A Light $L_{\mu} - L_{\tau}$ Gauge Boson for $g_{\mu} - 2$, Dark Matter etc

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(8)Summary

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 4.2σ deviation in $g_{\mu} - 2$: evidence for new interaction of muons!

 \implies consider $L_{\mu} - L_{\tau}!$

The Model

$$\mathcal{L} = \mathcal{L}_{SM} - g_X X_\rho \left(\bar{\mu} \gamma^\rho \mu - \bar{\tau} \gamma^\rho \tau + \overline{\nu_{\mu,L}} \gamma^\rho \nu_{\mu,L} - \overline{\nu_{\tau,L}} \gamma^\rho \nu_{\tau,L} \right) - \frac{1}{4} X_{\rho\sigma} X^{\rho\sigma} + \frac{1}{2} m_X^2 X_\rho X^\rho + \bar{\psi} \left(i \partial - m_\psi - g_X q_\psi X \right) \psi$$

 $X_{\rho\sigma} = \partial_{\rho} X_{\sigma} - \partial_{\sigma} X_{\rho}$: new field strength tensor

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Relevant parameters:

- New gauge coupling g_X
- Mass of new gauge boson m_X
- DM mass m_{ψ}
- DM charge q_{ψ}

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$$\delta \mathcal{L} = \epsilon e X_{\rho} J_{\text{em}}^{\rho}$$
, with $\epsilon = \frac{e g_X}{6\pi^2} \ln\left(\frac{m_{\tau}}{m_{\mu}}\right) \simeq \frac{g_X}{70}$

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Gives contribution to $a_{\mu} = (g_{\mu} - 2)/2$ (Leveille, 1977):

$$\Delta a_{\mu} = \frac{g_X^2}{8\pi^2} \int_0^1 dx \frac{2m_{\mu}^2 x^2 (1-x)}{x^2 m_{\mu}^2 + (1-x)m_X^2}$$



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Need:

$$g_X \simeq \begin{cases} 4.4 \cdot 10^{-4} & m_X^2 \ll m_\mu^2 \\ 5.4 \cdot 10^{-4} \frac{m_X}{m_\mu} & m_X^2 \gg m_\mu^2 \end{cases}$$

Neutrino Trident

Refers to $\nu_{\mu}p \rightarrow \nu_{\mu}\mu^{+}\mu^{-}p$. (3 charged tracks, hence "trident".) Has been observed in 1980's (CHARM, CCFR); rediscovered as BSM test by Pospelov (2006).



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Measurement agrees with SM prediction \implies strong upper bound on g_X for given m_X !

Dark Matter

Want to employ usual "thermal WIMP" mechanism:

$$\Omega_{\psi}h^2 \sim 0.1 \cdot \frac{1 \text{ pb}}{\langle v\sigma(\psi\bar{\psi} \to \text{anything else}) \rangle} \stackrel{!}{=} 0.12$$

- Ω : scaled mass density;
- *h*: scaled Hubble parameter;
- $\langle \dots \rangle$: thermal average.

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$$\begin{split} & \psi \bar{\psi} \to \nu_{\mu,\tau} \bar{\nu}_{\mu,\tau}; \\ & \psi \bar{\psi} \to \ell^+ \ell^- \ (\ell = \mu, \tau), \text{ if } m_{\psi} > m_{\ell}; \\ & \psi \bar{\psi} \to XX, \text{ if } m_{\psi} > m_X. \end{split}$$

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• $g_{\mu} - 2$ constraint: $\sigma(\psi \bar{\psi} \rightarrow \text{anything})$ is too small unless $m_{\psi} \simeq m_X/2!$ (Resonance enhancement.)

Relic Density



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Warning: MadDM did not work! Used method of Griest and Seckel (1991)

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Will need $m_{\psi} < m_{\mu} \implies$ scattering on nuclei is not sensitive!

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Scattering on electrons:

$$\sigma(\psi e \to \psi e) = \frac{\mu_e^2}{\pi} \frac{\epsilon^2 e^2 q_\psi^2 g_X^2}{(m_X^2 + \alpha^2 m_e^2)^2},$$

$$\simeq 6 \cdot 10^{-44} \text{ cm}^2 \left(\frac{g_X}{10^{-3}}\right)^4 \left(\frac{10 \text{ MeV}}{m_X}\right)^4 q_\psi^2$$

$$\mu_e = \frac{m_\psi m_e}{m_\psi + m_e}$$

Exptl bound (SENSEI): $\sigma < 5 \cdot 10^{-37}$ cm²: No problem!

$N_{\rm eff}$ and Hubble Tension

BBN (in particular, ⁴He fraction) depends on radiation density, which can be affected by additional "light" degrees of freedom: $\rho_{\rm rad} = \left[1 + N_{\rm eff} \frac{7}{8} \left(\frac{4}{11}\right)^{\frac{4}{3}}\right] \rho_{\gamma}$.

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- Light ψ annihilate into $\nu_{\mu,\tau}$: decouple from electrons at $T_{\rm d} \simeq 2.3 \text{ MeV}$

 $\implies \bar{\psi}\psi \rightarrow \nu_{\mu\tau}\bar{\nu}_{\mu,\tau}$ annihilations at $T < T_d$ increase $N_{\text{eff}}!$

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$$N_{\text{eff}} = N_{\nu} \left[1 + \frac{1}{N_{\nu}} \sum_{i=\psi,X} \frac{g_i}{2} F\left(\frac{m_i}{T_{\nu,D}}\right) \right]^{4/3},$$
$$F(x) = \frac{30}{7\pi^4} \int_x^\infty dy \frac{(4y^2 - x^2)\sqrt{y^2 - x^2}}{e^y \pm 1}.$$

Result

$$r = m_{\psi}/m_X$$



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BBN: $N_{\rm eff} < 3.4$

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Interpret $3.2 \le N_{\text{eff}} \le 3.4$ as relaxing Hubble tension.

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 - $pp \rightarrow 3\mu + \text{missing } E_T$: published LHC limits weaker than Trident bound. MD, M. Shi, Z. Zhang, 1811.12446

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Evidence for BSM contribution for WD cooling! Isern et al., 1805.00135

Evidence for additional WD cooling



Black: SM; red, blue: extra cooling (axions)

■ Requires kinetic X - γ mixing $\implies \sigma(\psi\bar{\psi} \to e^+e^-) \simeq 2 \cdot 10^{-5} \sigma(\psi\bar{\psi} \to \nu\bar{\nu})$

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 $10^{-3} \mathsf{fb} \le \left\langle \sigma(\psi \bar{\psi} \to e^+ e^-) v \right\rangle \cdot \left(\frac{m_{\psi}}{1 \text{ MeV}}\right)^{-2} / 2 \le 1 \mathsf{fb}$

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- For $m_{\psi} \simeq 10$ MeV: Only $\sim 50\%$ of required e^+ can come from DM annihilation; otherwise too many MeV photons. (Beacom and Yuksel, astro-ph/0512411)

Final Results



Grey: excluded by CMB; purple: excluded by BBN (N_{eff}) green: preferred by $g_{\mu} - 2$; blue: relaxes Hubble tension

Dependence on DM charge q_{ψ} $\delta = 4m_{\psi}^2/m_X^2 - 1$



Black: correct relic density; red: excluded by CMB; green: favored by 511 keV excess

Predictions

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- ~ $250 \gamma + \text{nothing events at BELLE-2}$, with $E_{\gamma} \simeq \sqrt{s}/2$ in cms system; comparable no. of events at super $\tau/\text{charm factory}$, if $\mathcal{L} = 2 \text{ ab}^{-1}/\text{year N}$. Borodatchenkova et al., hep-ph/0510147 Needs single photon trigger!

Contributing Diagram



Effective $g_{eeM} \simeq -eg_{\mu\mu M}/70$.

Summary

• $L_{\mu} - L_{\tau}$ model with $m_{\psi} \sim m_X/2 \simeq 10$ MeV, $g_X \cdot 4.5 \cdot 10^{-4}, q_{\psi} \simeq 1$ explains lots of things: Dark Matter; $g_{\mu} - 2$; easing of Hubble tension; contribution to 511 keV photon flux.

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- $L_{\mu} L_{\tau}$ model with $m_{\psi} \sim m_X/2 \simeq 10$ MeV, $g_X \cdot 4.5 \cdot 10^{-4}, q_{\psi} \simeq 1$ explains lots of things: Dark Matter; $g_{\mu} - 2$; easing of Hubble tension; contribution to 511 keV photon flux.
- Can be tested by future neutrino experiments (JUNO, COHERENT-II), possible high–luminosity low–energy e^+e^- colliders