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Non-Hermiticity: a new paradigm for model building in particle physics

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To **you** for listening in.

To my collaborators:

Jean Alexandre (King’s College London)

Carl M. Bender (University of Washington, St Louis)

Maxim N. Chernodub (CNRS, Université de Tours)

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Model building strategies for new physics

To go beyond the Standard Model of particle physics, we can:

- **Add new degrees of freedom**: extra gauge singlets, extra Higgs doublets, heavy neutrinos, SUSY partners, hidden sectors, …
- **Relax assumptions**: number of spatial dimensions, Lorentz invariance, locality, CPT invariance, ...

Non-Hermiticity?

All Hermitian matrices have real eigenvalues …

… but **matrices with real eigenvalues need not be Hermitian.**

Enter PT-symmetric QM

- We can relax Hermiticity in favour of the **weaker condition** of **PT-symmetry**, i.e., the combined action of parity and time-reversal (or indeed any antilinear symmetry).

  [Bender & Boettcher ’98, see also Mostafazadeh ’02 and Mannheim ’18a]

- This is sufficient to guarantee both **real eigenvalues**

  [Bender & Boettcher ’98]

  and **unitary evolution**.

  [Bender, Brody & Jones ’02]

For a recent review, see Bender, PT Symmetry in Quantum and Classical Physics, World Scientific ‘2019.

But we want to do **non-Hermitian quantum field theory** …

A scalar playground

A simple **scalar model** with ***c*-number Lagrangian** with **real parameters**:

[Alexandre, PM & Seynaeve '17]

$$\mathcal{L} = \partial_\nu \phi_i^* \partial^\nu \phi_i - m_i^2 \phi_i^* \phi_i - \mu^2 (\phi_1^* \phi_2 - \phi_2^* \phi_1) \qquad i = 1,2$$

"Naïve" $\mathcal{PT}$ symmetry if we have a complex scalar and a complex pseudoscalar, i.e.,

$$\mathcal{P}:\ \phi_1 \longrightarrow +\phi_1, \qquad \phi_2 \longrightarrow -\phi_2$$
$$\mathcal{T}:\ \phi_1 \longrightarrow +\phi_1^*, \qquad \phi_2 \longrightarrow +\phi_2^*$$

A scalar playground: matrix model

Non-Hermitian squared mass matrix: $M^2 = \begin{pmatrix} m_1^2 & \mu^2 \\ -\mu^2 & m_2^2 \end{pmatrix}$

With the $\mathcal{PT}$ symmetry of $\mathcal{L}$ translating into the **pseudo-Hermiticity** of M^2:

$$[P \cdot M^2 \cdot P]_{ij} = [M^2]_{ji} \qquad P = \begin{pmatrix} +1 & 0 \\ 0 & -1 \end{pmatrix}$$

Mass spectrum: $M_\pm^2 = \frac{m_1^2+m_2^2}{2} \pm \left[\left(\frac{m_1^2-m_2^2}{2}\right)^2 - \mu^4\right]^{1/2} \in \mathbb{R} \text{ if } \eta \equiv \frac{2|\mu^2|}{|m_1^2-m_2^2|} \leq 1$

M^2 is defective at the **exceptional point** at $\eta = 1$.

A scalar playground: matrix model

Eigenvectors $(m_1^2 - m_2^2, \mu^2 > 0)$:

$$\boldsymbol{e}_+ = N \begin{pmatrix} \eta \\ -1 + \sqrt{1-\eta^2} \end{pmatrix} \qquad \boldsymbol{e}_- = N \begin{pmatrix} -1 + \sqrt{1-\eta^2} \\ \eta \end{pmatrix} \qquad \eta = \frac{2\mu^2}{m_1^2 - m_2^2}$$

- **Not** orthogonal with respect to the Dirac inner product: $\boldsymbol{e}_\pm^* \cdot \boldsymbol{e}_\mp \neq 0$.
- **Not** orthonormal with respect to the $\mathcal{PT}$ inner product: $\boldsymbol{e}_\pm^* \cdot P \cdot \boldsymbol{e}_\pm \ngtr 0$.
- Orthonormal with respect to the $\mathcal{C}'\mathcal{PT}$ inner product: $\boldsymbol{e}_\pm^* \cdot C' \cdot P \cdot \boldsymbol{e}_\pm = 1$.

  [Bender, Brody & Jones ’02; see also Alexandre, Ellis & PM ’20b; cf. Mannheim ’18b]

A scalar playground: matrix model

$$C' = R \cdot P \cdot R^{-1} = \frac{1}{\sqrt{1-\eta^2}} \begin{pmatrix} 1 & -\eta \\ \eta & -1 \end{pmatrix} \qquad R \cdot M^2 \cdot R^{-1} = \widehat{M}^2 = \begin{pmatrix} M_+^2 & 0 \\ 0 & M_-^2 \end{pmatrix}$$

The structure of C' follows from its relation to the **similarity transformation** to the corresponding Hermitian model $\widehat{M}^2 = e^{-Q/2} \cdot M^2 \cdot e^{Q/2}$:

$$C' = e^{-Q} \cdot P \qquad e^{-Q} = C' \cdot P = R^2 = \frac{1}{\sqrt{1-\eta^2}} \begin{pmatrix} 1 & \eta \\ \eta & 1 \end{pmatrix}$$

$$Q = \ln R^2 = -\operatorname{arctanh}(\eta) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

[Alexandre, Ellis & PM ’20b]

But we want to work directly with the **non-Hermitian Hamiltonian** …

Variational procedure

The action is not Hermitian, so

$$\frac{\partial \mathcal{L}}{\partial \phi_i^*} - \partial_\nu \frac{\partial \mathcal{L}}{\partial \partial_\nu \phi_i^*} = 0 \;\not\Leftrightarrow\; \frac{\partial \mathcal{L}}{\partial \phi_i} - \partial_\nu \frac{\partial \mathcal{L}}{\partial \partial_\nu \phi_i} = 0$$

Prescription: choose one of these pairs of **Euler-Lagrange equations** to fix the dynamics.

The choices are **physically equivalent**; they are equivalent up to a field redefinition.

[Alexandre, PM & Seynaeve ‘17]

Noether's theorem

If the **Euler-Lagrange equations** are not mutually consistent, conserved currents **do not** correspond to symmetries of the Lagrangian.

$$\delta\mathcal{L} = \underbrace{\left(\frac{\partial\mathcal{L}}{\partial\phi_i} - \partial_\nu\frac{\partial\mathcal{L}}{\partial\partial_\nu\phi_i}\right)}_{\neq 0}\delta\phi_i + \delta\phi_i^*\underbrace{\left(\frac{\partial\mathcal{L}}{\partial\phi_i^*} - \partial_\nu\frac{\partial\mathcal{L}}{\partial\partial_\nu\phi_i^*}\right)}_{=0} + \partial_\nu j_\delta^\nu$$

The current is conserved if

[Alexandre, PM & Seynaeve '17]

$$\delta\mathcal{L} = \left(\frac{\partial\mathcal{L}}{\partial\phi_i} - \partial_\nu\frac{\partial\mathcal{L}}{\partial\partial_\nu\phi_i}\right)\delta\phi_i$$

Noether's theorem

For our scalar model, the **conserved current** is

[Alexandre, PM & Seynaeve '17]

$$j^\nu = i[\phi_1^* \partial^\nu \phi_1 - (\partial^\nu \phi_1^*)\phi_1] - i\,[\phi_2^* \partial^\nu \phi_2 - (\partial^\nu \phi_2^*)\phi_2]$$

corresponding to $(U(1) \times U(1))$

$$\begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \longrightarrow e^{i\alpha P} \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} = \begin{pmatrix} e^{+i\alpha} \phi_1 \\ e^{-i\alpha} \phi_2 \end{pmatrix}$$

with

$$\mathcal{L} \longrightarrow \partial_\nu \phi_i^* \partial^\nu \phi_i - m_i^2 \phi_i^* \phi_i - \mu^2 \left(e^{-2i\alpha} \phi_1^* \phi_2 - e^{+2i\alpha} \phi_2^* \phi_1 \right)$$

$$\delta\mathcal{L} = \left(\frac{\partial \mathcal{L}}{\partial \phi_j} - \partial_\nu \frac{\partial \mathcal{L}}{\partial \partial_\nu \phi_j} \right) \delta\phi_j = 2i\alpha\mu^2 (\phi_1^* \phi_2 - \phi_2^* \phi_1)$$

The Goldstone theorem

The existence of the conserved current is sufficient to ensure **the Goldstone theorem** continues to hold in the case of a **spontaneously broken global symmetry**:

[Alexandre, Ellis, PM & Seynaeve '18]

$$\mathcal{L} = \partial_\nu \phi_i^* \partial^\nu \phi_i + m_1^2|\phi_1|^2 - m_2^2|\phi_2|^2 - \mu^2(\phi_1^*\phi_2 - \phi_2^*\phi_1) - \frac{g}{4}|\phi_1|^4, \qquad m_1^2, m_2^2 > 0$$

$$\left.\begin{aligned} \left.\frac{\partial U}{\partial \phi_1^*}\right|_{\phi_a = v_a} &= \frac{g}{2}|v_1|^2 v_1 - m_1^2 v_1 + \mu^2 v_2 = 0 \\ \left.\frac{\partial U}{\partial \phi_2^*}\right|_{\phi_a = v_a} &= m_2^2 v_2 - \mu^2 v_1 = 0 \end{aligned}\right\} \Rightarrow \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \sqrt{2\frac{m_1^2 m_2^2 - \mu^4}{g m_2^2}} \begin{pmatrix} 1 \\ \frac{\mu^2}{m_2^2} \end{pmatrix} e^{i\alpha}$$

Away from the exceptional point, we have a **single massless, Goldstone mode**.

The Englert-Brout-Higgs mechanism

Gauging the global $U(1)$ symmetry (which is itself subtle):

$$\mathcal{L} = -\frac{1}{4}F_{\nu\rho}F^{\nu\rho} + D_\nu^*\phi_i^*D^\nu\phi_i + m_1^2|\phi_1|^2 - m_2^2|\phi_2|^2 - \mu^2(\phi_1^*\phi_2 - \phi_2^*\phi_1) - \frac{g}{4}|\phi_1|^4$$
$$D_\nu = \partial_\nu - iqA_\nu$$

The **Englert-Brout-Higgs mechanism** is still borne out …

… all the way to the exceptional point:

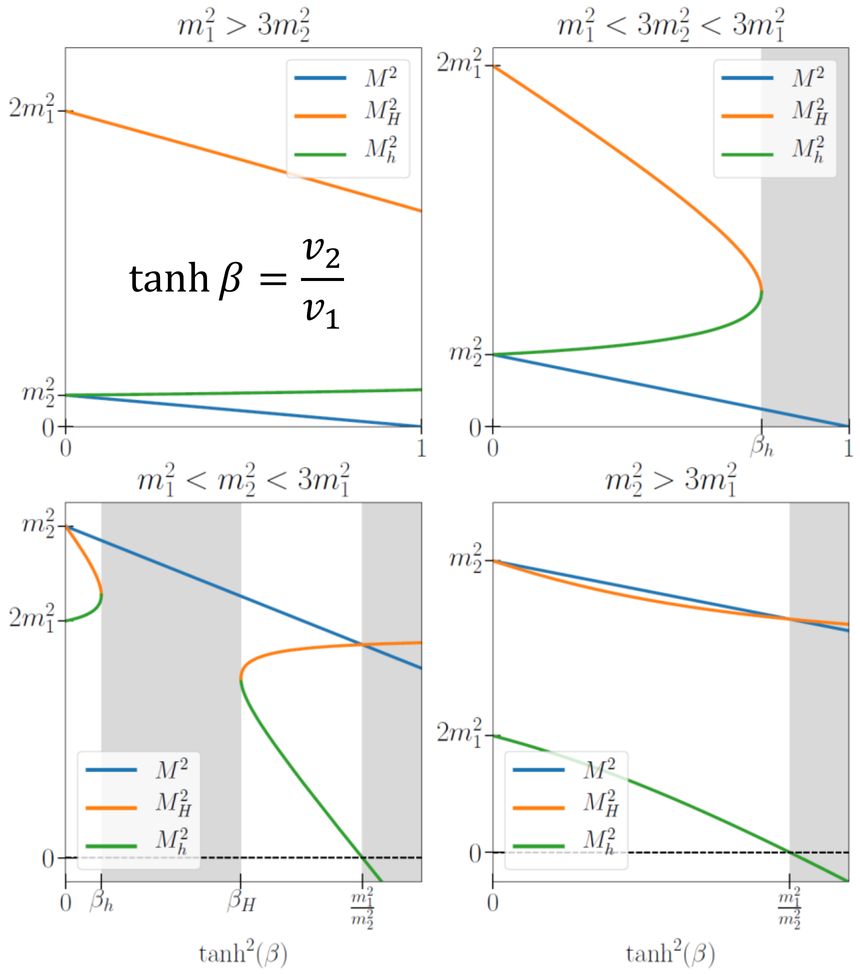
[Alexandre, Ellis, PM & Seynaeve ’19 & ‘20]

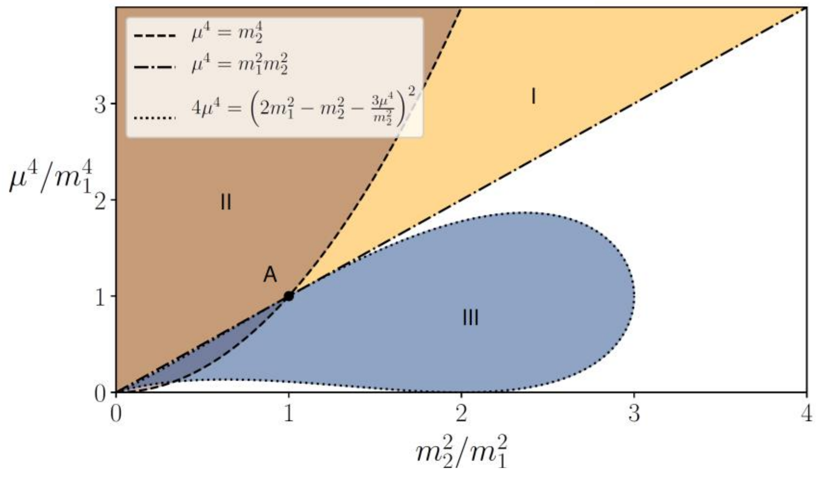
$$M_A^2 = 2q^2\big(|v_1|^2 + |v_2|^2\big) \longrightarrow 2q^2\big(|v_1|^2 + |v_1|^2\big) = \frac{4}{g}\big(m_1^2 - m_2^2\big)\text{ at }\mu^2 = \pm m_2^2$$

The $\mathcal{C}'\mathcal{PT}$ norm of **the Goldstone mode** is ill defined, but its Hermitian norm is defined!

[cf. Mannheim ’19, Fring & Taira ’20a, ‘20b & ’20c, based on a similarity transformation of this model, for which $|v_2|^2 \longrightarrow -|v_2|^2$.]

(Non-Abelian) Englert-Brout-Higgs mechanism

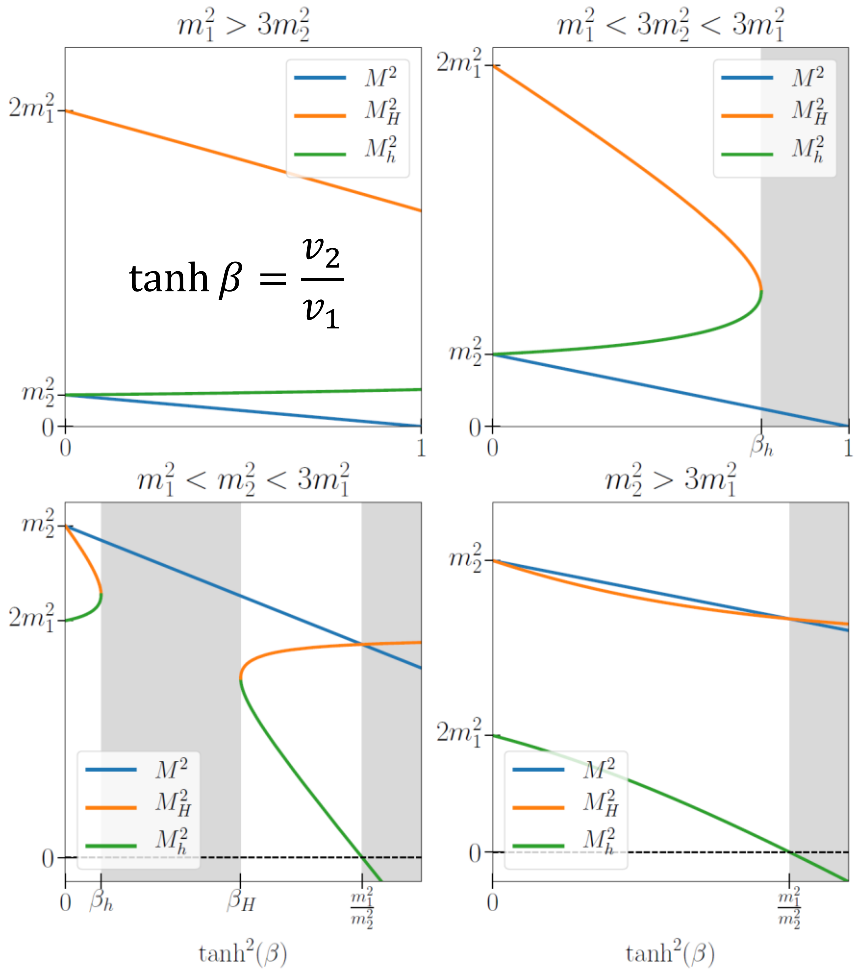




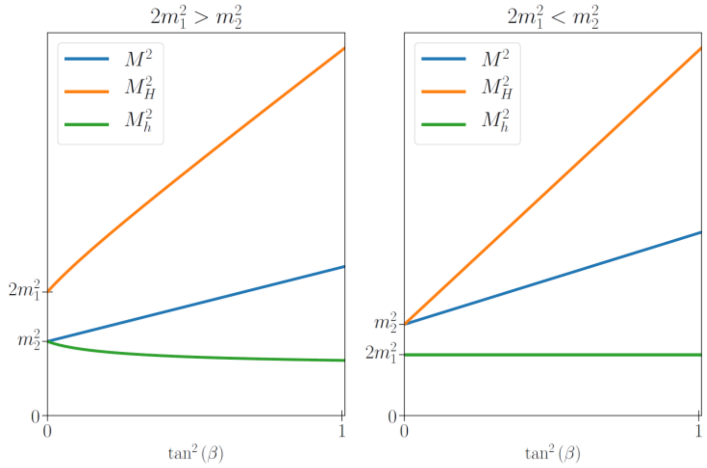
I: symmetric $SU(2) \times U(1)$ phase
II: $\mathcal{PT}$ broken phase ($M^2 < 0$ for $H^\pm, D$)
III: $\mathcal{PT}$ broken phase ($M_h^2, M_H^2 \notin \mathbb{R}$ for h, H)
Unshaded: physical SSB phase

[Alexandre, Ellis, PM & Seynaeve '20]

(Non-Abelian) Englert-Brout-Higgs mechanism



$$\tanh\beta = \frac{v_2}{v_1}$$



cf. Hermitian case

[Alexandre, Ellis, PM & Seynaeve '20]

Canonical variables

$\Phi \equiv (\phi_1, \phi_2)$ evolves with H $\Leftrightarrow$ $\dot{\Phi}^\dagger \equiv (\dot{\phi}_1, \dot{\phi}_2)^\dagger$ evolves with $H^\dagger \neq H$

But canonical variables must both evolve with **the same** H (or $H^\dagger$)!

Canonical variables

The **self-consistent non-Hermitian deformation** is

[Alexandre, Ellis & PM ’20b]

$$\mathcal{L} = \partial_\nu \tilde{\phi}_i^* \partial^\nu \phi_i - m_i^2 \tilde{\phi}_i^* \phi_i - \mu^2 \left(\tilde{\phi}_1^* \phi_2 - \tilde{\phi}_2^* \phi_1 \right)$$

The Euler-Lagrange equations are now mutually consistent:

$$\frac{\partial \mathcal{L}}{\partial \tilde{\phi}_i^*} - \partial_\nu \frac{\partial \mathcal{L}}{\partial \partial_\nu \tilde{\phi}_i^*} = 0 \iff \frac{\partial \mathcal{L}}{\partial \phi_i} - \partial_\nu \frac{\partial \mathcal{L}}{\partial \partial_\nu \phi_i} = 0$$

The **tilde-conjugated fields** are defined via

$$\mathcal{P}\colon\ \phi_i(t, \boldsymbol{x}) \longrightarrow \phi_i'(t, -\boldsymbol{x}) = P_{ij} \tilde{\phi}_j(t, \boldsymbol{x})$$

But there is still a choice!

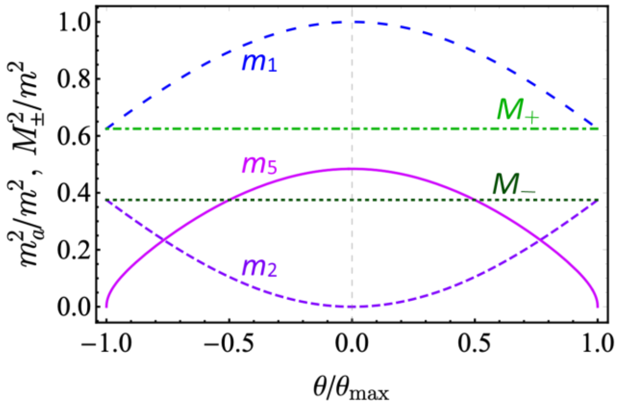
A unique feature: a local similarity map

[Chernodub & PM '21]

Choose a **local non-Hermitian mass matrix** with globally constant eigenvalues ($\mu^2 \to m_5^2$):

$$\left.\begin{array}{c} m_1^2(x) + m_2^2(x) = M_0^2 \\ [m_1^2(x) - m_2^2(x)]^2 - 4m_5^4(x) = m_0^4 \end{array}\right\} M_\pm^2 = \frac{1}{2}\left(M_0^2 \pm m_0^2\right)$$

$$\theta(x) = \arctan\frac{m_2(x)}{m_1(x)}$$



The mass matrix is diagonalised by a **local similarity transformation**, leading to:

$$\mathcal{L}_{C,\mathrm{H}} = \partial_\mu\tilde{\phi}_+^*\partial^\mu\phi_+ - M_+^2\tilde{\phi}_+^*\phi_+ + \partial_\mu\tilde{\phi}_-^*\partial^\mu\phi_- - M_-^2\tilde{\phi}_-^*\phi_- + C_\mu C^\mu\left(\tilde{\phi}_+^*\phi_+ + \tilde{\phi}_-^*\phi_-\right) + C_\mu j^\mu$$

$$j^\mu = -\tilde{\phi}_+^*\partial^\mu\phi_- + \phi_+\partial^\mu\tilde{\phi}_-^* - \tilde{\phi}_-^*\partial^\mu\phi_+ + \phi_-\partial^\mu\tilde{\phi}_+^* - 2C^\mu\left(\tilde{\phi}_+^*\phi_+ + \tilde{\phi}_-^*\phi_-\right)$$

C^μ is the **similarity "gauge" field**.

A unique feature: a local similarity map

[Chernodub & PM ‘21]

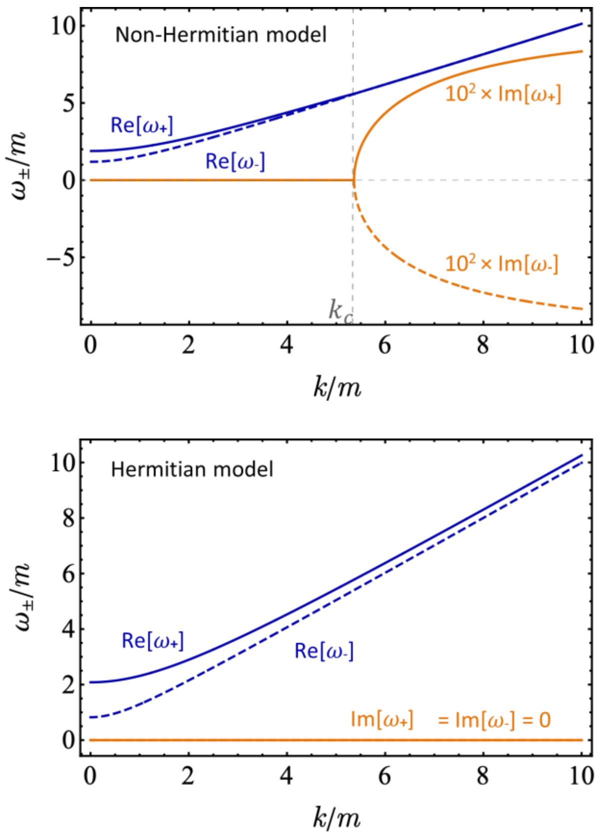
Energy spectrum for a constant similarity “gauge” field:

$$\left(\omega^2 - \boldsymbol{k}^2 - C^2 - m_1^2\right)\left(\omega^2 - \boldsymbol{k}^2 - C^2 - m_2^2\right) + 4(C^0\omega - \boldsymbol{k}\cdot\boldsymbol{C})^2 + m_5^4 = 0$$

Modes with $k < k_c$ are in the $\mathcal{PT}$ **unbroken** regime; modes with $k > k_c$ are $\mathcal{PT}$ **broken**. For the **time-like case**,

$$k_c^2 = \frac{(m_1^2 - m_2^2)^2 - 4m_5^4}{16C_0^2} - \frac{m_1^2 + m_2^2}{2}$$

Novel IR/UV connection.



Second quantisation

In the flavour basis, the matrix-valued energy is non-Hermitian, and we need:

[see Alexandre, Ellis & PM ’20b, also for full details of how the discrete symmetry properties $\mathcal{C}, \mathcal{C}', \mathcal{P}, \mathcal{T}$ are borne out.]

$$\hat{\phi}_i(x) = \int \frac{d^3\boldsymbol{p}}{(2\pi)^3} [2E(\boldsymbol{p})]_{ij}^{-1/2} \Big[\big(e^{-ip\cdot x}\big)_{jk} \hat{a}_{k,\boldsymbol{p}}(0) + \big(e^{ip\cdot x}\big)_{jk} \check{c}^\dagger_{k,\boldsymbol{p}}(0) \Big]$$

$$\check{\phi}_i^\dagger(x) = \int \frac{d^3\boldsymbol{p}}{(2\pi)^3} [2E(\boldsymbol{p})]_{ji}^{-1/2} \Big[\big(e^{-ip\cdot x}\big)_{kj} \hat{c}_{k,\boldsymbol{p}}(0) + \big(e^{ip\cdot x}\big)_{kj} \check{a}^\dagger_{k,\boldsymbol{p}}(0) \Big]$$

The **hatted** (ˆ) and **checked** (ˇ) **fields** are related via **parity**:

$$P_{ij}\check{\phi}_j(\mathcal{P}x) = \hat{\mathcal{P}}\hat{\phi}_i(x)\hat{\mathcal{P}}^{-1}$$

The **second-quantised Lagrangian** is

$$\hat{\mathcal{L}} = \partial_\nu \check{\phi}_i^\dagger \partial^\nu \hat{\phi}_i - m_i^2 \check{\phi}_i^\dagger \hat{\phi}_i - \mu^2 \left(\check{\phi}_1^\dagger \hat{\phi}_2 - \check{\phi}_2^\dagger \hat{\phi}_1 \right)$$

Flavour oscillations

Mass eigenstates: $|\boldsymbol{p}, +(-), t\rangle, \quad (|\boldsymbol{p}, +(-), t\rangle)^{\S}, \quad \S \equiv \mathcal{C}'\mathcal{P}\mathcal{T} \circ \mathrm{T}$

Flavour eigenstates: [Alexandre, Ellis & PM ‘20b]

$$|\breve{\boldsymbol{p}}, 1(2), t\rangle = N\left\{\eta|\boldsymbol{p}, +(-), t\rangle - \left[1 - \sqrt{1-\eta^2}\right]|\boldsymbol{p}, -(+), t\rangle\right\}$$

$$\langle\widehat{\boldsymbol{p}}, 1(2), t| = N\left\{\eta(|\boldsymbol{p}, +(-), t\rangle)^{\S} + \left[1 - \sqrt{1-\eta^2}\right](|\boldsymbol{p}, -(+), t\rangle)^{\S}\right\}$$

Orthonormality:

$$\langle\widehat{\boldsymbol{p}}, i, t|\breve{\boldsymbol{p}}', j, t\rangle = (2\pi)^3\delta_{ij}\delta^3(\boldsymbol{p} - \boldsymbol{p}')$$

$$(|\boldsymbol{p}, \pm, t\rangle)^{\S}|\boldsymbol{p}', \pm, t\rangle = (2\pi)^3\delta^3(\boldsymbol{p} - \boldsymbol{p}')$$

$$(|\boldsymbol{p}, \pm, t\rangle)^{\S}|\boldsymbol{p}', \mp, t\rangle = 0$$

Flavour oscillations

Transition "probability":

$$\Pi_{i\to j}(t) = \frac{1}{V}\int \frac{\mathrm{d}^3\boldsymbol{p}'}{(2\pi)^3}\langle \hat{\boldsymbol{p}}, j, t|\check{\tilde{\boldsymbol{p}}}', i, 0\rangle\langle \hat{\boldsymbol{p}}', i, 0|\check{\tilde{\boldsymbol{p}}}, j, t\rangle, \qquad V \equiv (2\pi)^3\delta^3(\mathbf{0})$$

Oscillation and **"survival probabilities":** [Alexandre, Ellis & PM '20b]

$$\Pi_{1(2)\to 2(1)}(t) = -\frac{\eta^2}{1-\eta^2}\sin^2\left[\frac{1}{2}\big(E_+(\boldsymbol{p}) - E_-(\boldsymbol{p})\big)t\right] \notin [0,1]$$

$$\Pi_{1(2)\to 1(2)}(t) = 1 + \frac{\eta^2}{1-\eta^2}\sin^2\left[\frac{1}{2}\big(E_+(\boldsymbol{p}) - E_-(\boldsymbol{p})\big)t\right] \notin [0,1]$$

Unitarity: $\Pi_{1(2)\to 1(2)}(t) + \Pi_{1(2)\to 2(1)}(t) = 1$

[cf. the similar issue found in Ohlsson & Zhou '20 & '21]

Flavour oscillations

Resolution: experimental **observables** are **scattering matrix elements**.

[Alexandre, Ellis & PM '20b]

We must source states consistent with the $\mathcal{PT}$ symmetry:

$$\mathcal{L}_{\text{int}} = J_A \check{\phi}_1^\dagger + J_A^\dagger \hat{\phi}_1 - J_B \check{\phi}_2^\dagger + J_B^\dagger \hat{\phi}_2$$

"Squared" matrix element for $A \longrightarrow B$:

$$\mathcal{M}_{A\longrightarrow B}^{\mathcal{C}'\mathcal{PT}} \mathcal{M}_{A\longrightarrow B} = VT(2\pi)^4 \delta^4(p_A - p_B) \frac{\mu^4}{\left(p_A^2 - M_+^2\right)^2 \left(p_A^2 - M_-^2\right)^2} > 0$$

remaining real and perturbative all the way up to the exceptional point ($M_+^2 = M_-^2$).

What about fermions?

An example: non-Hermitian extension of the **Dirac theory**

[Bender, Jones & Rivers ’05; Alexandre, Bender & PM ‘15; Alexandre, PM & Seynaeve ‘17]

$$\mathcal{L} = \bar{\psi}(i\gamma^\nu \partial_\nu - m - \mu\gamma^5)\psi, \qquad \gamma^5 = (\gamma^5)^\dagger$$

Eigenmasses: $M^2 = m^2 - \mu^2$

The **conserved current** is [Alexandre & Bender ‘15]

$$j^\nu = \bar{\psi}\gamma^\nu\left(1 + \frac{\mu}{m}\gamma^5\right)\psi = \left(1 - \frac{\mu}{m}\right)\psi_L^\dagger \bar{\sigma}^\nu \psi_L + \left(1 + \frac{\mu}{m}\right)\psi_R^\dagger \bar{\sigma}^\nu \psi_R$$

corresponding to

$$\psi \longrightarrow \psi' = \exp\left[i\alpha\left(1 + \frac{\mu}{m}\gamma^5\right)\right]\psi$$

with

$$\delta\mathcal{L} = -2\mu\bar{\psi}\gamma^5\delta\psi \neq 0$$

What about fermions?

- **Exceptional points** $\mu = +(-)m \implies$ a **massless theory** of right(left) chiral Weyl fermions.
  [Alexandre, Bender & PM '15, cf. Chernodub '17]
- Gauging this model, the full **vector plus axial vector gauge symmetry** is recovered at the **exceptional points**.
  [Alexandre, Bender & PM '15; Alexandre, PM & Seynaeve '17]
- Massless fermions can undergo **flavour oscillations**.
  [Jones-Smith & Mathur '14]
- The same model can be obtained from a **non-Hermitian Higgs-Yukawa theory**.
  [Alexandre, Bender & PM '15 & '17; see also Alexandre & Mavromatos '20]

$$\mathcal{L}_{\mathrm{Yuk}} = -y_- \bar{L}_L \widetilde{H} \nu_R - y_+ \bar{\nu}_R \widetilde{H}^\dagger L_L$$

- A non-Hermitian explanation for the smallness of the **light neutrino masses**?

SUSY embedding?

Two $\mathcal{N} = 1$ scalar chiral superfields: Φ_1, Φ_2

Superpotential: [Alexandre, Ellis & PM '20a]

$$W_\pm = \frac{1}{2} m_{11} \Phi_1^2 \mp m_{12} \Phi_1 \Phi_2 + \frac{1}{2} m_{22} \Phi_2^2, \qquad \mathcal{L} = \mathcal{L}_K + \int d^2\theta \; W_+ + \int d^2\theta^\dagger \, W_-^\dagger \neq \mathcal{L}^\dagger$$

On-shell (two complex scalars, two Majorana fermions, $a = 1,2$):

$$\mathcal{L}_{\text{scal}} = \partial_\nu \phi_a^* \partial^\nu \phi_a - \left(m_{aa}^2 - m_{12}^2\right) \phi_a^* \phi_a - m_{12}(m_{22} - m_{11})(\phi_1^* \phi_2 - \phi_2^* \phi_1)$$

$$\mathcal{L}_{\text{ferm}} = \frac{1}{2} \bar{\psi}_a i \gamma^\nu \partial_\nu \psi_a - \frac{1}{2} m_{aa} \bar{\psi}_a \psi_a - \frac{1}{2} m_{12} \left(\bar{\psi}_1 \gamma^5 \psi_2 + \bar{\psi}_2 \gamma^5 \psi_1\right), \qquad \gamma^5 = \left(\gamma^5\right)^\dagger$$

But $M_{\text{scal},\pm}^2 \neq M_{\text{ferm},\pm}^2 \;\Rightarrow$ supersymmetry breaking!

Closing remarks

- **Noether's theorem**, **Goldstone theorem** and **Englert-Brout-Higgs mechanism** borne out.
- **Unique features** and **parametric dependence** very different to similar Hermitian models.
- **Second quantisation** in the non-Hermitian "frame" is subtle.
- Formulate **non-Hermitian flavour oscillations** consistent with perturbative unitarity.
- A new possibility for **SUSY breaking**.
- Potential implications for the **neutrino sector**.
- **CP violation?** [Dale, Mason & PM in prep]; **fermionic second quantisation?** [Alexandre, Ellis & PM in prep].

Thank you and stay well

Questions or **comments**?

Message me on **twitter @pwmillington** or **email** me at **p.millington@nottingham.ac.uk**.

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Please see references therein.

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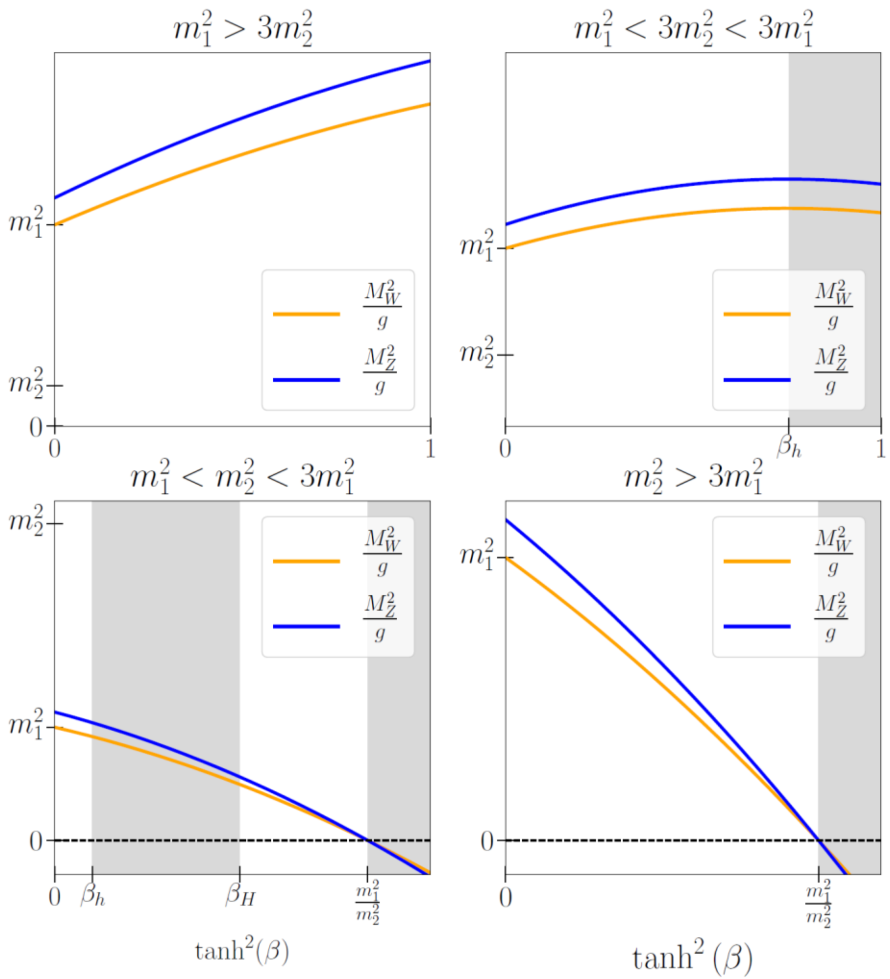
Back up slides

(Non-Abelian) Englert-Brout-Higgs mechanism



[Alexandre, Ellis, PM & Seynaeve '20]

Second quantisation

Why this doubling?

- To ensure a **consistent variational procedure**, and to construct **canonical conjugate variables**.
- In terms of the **mass "eigenfields"** or the **similarity transformation** to the Hermitian "frame":

$$\hat{\xi}_i = R_{ij}\hat{\phi}_j \not\Leftrightarrow \hat{\xi}_i^\dagger = \hat{\phi}_j^\dagger R_{ji}^{-1}$$

instead:

$$\hat{\xi}_i = R_{ij}\hat{\phi}_j \Leftrightarrow \hat{\xi}_i^\dagger = \check{\phi}_i^\dagger R_{ji}^{-1}$$

Maxwell equations

Since we must couple to a **non-conserved current**

$$\partial_\nu j^\nu = -2iq\mu^2(\phi_1^*\phi_2 - \phi_2^*\phi_1)$$

The **Maxwell equations** are inconsistent, since $\partial_\nu\partial_\rho F^{\nu\rho} = 0$ identically.

Resolution: Adding a covariant **gauge fixing term** is sufficient: [Alexandre, Ellis, PM & Seynaeve '19]

$$\mathcal{L} \supset -\frac{1}{2\zeta}(\partial_\nu A^\nu)^2$$

This leads to the constraint (cf. Stueckelberg case)

$$\frac{1}{\zeta}\Box\,\partial_\nu A^\nu = -2iq\mu^2(\phi_1^*\phi_2 - \phi_2^*\phi_1)$$

precluding the **Lorenz gauge condition**.