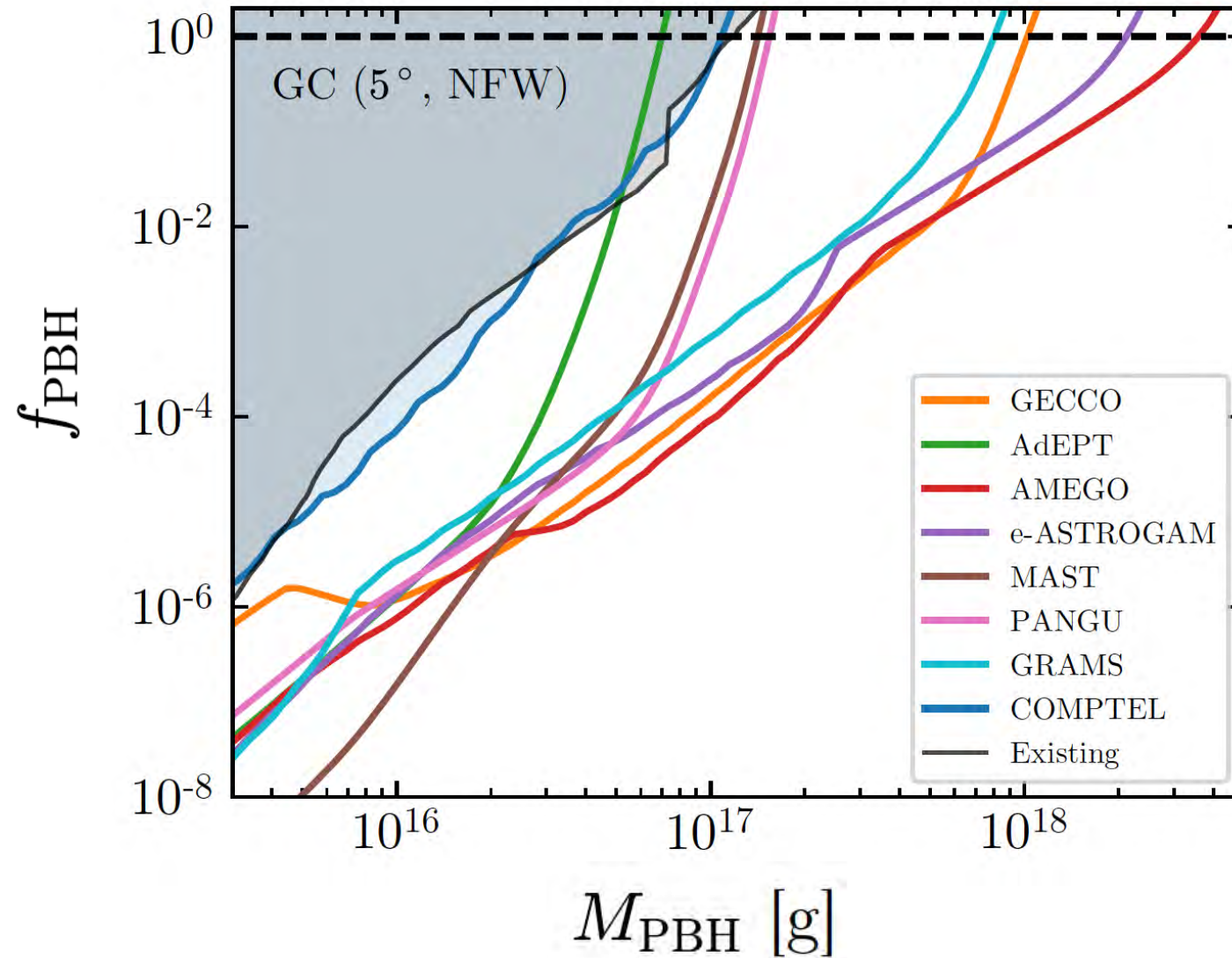
$$\tau(M) \simeq 200 \tau_U \left(\frac{M}{10^{15} \text{ g}} \right)^3 \simeq 200 \tau_U \left(\frac{10 \text{ MeV}}{T_H} \right)^3$$

- are ~ asteroid/comet/**PYRAMID** mass
- can't be much hotter than **10 MeV**





New **MeV Telescopes** could discover Hawking evaporation!



New MeV Telescopes could discover Hawking evaporation!



Snowmass2021 - Letter of Interest

Searching for Dark Matter and New Physics with GECCO

Thematic Areas:

- (CF1) Dark Matter: Particle Like
- (CF2) Dark Matter: Wavelike
- (CF3) Dark Matter: Cosmic Probes
- (CF4) Dark Energy and Cosmic Acceleration: The Modern Universe
- (CF5) Dark Energy and Cosmic Acceleration: Cosmic Dawn and Before
- (CF6) Dark Energy and Cosmic Acceleration: Complementarity of Probes and New Facilities
- (CF7) Cosmic Probes of Fundamental Physics

Contact Information:

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Stefano Profumo (UC Santa Cruz) [profumo@ucsc.edu];

Authors: Alexander Moiseev (CRESST, Greenbelt and NASA, Goddard and Maryland University), Stefano Profumo (UC Santa Cruz and Santa Cruz Institute for Particle Physics), Adam Coogan (Gravitation Astroparticle Physics Amsterdam (GRAPPA), Institute for Theoretical Physics Amsterdam and Delta Institute for Theoretical Physics, University of Amsterdam), Logan Morrison (UC Santa Cruz and Santa Cruz Institute for Particle Physics)

Abstract: We outline the potential science opportunities offered by a future MeV gamma-ray telescope. We point out that such an instrument would play a critical role in opening up a discovery window for particle dark matter with mass in the MeV or sub-MeV range, in disentangling the origin of the mysterious 511 keV line emission in the Galactic center region, and in potentially discovering Hawking evaporation from light primordial black holes. We refer to a new, proposed MeV gamma-ray telescope, the Galactic Explorer with a Coded Aperture Mask Compton Telescope (GECCO) that could deliver on all of those science objectives in the search for new physics and specifically for the nature of dark matter.

Hunting for Dark Matter and New Physics with (a) GECCO

Adam Coogan,^a Alexander Moiseev,^b Logan Morrison, and^{c,d} Stefano Profumo^{c,d}

^aGRAPPA, Institute of Physics, University of Amsterdam, 1098 XH Amsterdam, The Netherlands

^bCRESST, Greenbelt and NASA, Goddard and Maryland University

^cDepartment of Physics, 1156 High St., University of California Santa Cruz, Santa Cruz, CA 95064, USA

^dSanta Cruz Institute for Particle Physics, 1156 High St., Santa Cruz, CA 95064, USA

E-mail: a.m.coogan@uva.nl, amoiseev@umd.edu, loanmorr@ucsc.edu, profumo@ucsc.edu

ABSTRACT: We outline the science opportunities in the areas of searches for dark matter and new physics offered by a proposed future MeV gamma-ray telescope, the Galactic Explorer with a Coded Aperture Mask Compton Telescope (GECCO). We point out that such an instrument would play a critical role in opening up a discovery window for particle dark matter with mass in the MeV or sub-MeV range, in disentangling the origin of the mysterious 511 keV line emission in the Galactic center region, and in potentially discovering Hawking evaporation from light primordial black holes.

arXiv:2101.10370v1 [astro-ph.HE] 25 Jan 2021

$$M_{\text{PBH}} \text{ [g]}$$



Constraints from **Positron** production also heavily affected

BlackHawk

By Alexandre Arbey and Jérémy Auffinger

Calculation of the Hawking evaporation spectra of any black hole distribution

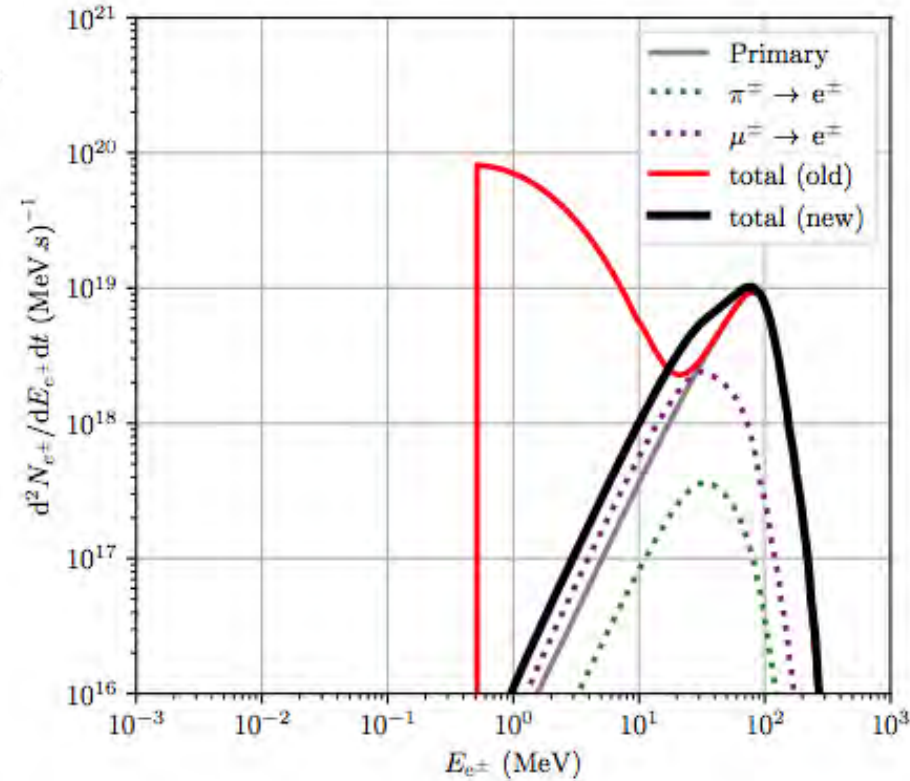
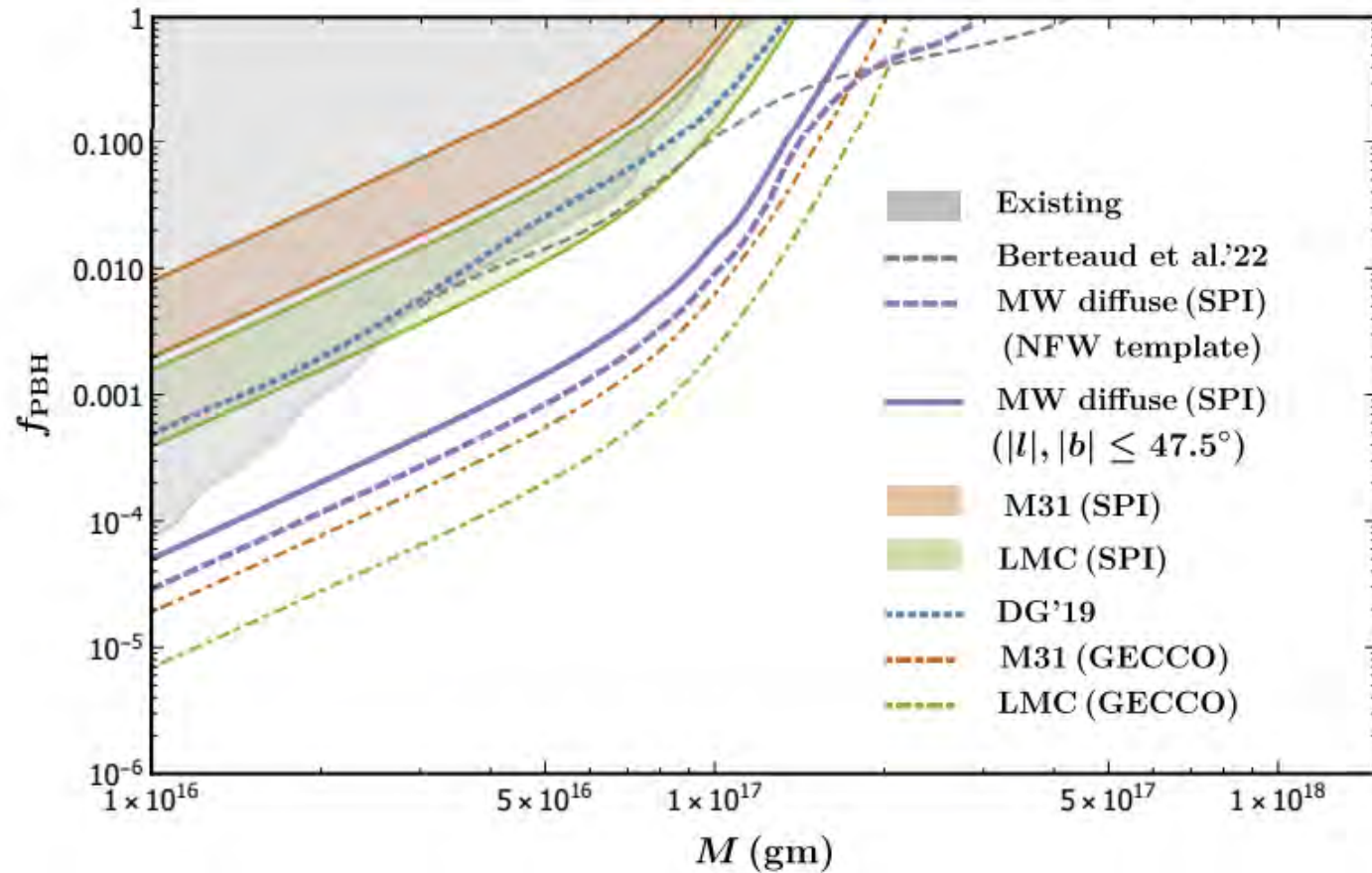
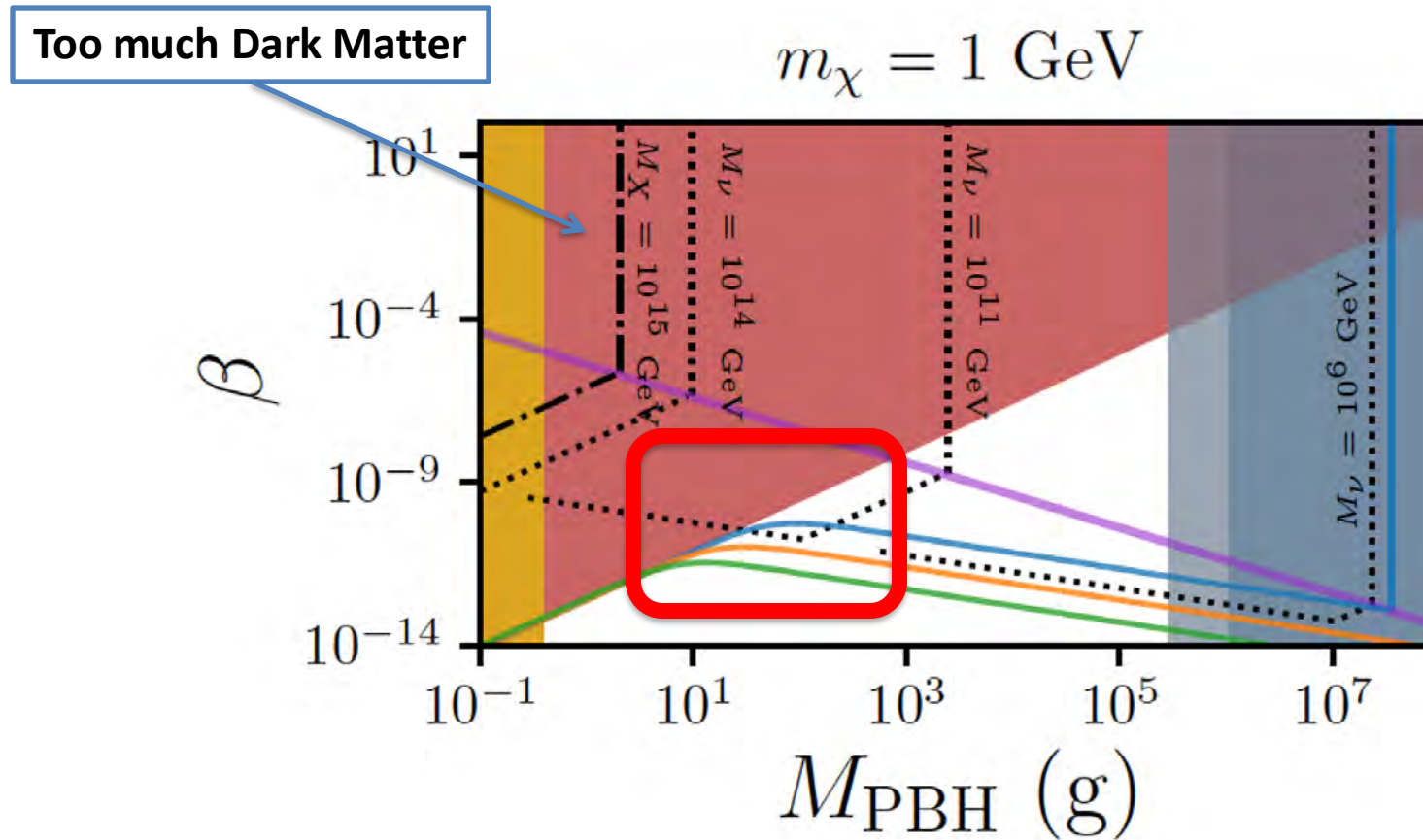


FIG. 5. Comparison between the old PYTHIA extrapolated (solid red) and new Hazma computed (solid black) BlackHawk instantaneous total electron spectrum for a $M_{\text{BH}} = 5.3 \times 10^{14} \text{ g}$. From left to right: primary electron spectrum (grey solid); muon decay (purple dotted); charged pion decay (green dotted).

...speaking of which, most **up-to-date** constraints
from Galactic 511 keV emission from **INTEGRAL/SPI***



Dark Matter can be a **mix** of **Planck-scale relics** from PBH evaporation, and stuff the PBH **evaporated into!**



Hawking-Radiation Recoil of Microscopic Black Holes

Samuel Kováčik ¹

¹*Faculty of Mathematics, Physics and Informatics, University of Žilina, Slovakia*

¹*Department of Theoretical Physics, University of Žilina, Slovakia*

Abstract

The Hawking radiation would cause black holes evaporate rapidly. From many astrophysical considerations it has been argued that the space would alter this behavior of a Planck-size black hole. If left behind is a Planck-mass remnant with a section on the order of $10^{-70} m^2$, detection nearly impossible. Such black hole remnants have been identified as possible dark matter candidates. Here we argue that the final stage of the evaporation has a recoil effect which would give the microscopic black hole velocity on the order of $10^{-1}c$ which is in disagreement with the cold dark matter cosmological model.

Black hole remnants are not too fast to be dark matter

Benjamin V. Lehmann^{1,2,*} and Stefano Profumo^{1,2,†}

¹*Department of Physics, University of California Santa Cruz, 1156 High St., Santa Cruz, CA 95064, USA*

²*Santa Cruz Institute for Particle Physics, 1156 High St., Santa Cruz, CA 95064, USA*

We comment on recent claims that recoil in the final stages of Hawking evaporation gives black hole remnants large velocities, rendering them inviable as a dark matter candidate. We point out that due to cosmic expansion, such large velocities at the final stages of evaporation are not in tension with the cold dark matter paradigm so long as they are attained at sufficiently early times. In particular, the predicted recoil velocities are robustly compatible with observations if the remnants form before the epoch of big bang nucleosynthesis, a requirement which is already imposed by the physics of nucleosynthesis itself.

of the striking difference compared to the ordinary black hole theory is that the Hawking temperature [9] defined to be proportional to the surface gravity at the horizon does not grow indefinitely but instead drops to zero at small but positive mass, resulting in a microscopic black hole remnant.

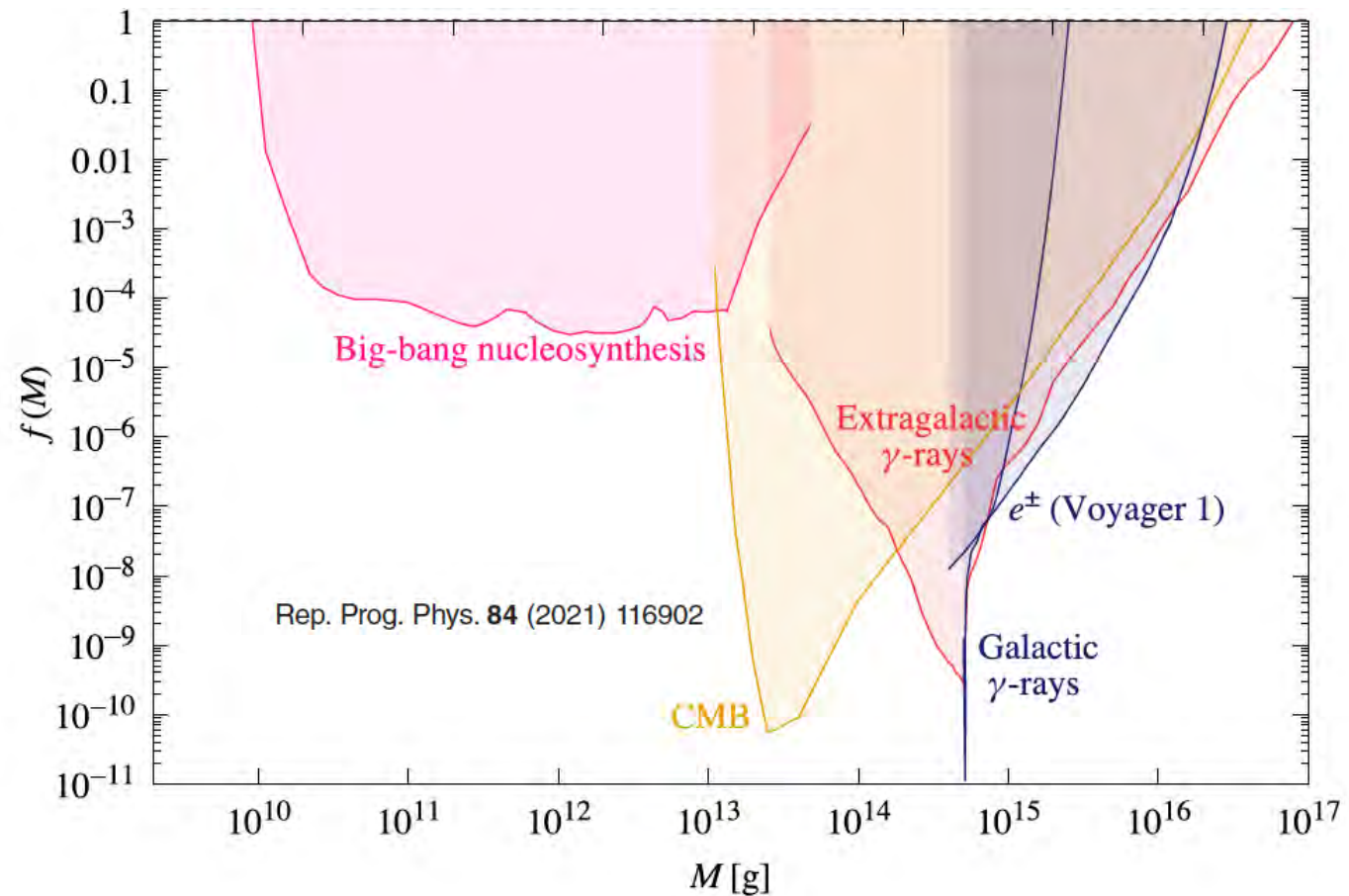
Black holes remnants have been considered as

...true only if **evaporation** stops **very late**
(much later than **BBN**), which
cannot happen!

* Lehmann and Profumo, 2105.01627

At earlier times, evaporation perturbs BBN, CMB

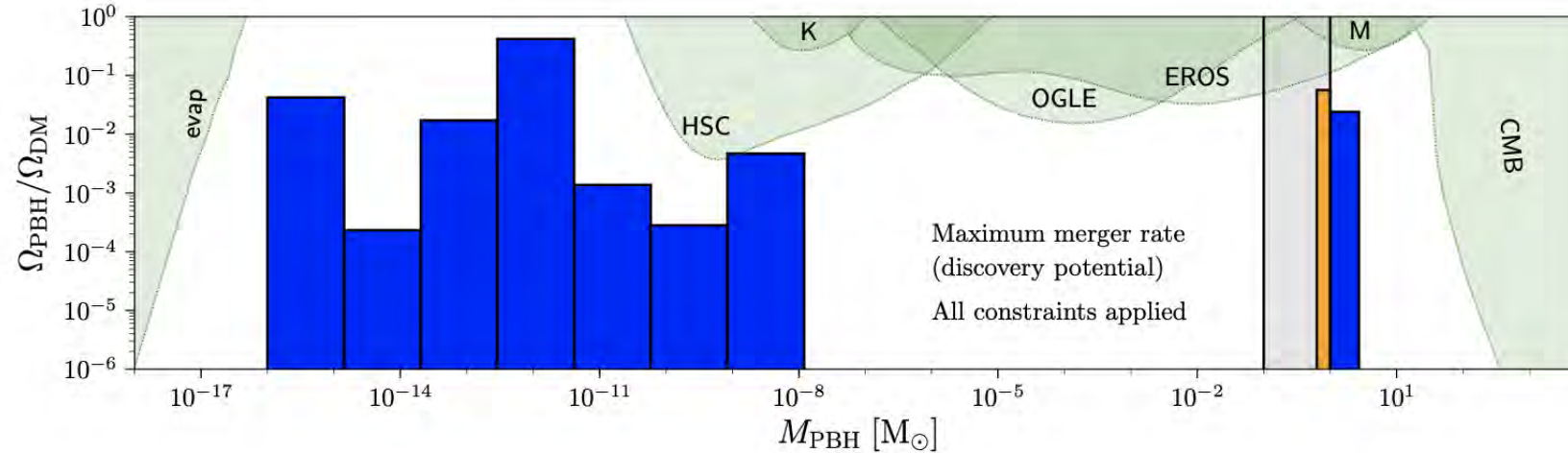
Here, perhaps,
gravity waves!
(but very high
frequency!)*



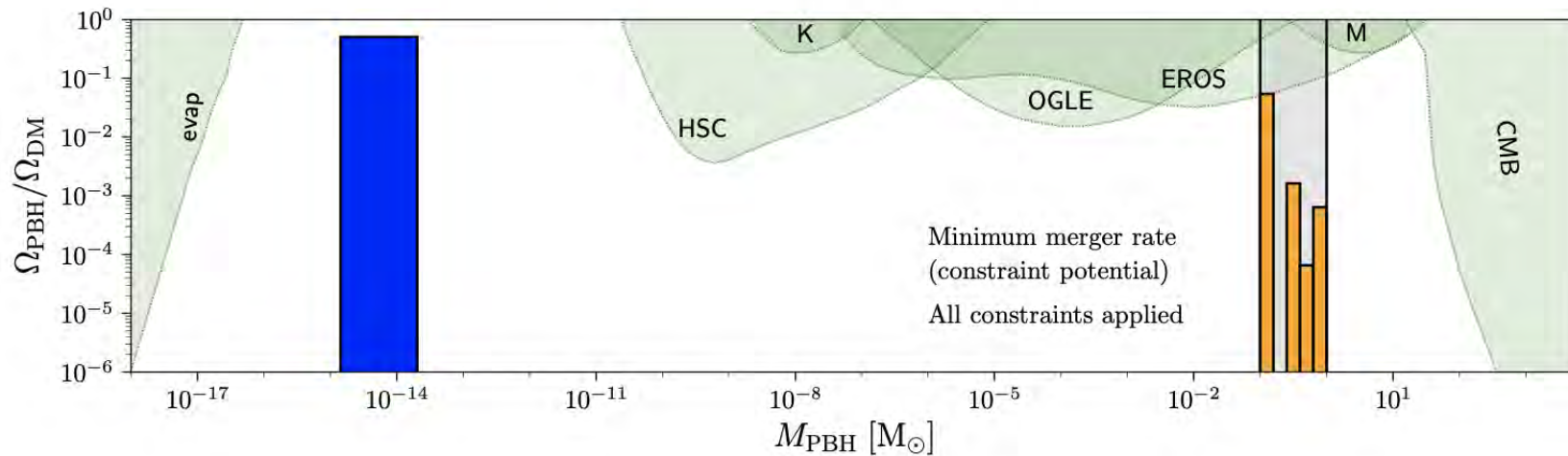
Notice that here M [g] is the initial mass at PBH formation

*Ireland, Profumo and Scharnhorst, 2302.10188

We can **numerically** compute the **maximal** and **minimal** possible “**goldilocks event rate**” for a given mass fraction of **light+detectable BHs**

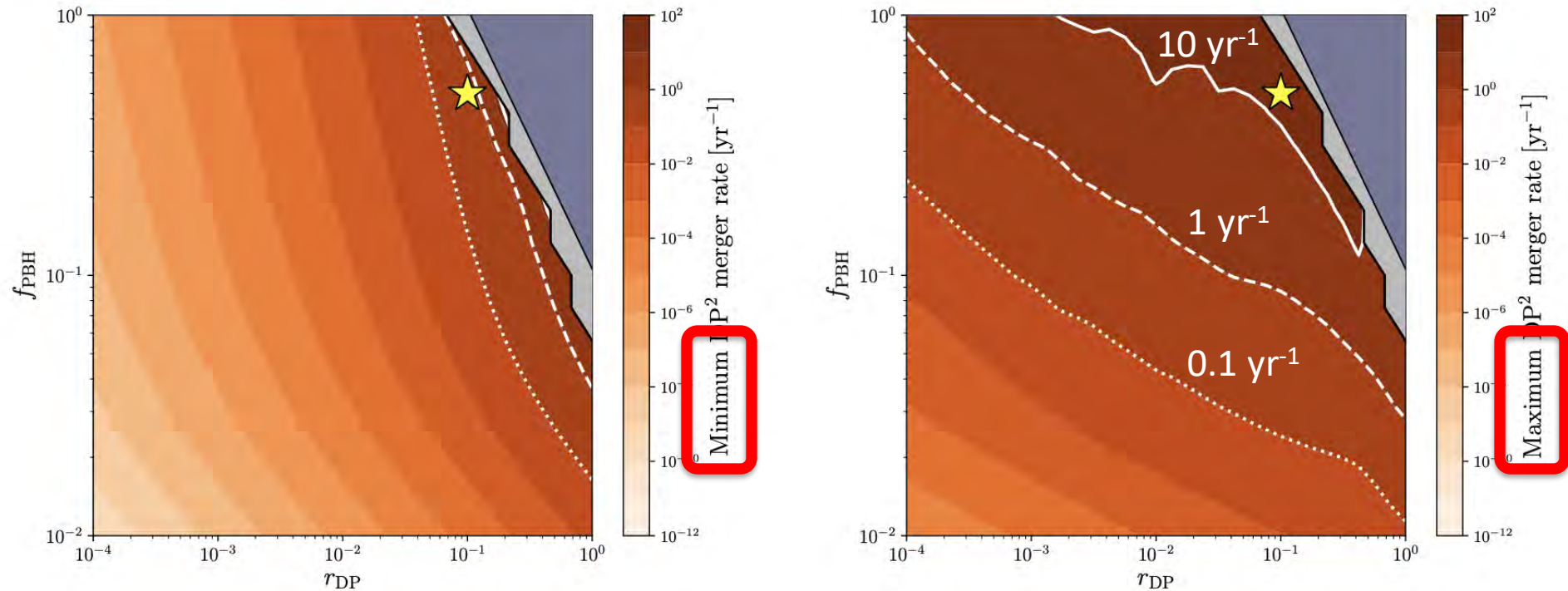


Maximum
Merger Rate
(@fixed $f_{\text{PBH}}, r_{\text{DP}}$)



Minimum
Merger Rate
(@same $f_{\text{PBH}}, r_{\text{DP}}$)

...if at least some PBH are **light+detectable**,
the **minimum/maximum** rate of “**Goldilocks events**”
are encouraging! **LIGO** searches are ongoing!



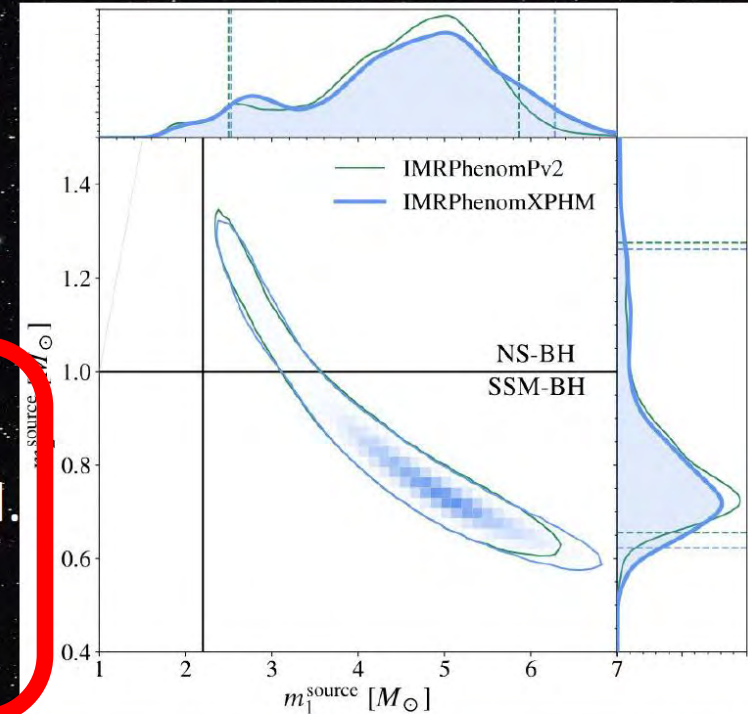
$$\int dm \psi(m) = \frac{\Omega_{\text{PBH}}}{\Omega_{\text{DM}}} \equiv f_{\text{PBH}}.$$

$$r_{\text{DP}} = \frac{1}{f_{\text{PBH}}} \int_{m_{\text{DP}}^{\text{min}}}^{m_{\text{DP}}^{\text{max}}} dm \psi(m).$$

...if at least some PBH are **light+detectable**,
the **minimum/maximum** rate of “**Goldilocks events**”
are encouraging! **LIGO** searches are ongoing! [**SSM170401**]**

- For 16% of the posterior we have $m_1 > 2.2M_\odot$ and $m_2 > 1.0M_\odot$, i.e. masses compatible with an NS-BH.
 - The primary would be a BH in the mass gap
 - The secondary could be a NS or a BH

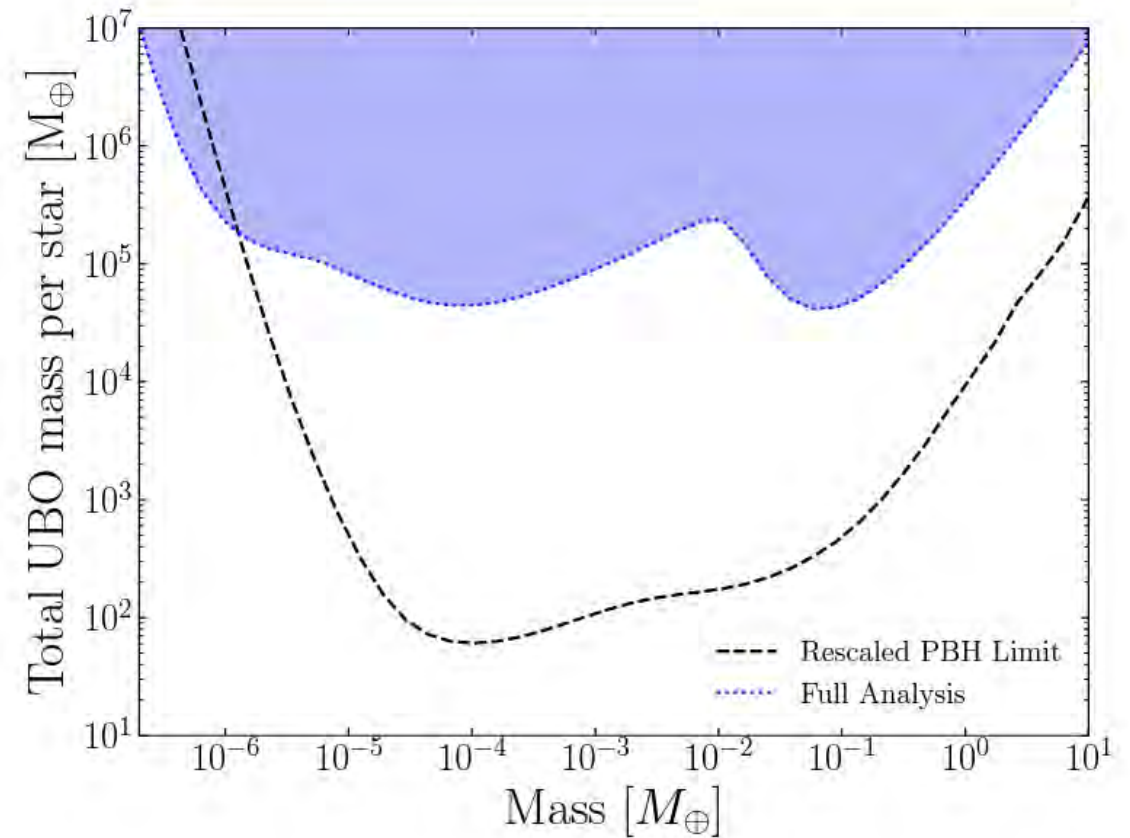
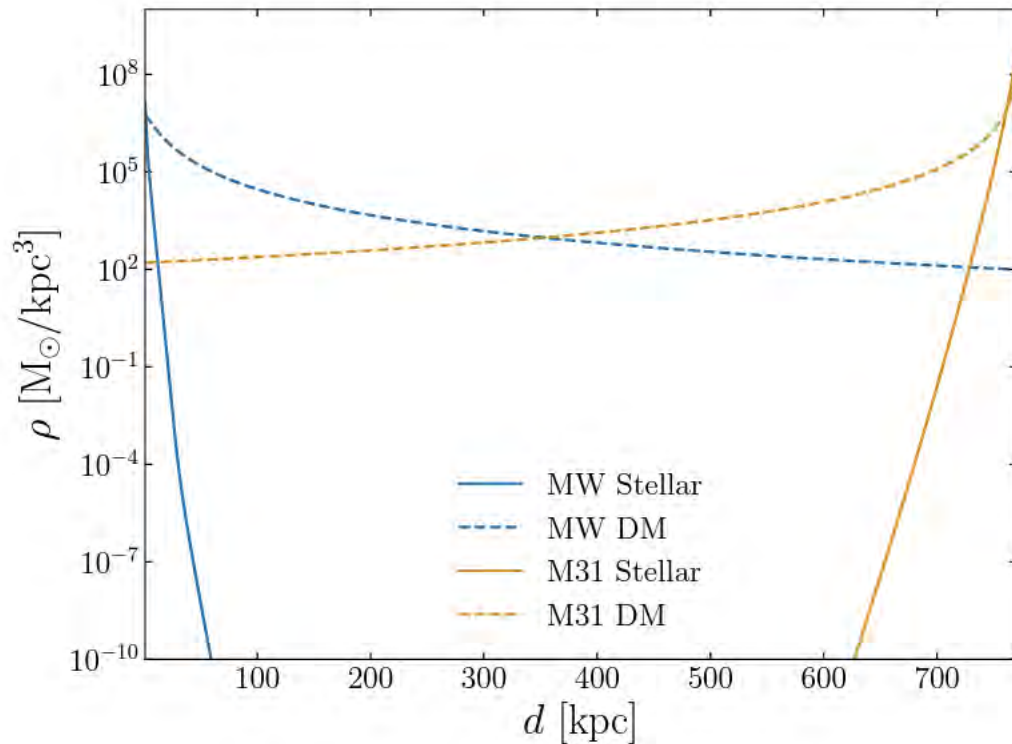
- For 84% of the posterior we have $m_1 > 2.2M_\odot$ and $m_2 < 1.0M_\odot$, i.e. masses compatible with an SSM-BH.
 - The primary can be in the mass gap or above
 - The secondary is an SSM black hole.



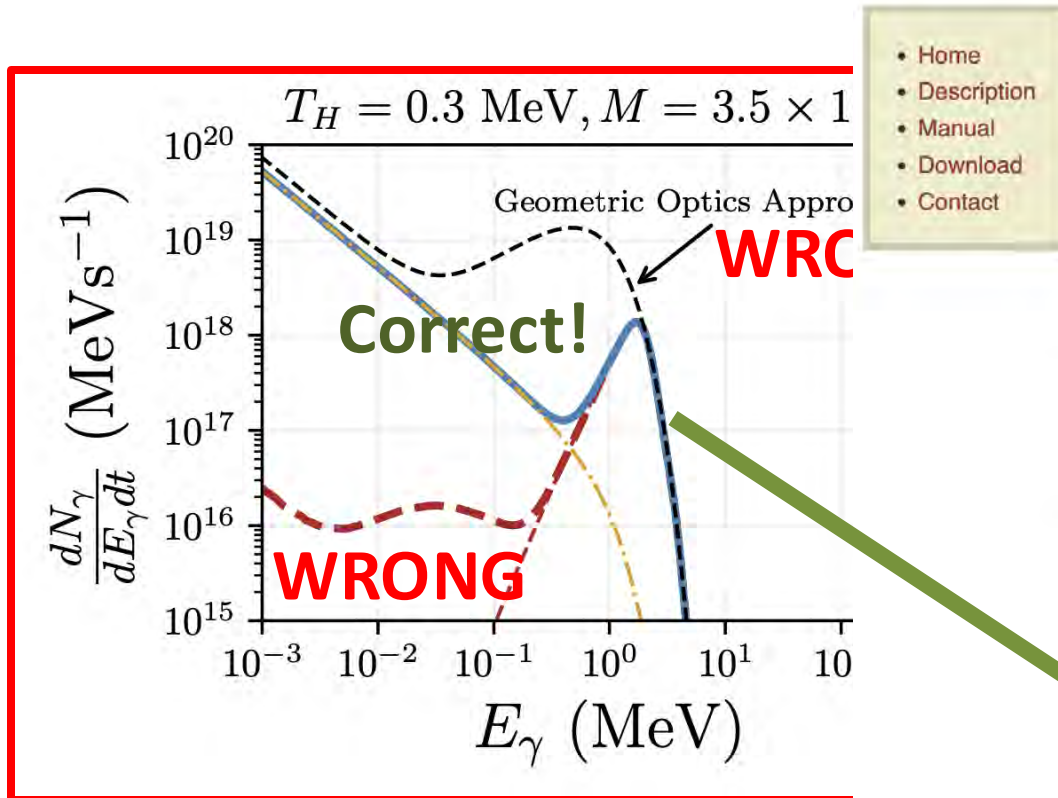
- In all cases, astrophysical models have problems generating a system with these masses

Do the **HSC** observation provide any constraints on **FFPs***?

Simply rescaling the PBH limit vastly overestimate the HSC constraints!



Evaporation products (gamma rays, cosmic-ray positrons) are **detectable**, constraining the fraction of light PBH that can be the DM



BlackHawk

By **Alexandre Arbey** and **Jérémy Auffinger**

Calculation of the Hawking evaporation spectra of any black hole distribution

BlackHawk is a public C program for calculating the Hawking evaporation spectra of any black hole distribution. This program enables the users to compute the primary and secondary spectra of stable or long-lived particles generated by Hawking radiation of the distribution of black holes, and to study their evolution in time.

If you use BlackHawk to publish a paper, please cite:

- A. Arbey and J. Auffinger, *Eur. Phys. J. C* 79 (2019) 693, arXiv:1905.04268 [gr-qc]
- A. Arbey and J. Auffinger, *Eur. Phys. J. C* 81 (2021) 910, arXiv:2108.02737 [gr-qc]

If you use the hadronized spectra of BlackHawk, we also advise that you cite the corresponding particle physics code:

- PYTHIA spectra (hadronization_choice = 0 or 2): T. Sjöstrand *et al.*, *Comput. Phys. Commun.* 191 (2015) 159-177, arXiv:1410.3012 [hep-ph]
- HERWIG spectra (hadronization_choice = 1): J. Bellm *et al.*, *Eur. Phys. J. C* 76 (2016) 4, 196, arXiv:1512.01178 [hep-ph]
- Hazma spectra (hadronization_choice = 3): A. Coogan, L. Morrison, S. Profumo, *JCAP* 01 (2020) 056, arXiv:1907.11846 [hep-ph]
- HDMSpectra spectra (hadronization_choice = 4): C. W. Bauer, N. L. Rodd, B. R. Webber, *JHEP* 06 (2021) 121, arXiv:2007.15001 [hep-ph]

Coogan, Morrison and Profumo, 2010.04797, PRL 126 (2021) 17, 171101

An exciting possibility is to see the **terminal PBH explosion!**

The **rate of explosions** today depends on the PBH **mass function**

Considering an initial mass function Ψ_i
the rate of PBH explosions today reads

$$\dot{n}_{\text{PBH}} = \rho_{\text{DM}} \frac{\psi_i(M_U)}{3t_U}$$

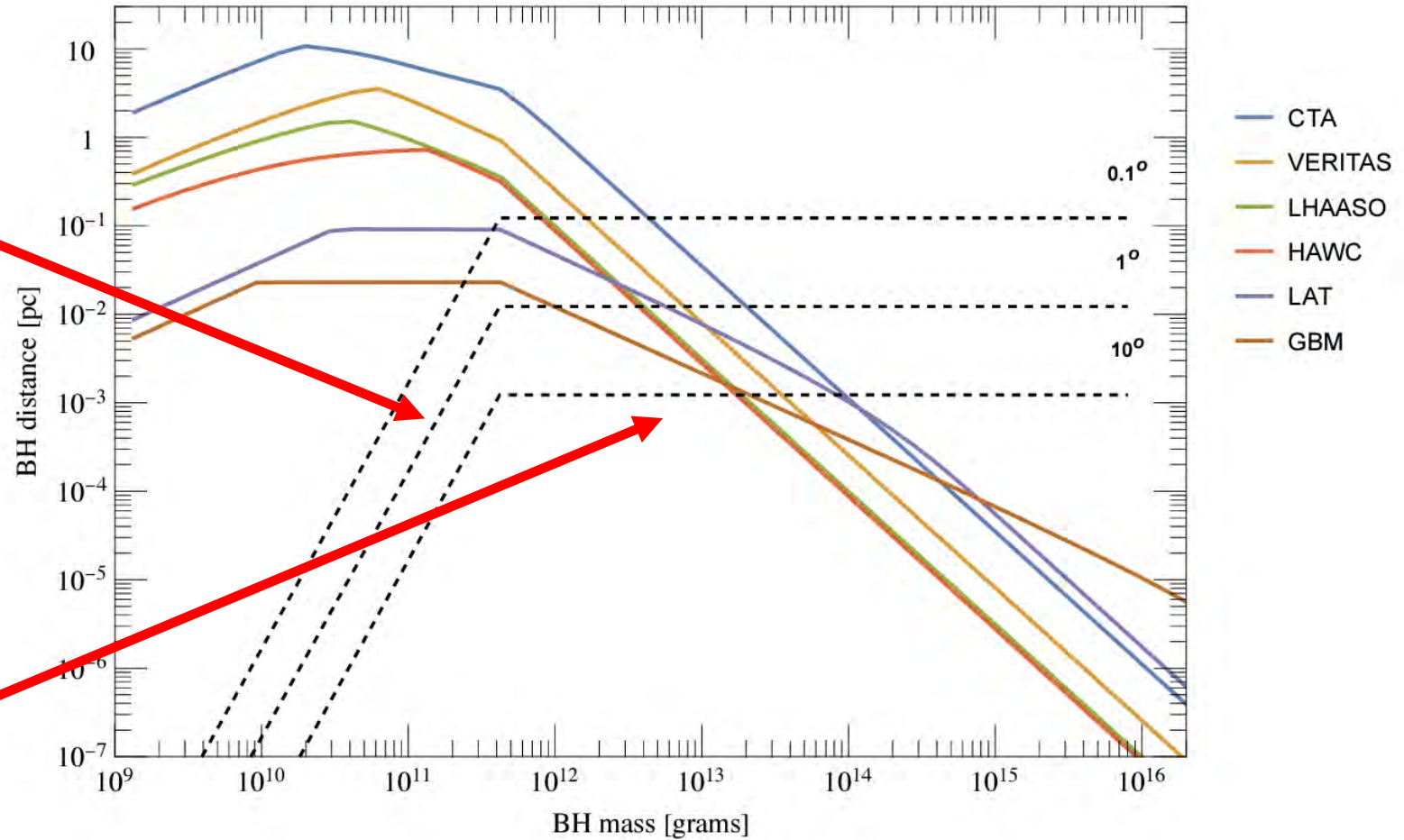
Additionally, PBH can be “**spawned**” at late times**

How can we search for **PBH explosions** today?

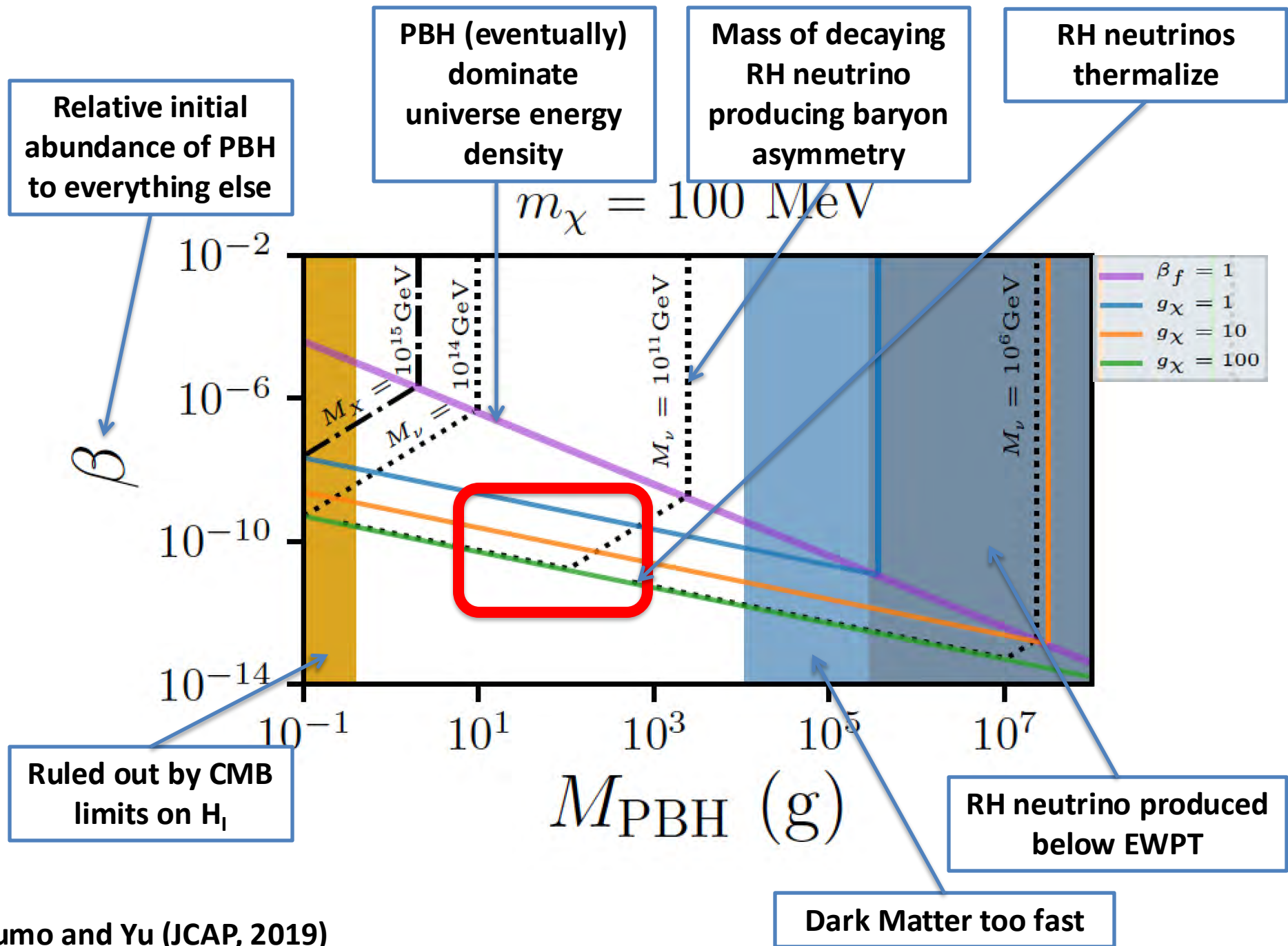
Proper motion

Distance

Evaporation time
scale > Max
observation time

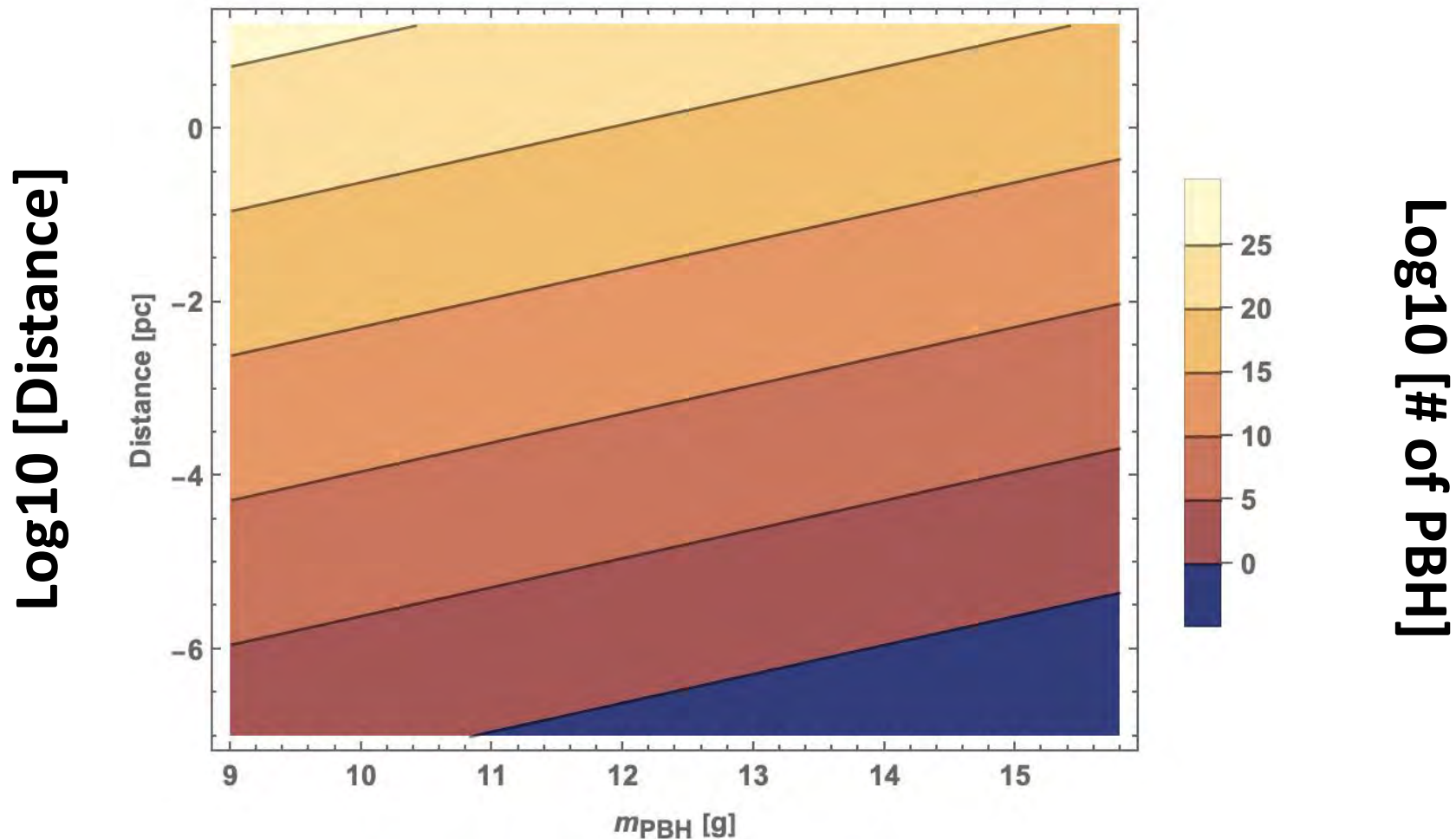


Mass/Lifetime



* Morrison, Profumo and Yu (JCAP, 2019)

[could there actually be **PBH** around, so **close**? Yes.]
[is the rate **significant**? It depends on the **mass function**]



$\text{Log}_{10} (\text{Mass/Lifetime})$ [today, at detection]

Alternately, evaporation can seed gravitational baryo- (or DM-) genesis



L. Santos-Olmsted

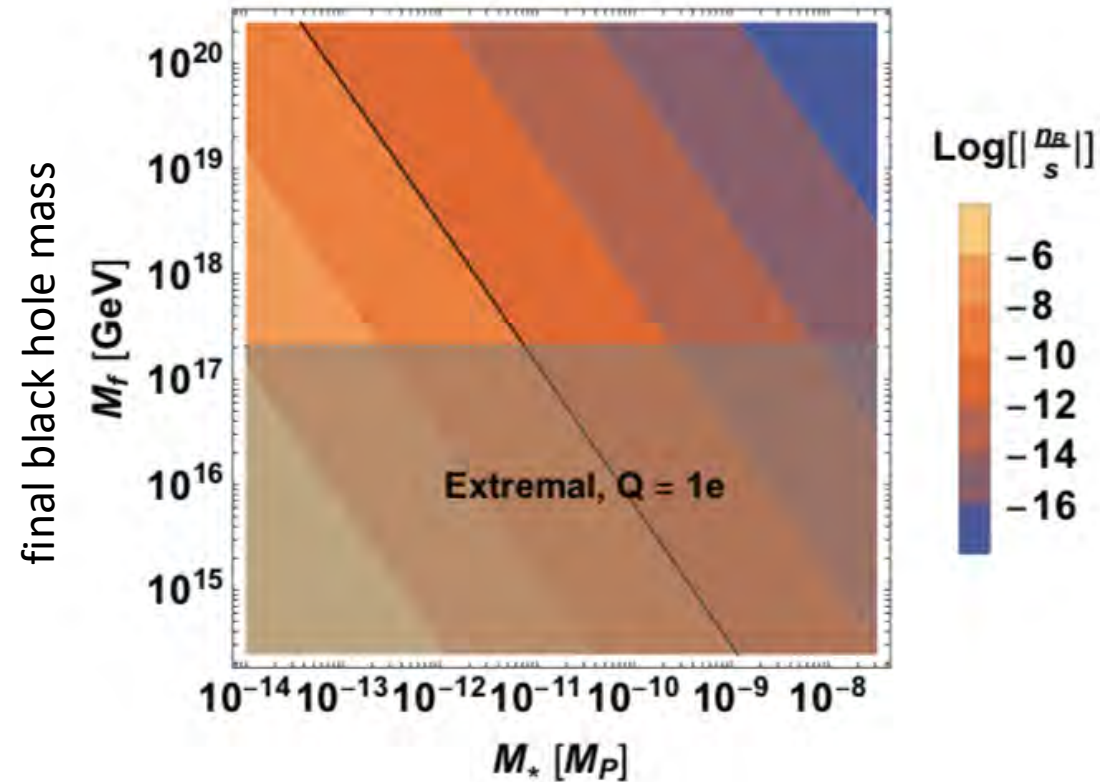


N. Smyth

$$\mathcal{S} \supset \frac{1}{M_*^4} \partial_\alpha (\mathcal{R}_{\mu\nu\rho\sigma} \mathcal{R}^{\mu\nu\rho\sigma}) J_{\text{B-L}}^\alpha$$

Kretschmann scalar

Relic PBH

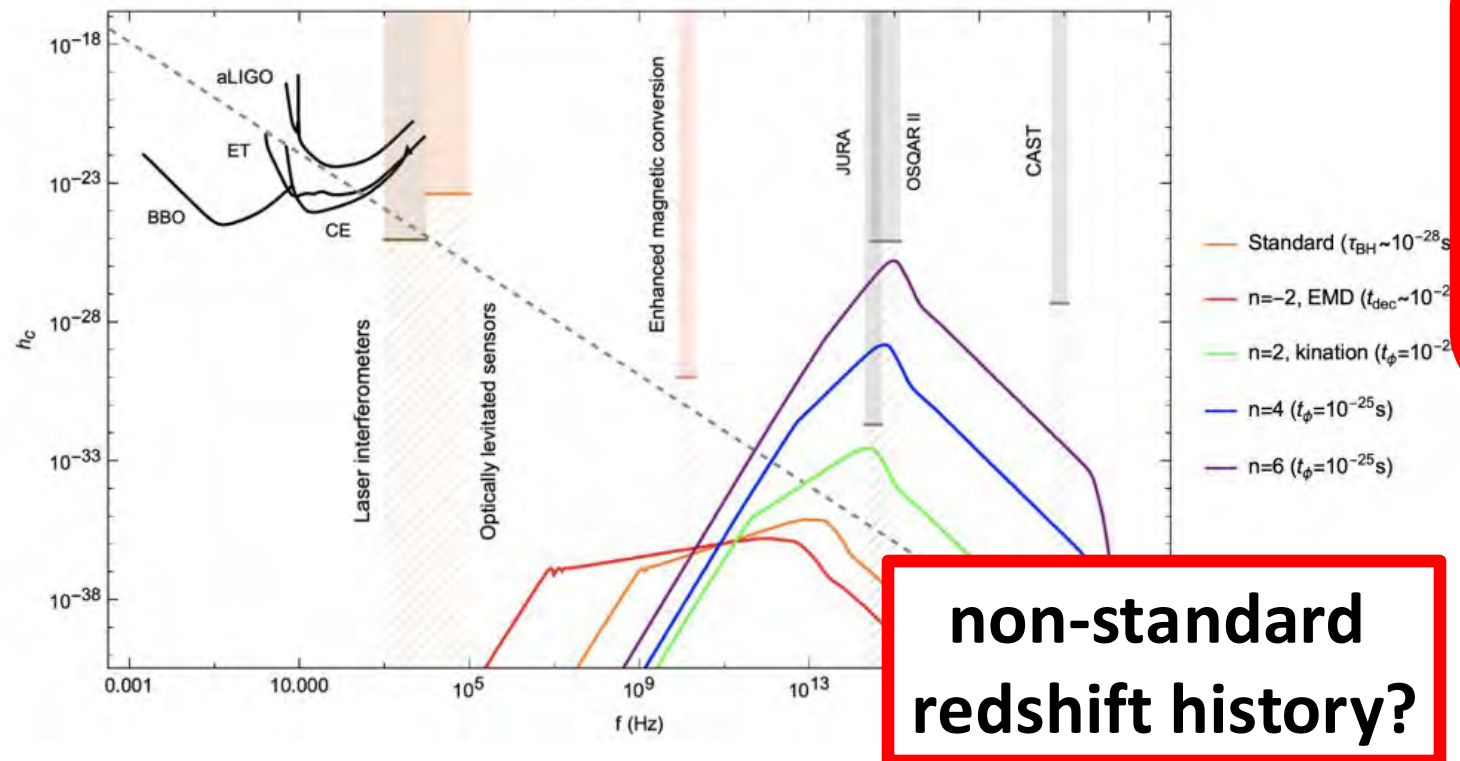


...even if PBH are **NOT** the dark matter, they can **PRODUCE**
the **dark matter** AND the **baryon asymmetry**
via **Hawking evaporation!**

Mass (g)	T_H (GeV)	τ (s)	$T_{\text{evap}} = T(\tau)$ (GeV)
$5M_P \simeq 10^{-4}$	1.7×10^{17}	10^{-41}	2×10^{17}
1	1.7×10^{13}	4×10^{-29}	2×10^{11}
10^3	1.7×10^{10}	4×10^{-20}	6×10^6
10^6	1.7×10^7	4×10^{-11}	200
10^9	1.7×10^4	0.04	0.006
10^{12}	17	$4 \times 10^7 \sim 1 \text{ yr}$	$\sim 1 \text{ keV}$

ruled out by **BBN**

If evaporation **completes** in the very early universe ($\tau \ll t_{\text{BBN}}$), the only **relic** that could be observable today is **gravitational waves**!

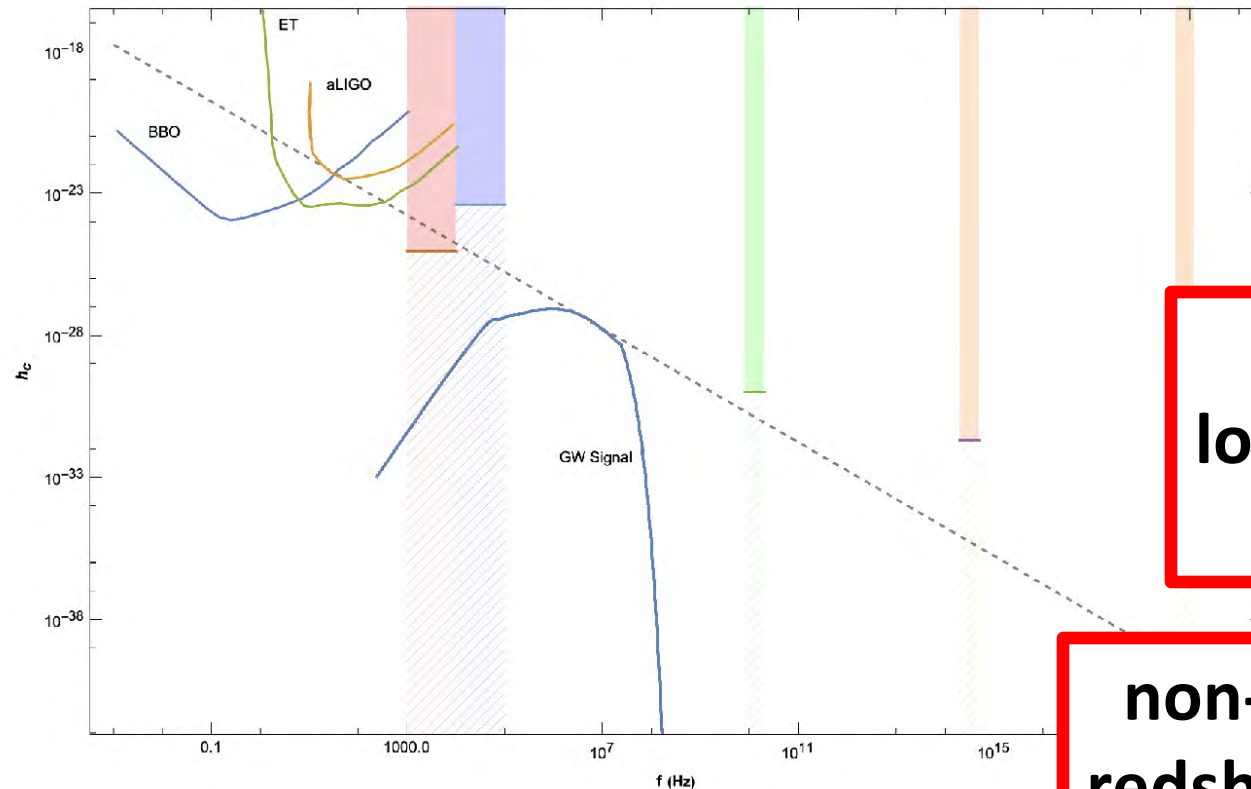


A. Ireland



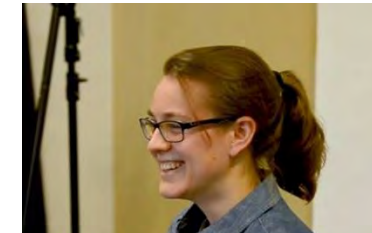
J. Sharnhorst

If evaporation **completes** in the very early universe ($\tau \ll t_{\text{BBN}}$), the only **relic** that could be observable today is **gravitational waves**!



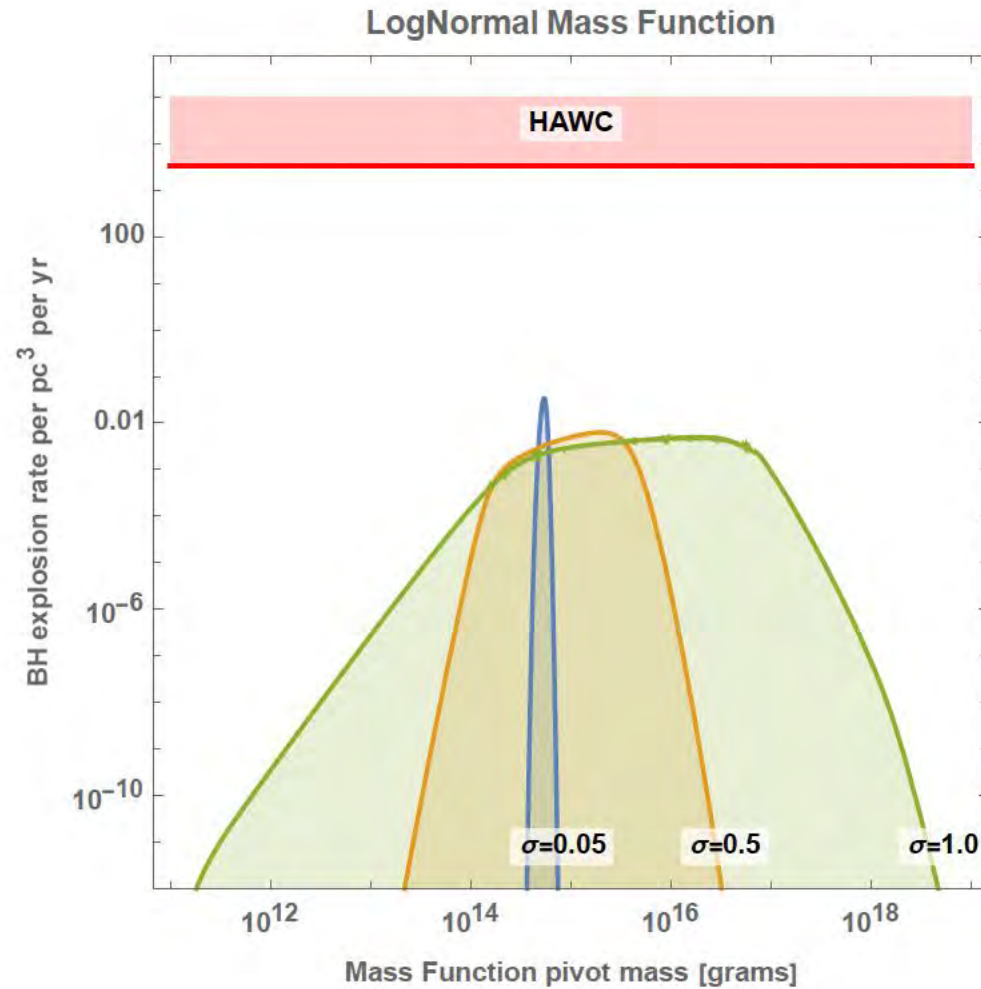
with extra dimensions,
lower M_{Planck} peak can shift
to low frequencies!

non-standard
redshift history?



J. Sharnhorst

The rate of explosions today depends on the PBH mass function



$$\dot{n}_{\text{PBH}} \simeq \frac{1.2 \times 10^{-3} \text{ pc}^{-3} \text{ yr}^{-1}}{\sigma},$$

...we need to **stack** lots of particles!

$$N \sim 10^8 M_{\text{sun}} / (10^{-22} \text{ eV}) \sim 10^{90}$$

→ this mass limit can only apply to **bosons**, $J=n$

what if $J=(2n+1)/2$, i.e. **fermion**?

→ we expect fermions respecting Pauli's **exclusion principle**
to be $10^{90/4}$ **fewer**, thus $m > 10^{22.5} 10^{-22} \sim$ **few eV**

Phase space density is bounded : $f = gh^{-3}$

Upper limit: highest **observed** phase space density: **dSph!**

$$\frac{g}{h^3} \geq n \cdot f_p \geq \frac{\rho_{\text{DM}}}{m} \frac{1}{\left(m \cdot \sqrt{2\pi\sigma^2}\right)^3} \quad (MB \text{ with } exp=1)$$

$$m^4 \geq \frac{\rho_{\text{DM}} h^3}{[g(2\pi\sigma^2)^{3/2}]} \sim (25 \text{ eV})^4$$

Tremaine-Gunn limit (1979)



m > 25 eV

Asteroid mass PBH as dark matter **well established** ...even **string theorists** take this seriously!

Monochromatic Mass Spectrum of Primordial Black Holes

Matthew Kleban and Cameron Norton

Center for Cosmology and Particle Physics, New York University, New York, NY

October 17, 2023

Abstract

During slow-roll inflation, non-perturbative transitions can produce bubbles of metastable vacuum. These bubbles expand exponentially during inflation to super-horizon size, and later collapse into black holes when the expansion of the universe is decelerating. In a broad class of models, the inflationary fine-tuning that gives rise to small density fluctuations causes these bubbles to appear only during a time interval that is short compared to the inflationary Hubble time. As a result, despite the fact that the final mass of the black hole is exponentially sensitive to the moment bubbles form during inflation, the resulting primordial black hole mass spectrum can be nearly monochromatic. If the transition occurs near the middle of inflation, the mass can fall in the “asteroid” range $10^{17} - 10^{22}$ g in which all known observations are compatible with black holes comprising 100% of dark matter.

1 Introduction

One of the greatest mysteries in modern physics is the nature of dark matter. Despite accounting for over 25% of the energy density of our universe, its nature and origin remains uncertain. Decades of

The rate of explosions today depends on the PBH **mass function**

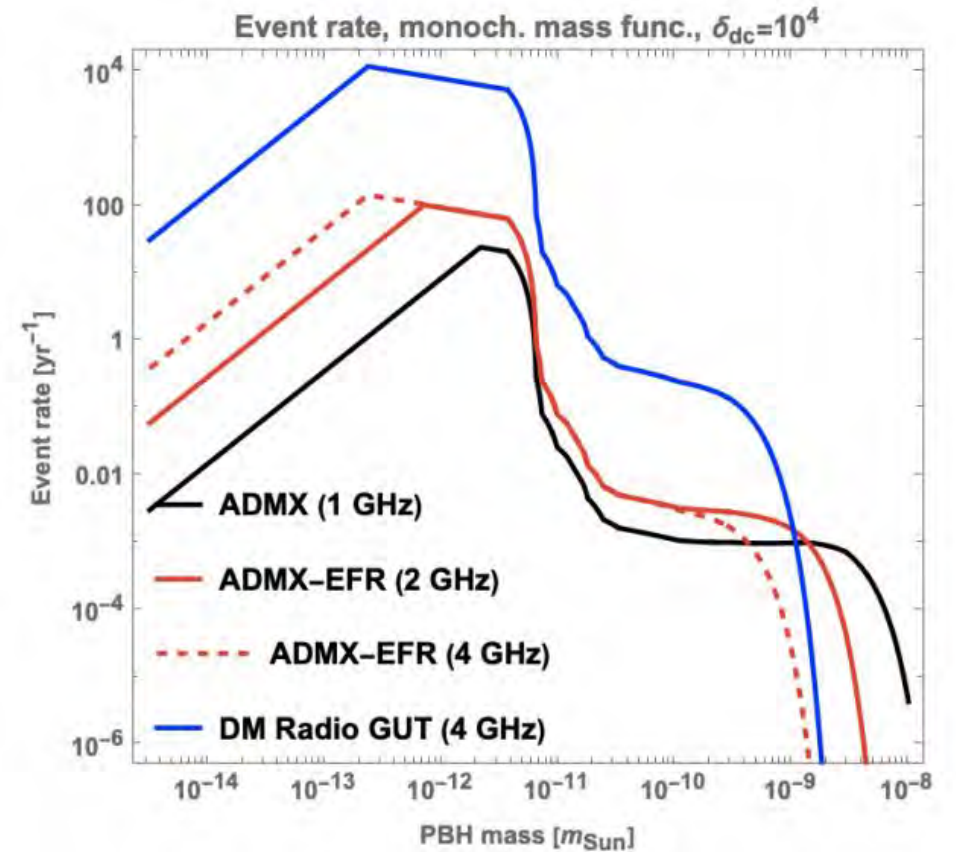
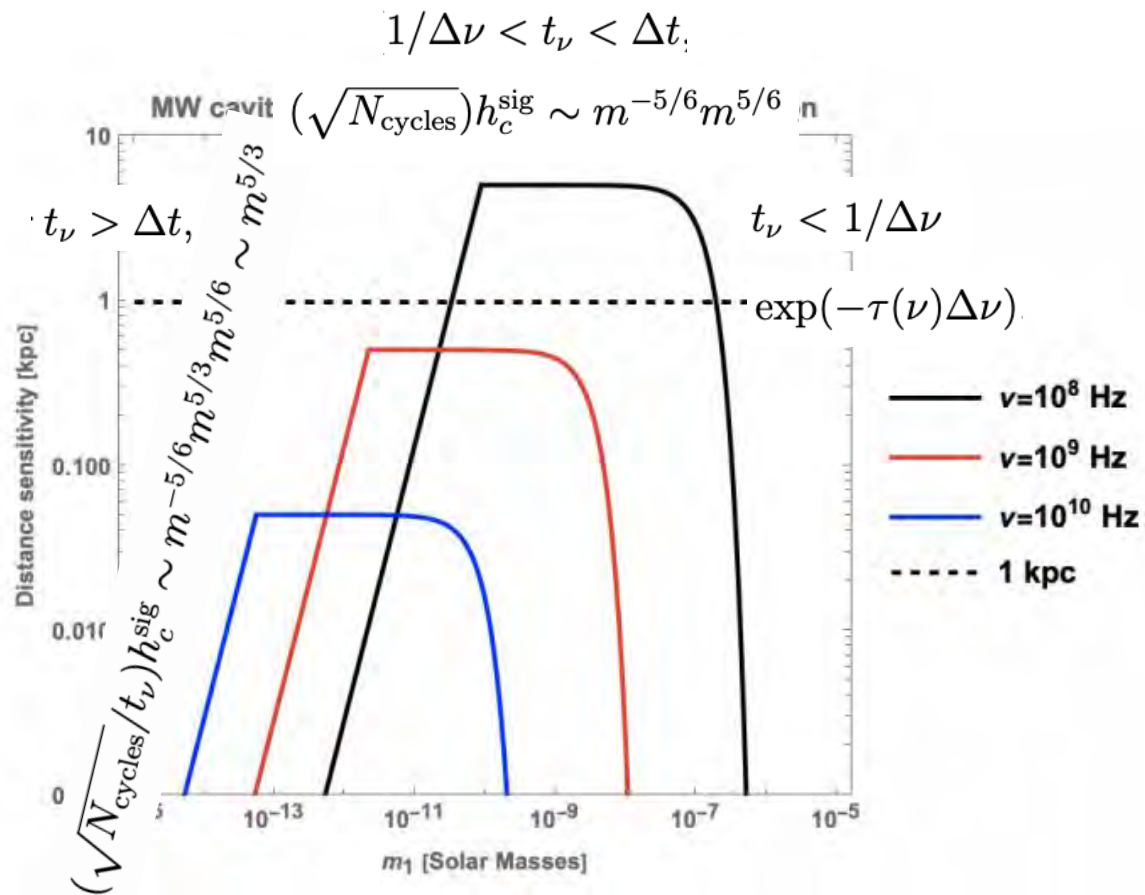
$$\psi(M) \propto M \frac{dn_{\text{PBH}}}{dM} \quad \int dM \psi(M) = f_{\text{PBH}}$$

$$\dot{n}_{\text{PBH}} = \frac{dn_{\text{PBH}}}{dM} \left(-\frac{dM}{dt} \right) = \rho_{\text{DM}} \frac{\psi(M)}{M} \frac{\alpha(M)}{M^2} \quad \frac{\alpha(M)}{M^2} \simeq 4 \times 10^{-4} \text{g/sec}$$

Considering an initial mass function Ψ_i
the **rate of PBH explosions** today reads

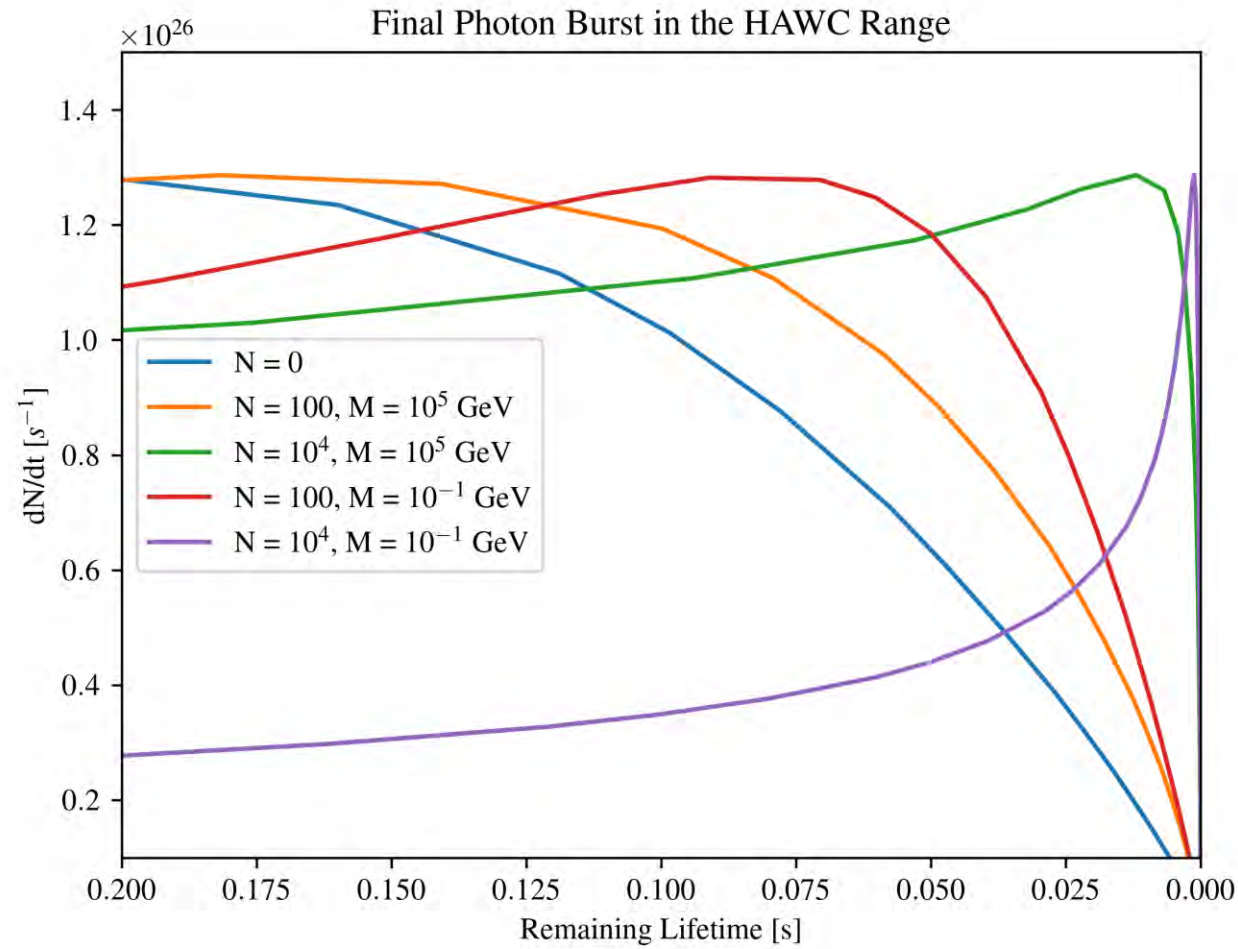
$$\dot{n}_{\text{PBH}} = \rho_{\text{DM}} \frac{\psi_i(M_U)}{t_U}$$

Critical distance and event rate



The (**non-universal**) Black Hole **Light-curve**

Photon Emission Rate
at **HAWC** telescope



Time to Explosion

new degrees
of freedom

Mass-scale of
new d.o.f

SM emission
from new d.o.f

